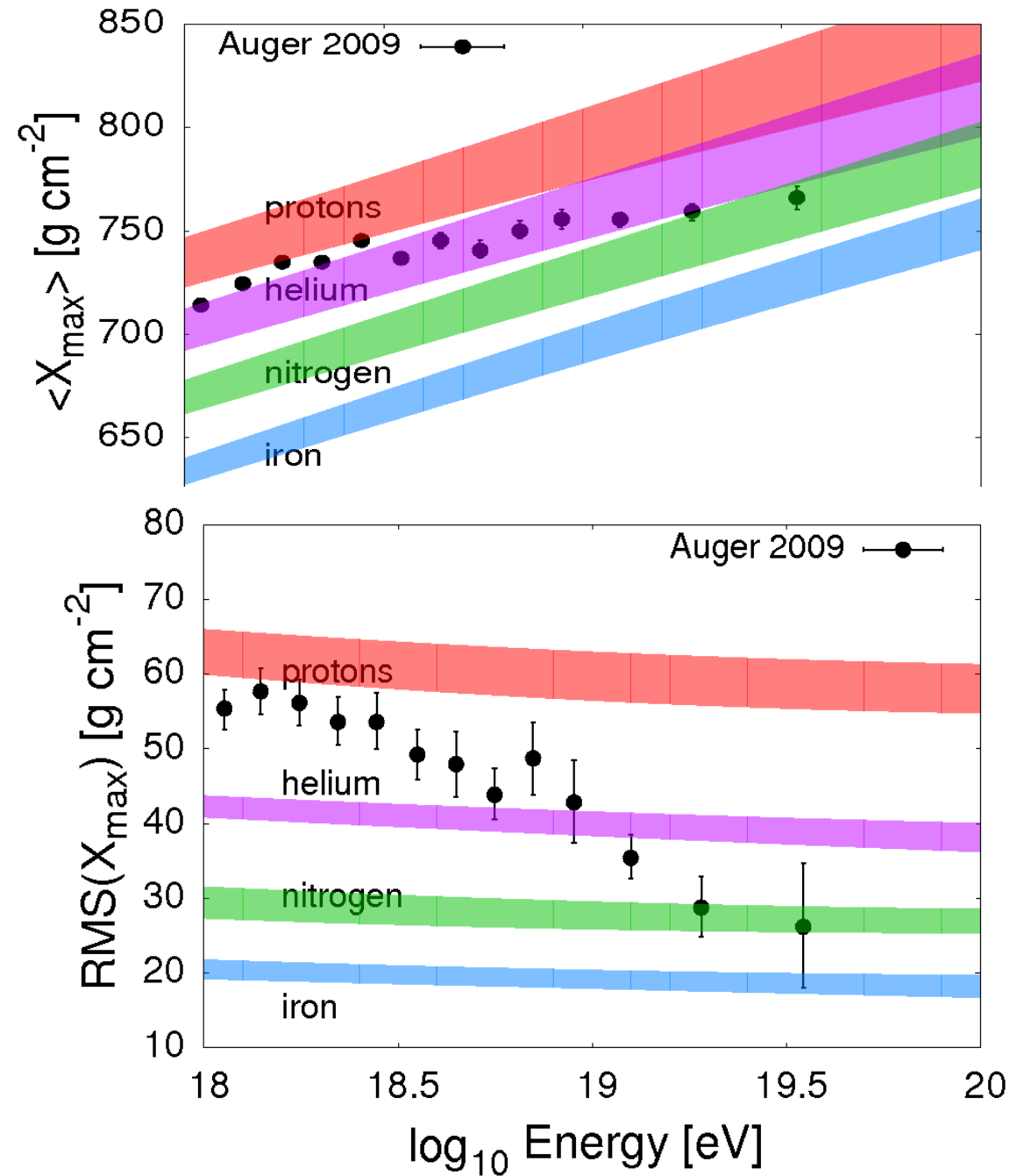


# Implications of Auger Composition Results



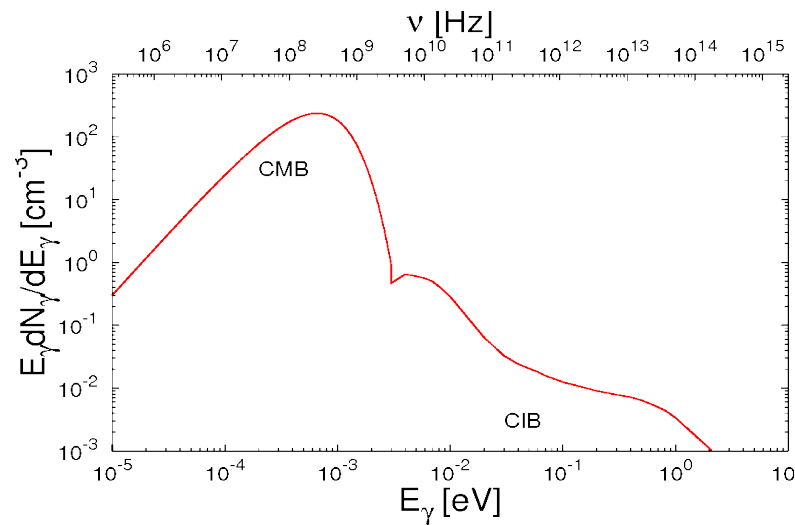
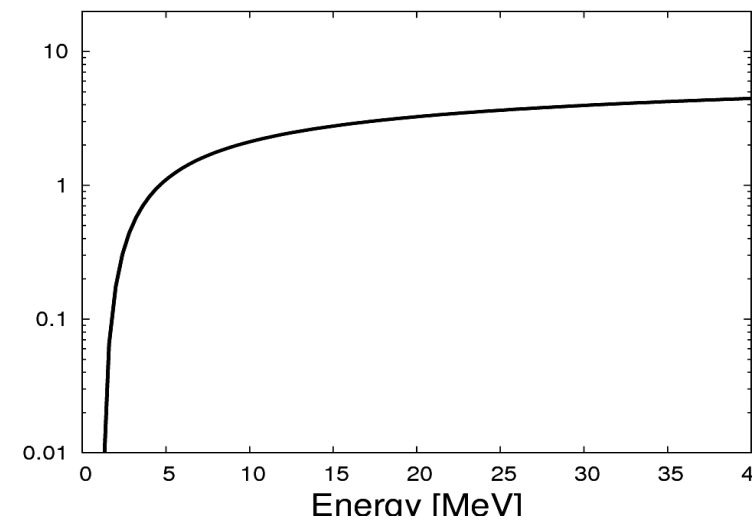
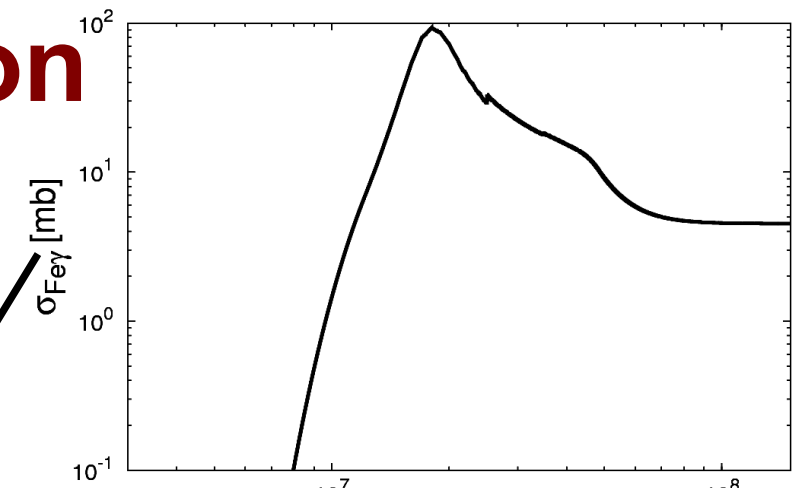
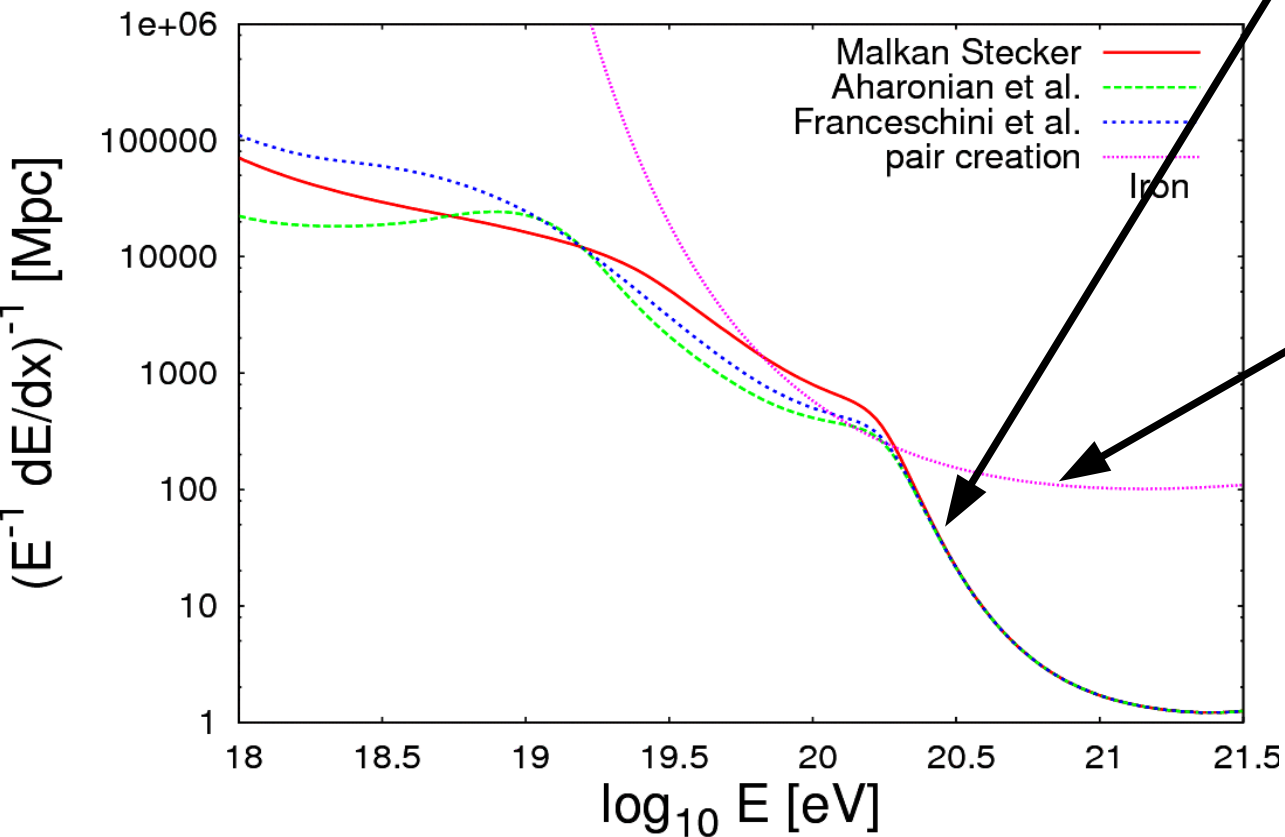
1) What source injection spectra + composition is consistent with Auger Results?

2) What constraints do presence of nuclei place on the local sources?

3) Local VHE emitters which might fit the bill

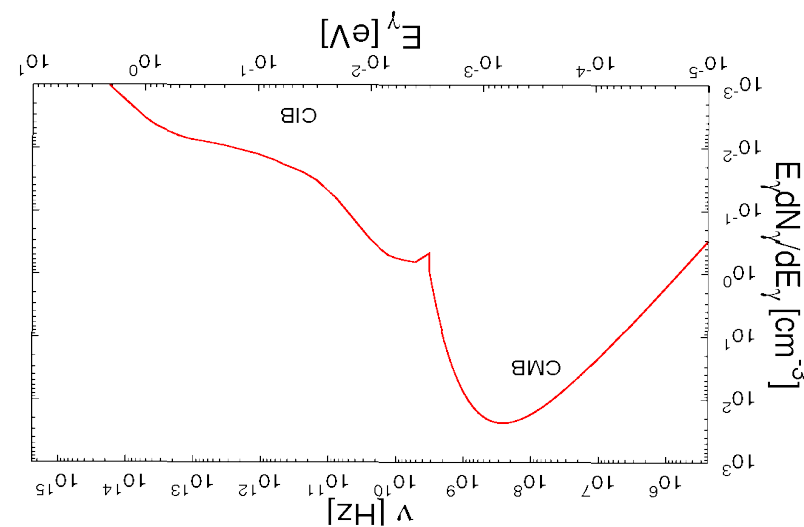
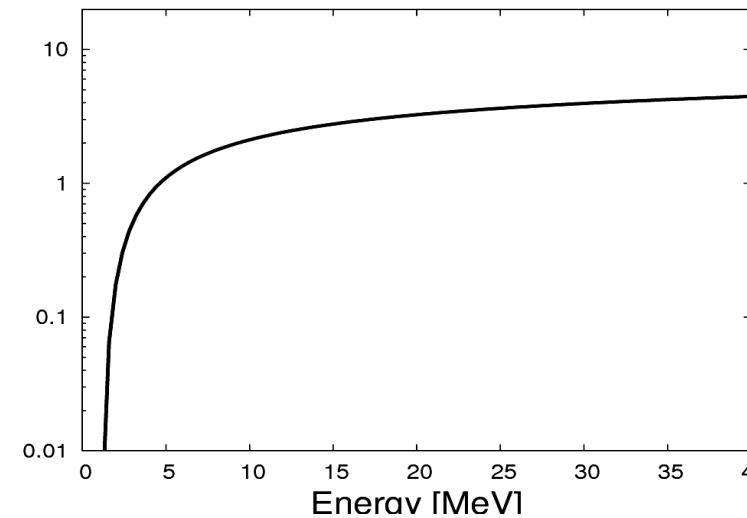
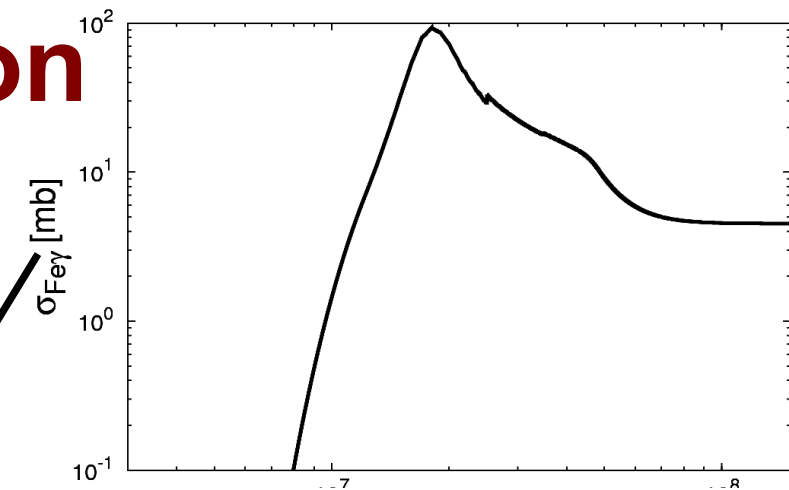
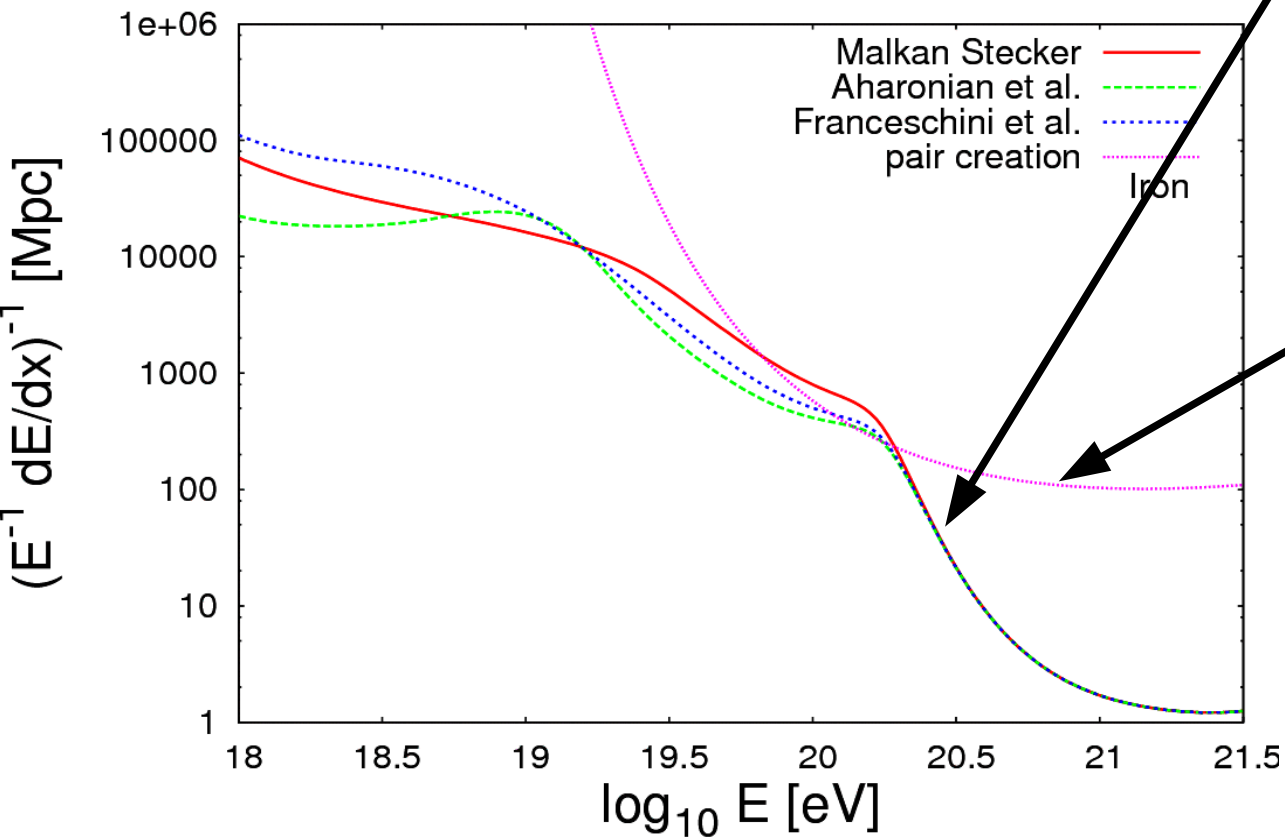
4) What hope do we have to identify these sources?

# UHECR Nuclei Propagation From Source to Earth



Andrew Taylor

# UHECR Nuclei Propagation From Source to Earth



Andrew Taylor

# 1<sup>st</sup> Order: Cascade of Nuclei Through Species- single nucleon loss

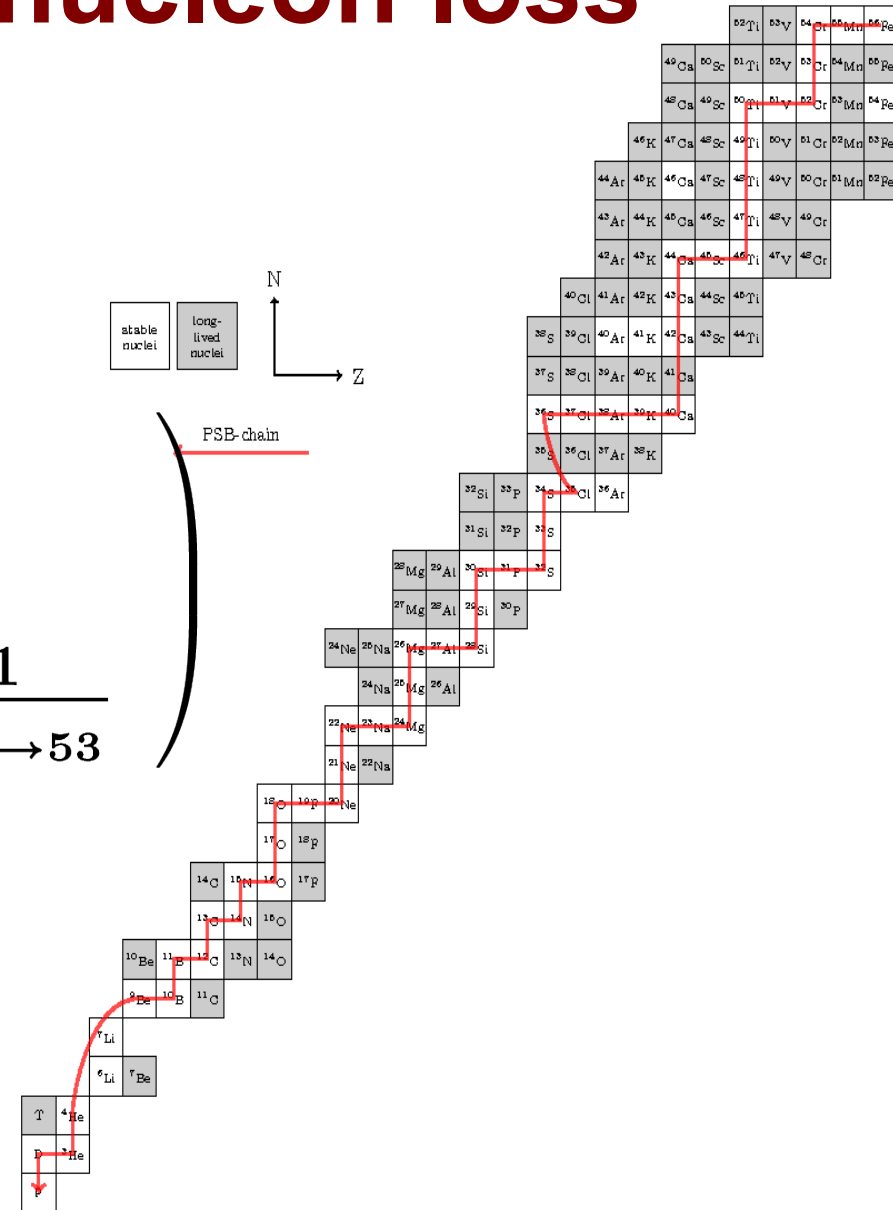
$$\frac{d}{dt} \begin{pmatrix} f_{56} \\ f_{55} \\ f_{54} \end{pmatrix} = \Lambda \begin{pmatrix} f_{56} \\ f_{55} \\ f_{54} \end{pmatrix}$$

$$\Lambda = \begin{pmatrix} -\frac{1}{\tau_{56 \rightarrow 55}} & 0 & 0 \\ \frac{1}{\tau_{56 \rightarrow 55}} & -\frac{1}{\tau_{55 \rightarrow 54}} & 0 \\ 0 & \frac{1}{\tau_{55 \rightarrow 54}} & -\frac{1}{\tau_{54 \rightarrow 53}} \end{pmatrix}$$

Whose eigenvalues are

$$f_q = \sum_{n=q}^{56} \frac{\tau_q \tau_n^{56-q-1}}{\prod_{p=q}^{56} (\tau_n - \tau_p)} e^{-t/\tau_n}$$

Andrew Taylor





# Cascade of Nuclei Through Species- single nucleon loss

Since nuclei Lorentz factor remains  
~conserved, and cross-section varies mildly  
with A (nuclear mass)

$$\tau_{56 \rightarrow 55} \approx \tau_{55 \rightarrow 54} \dots$$

For the case  $\tau_{56 \rightarrow 55} = \tau_{55 \rightarrow 54} \dots$

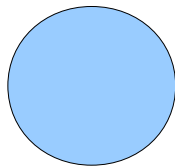
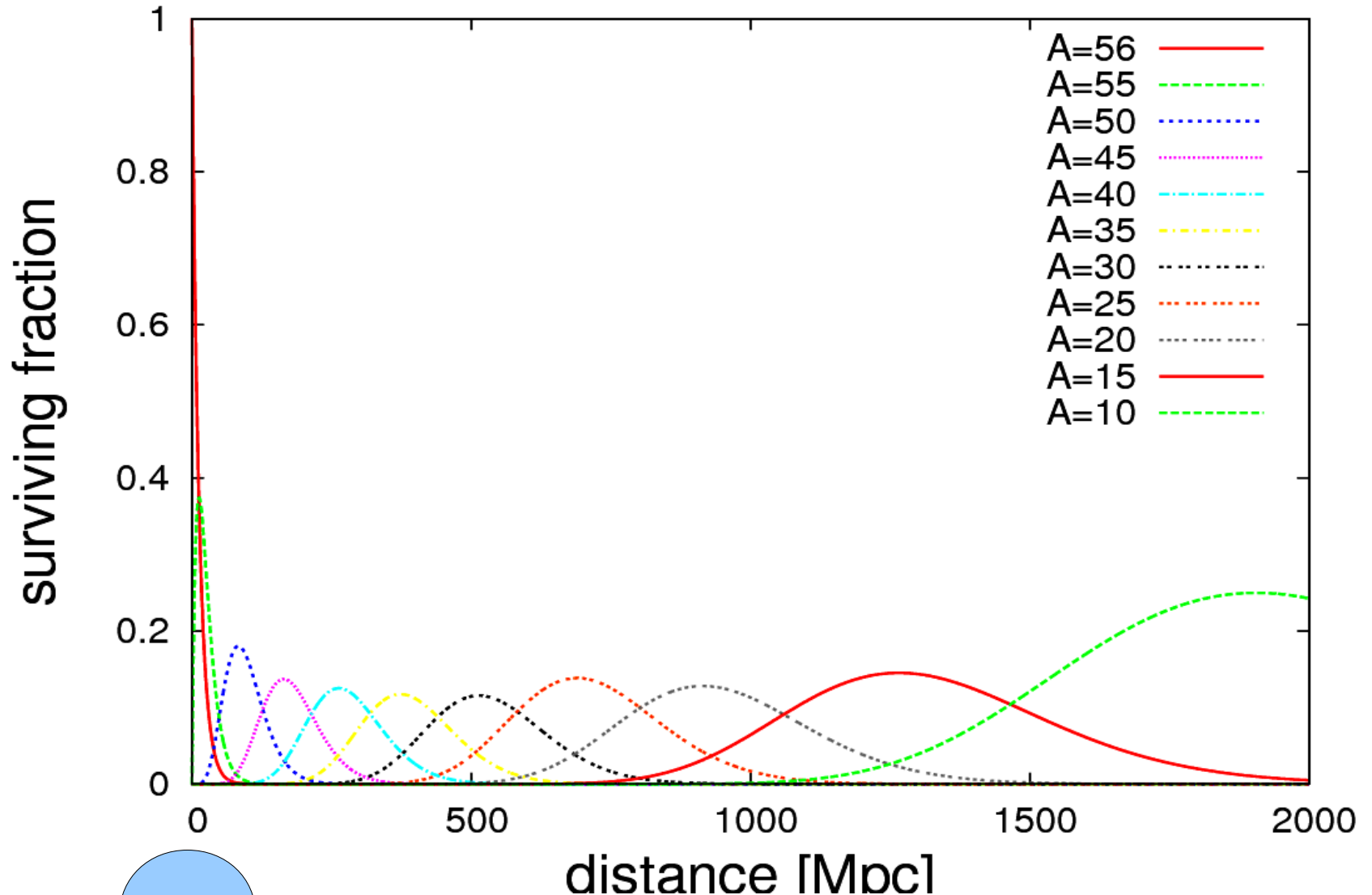
$$f_q = \frac{t^{(q_{max} - q)}}{\tau_q (q_{max} - q)!} e^{-t/\tau_q}$$

ie. Gaisser-Hillas  
type function!

(used to describe air showers)

# Nuclei Propagation Away from their Source + their Transmutation

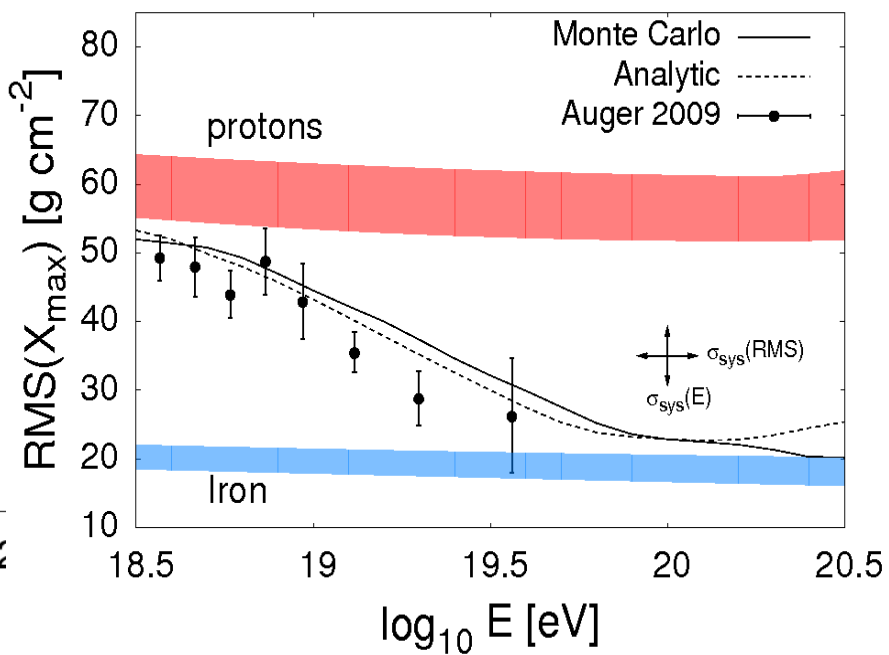
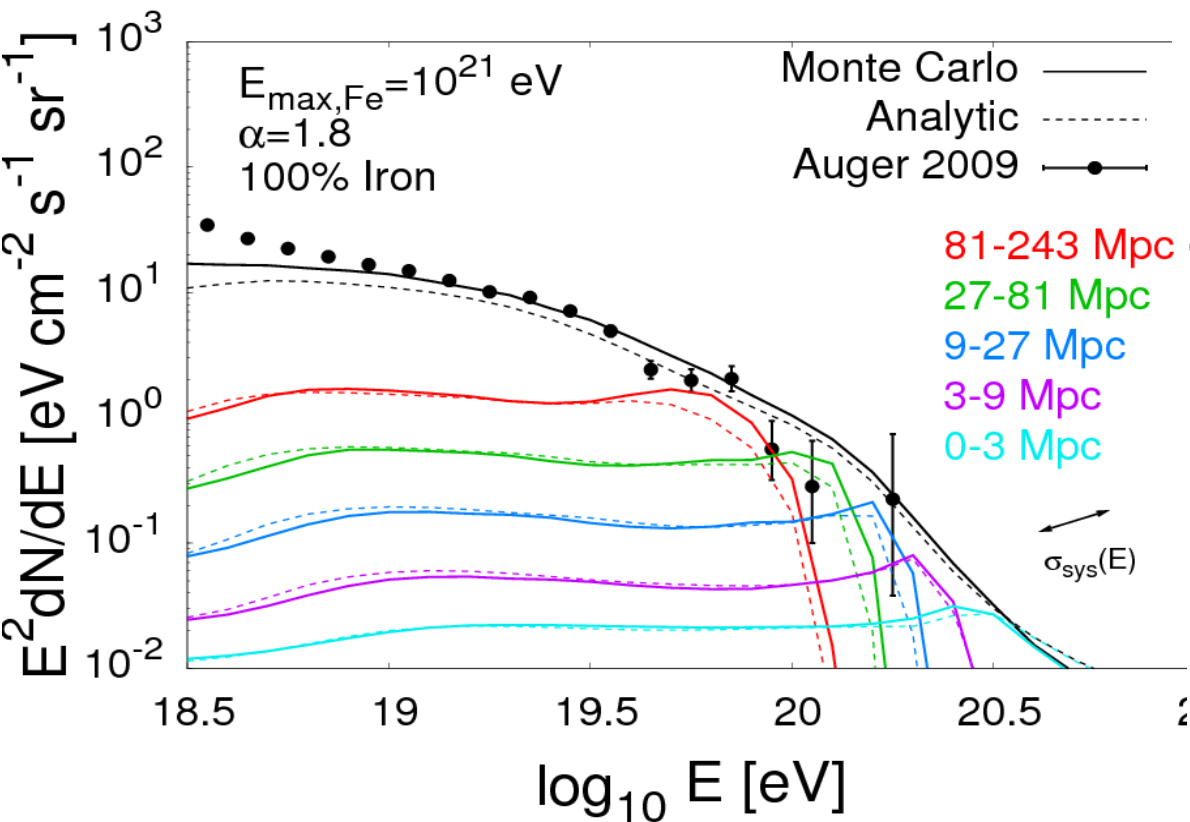
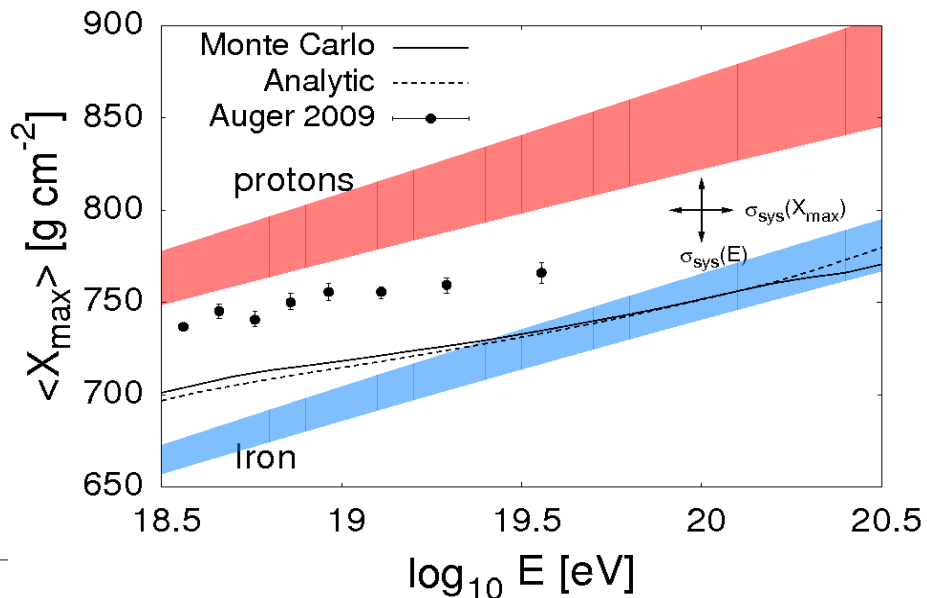
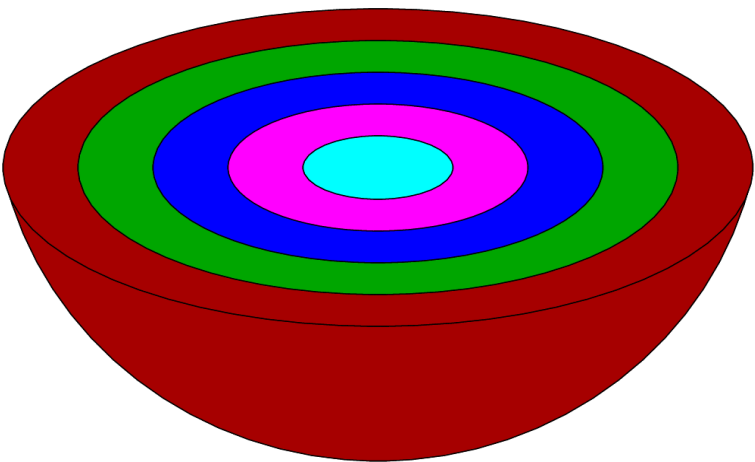
For  $10^{20}$  eV Iron Nuclei at source-



# Local Scales Effect Highest Energies

0 3 9 27 81 243 Mpc

(logarithmic scale)



# What is the Source Composition?

Keep It Simple

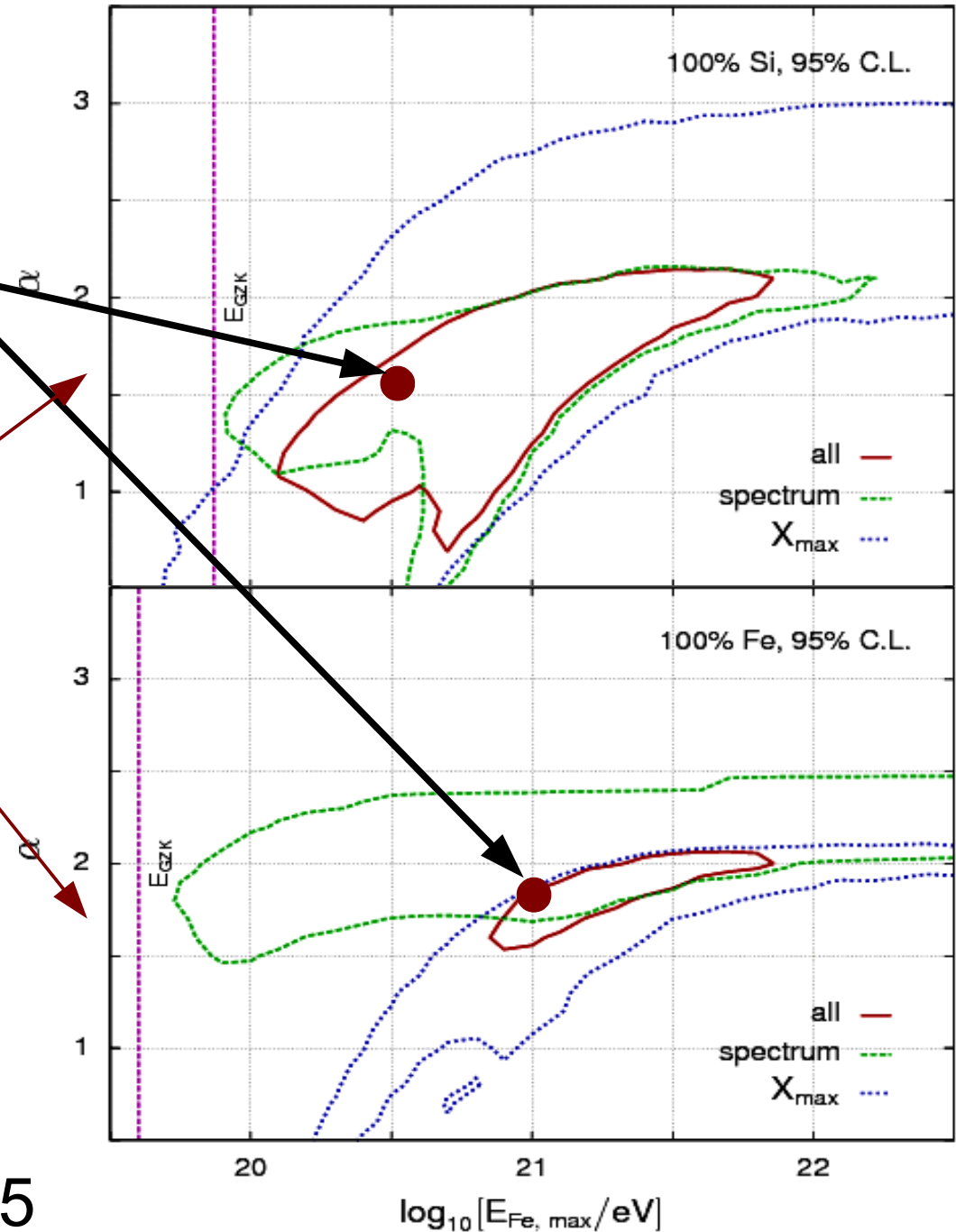
- **Single Composition**

**Example Best-Fit Models**

**Silicon** →

**Hard Spectra preferred**

**Iron** →



astro-ph/1107.2055

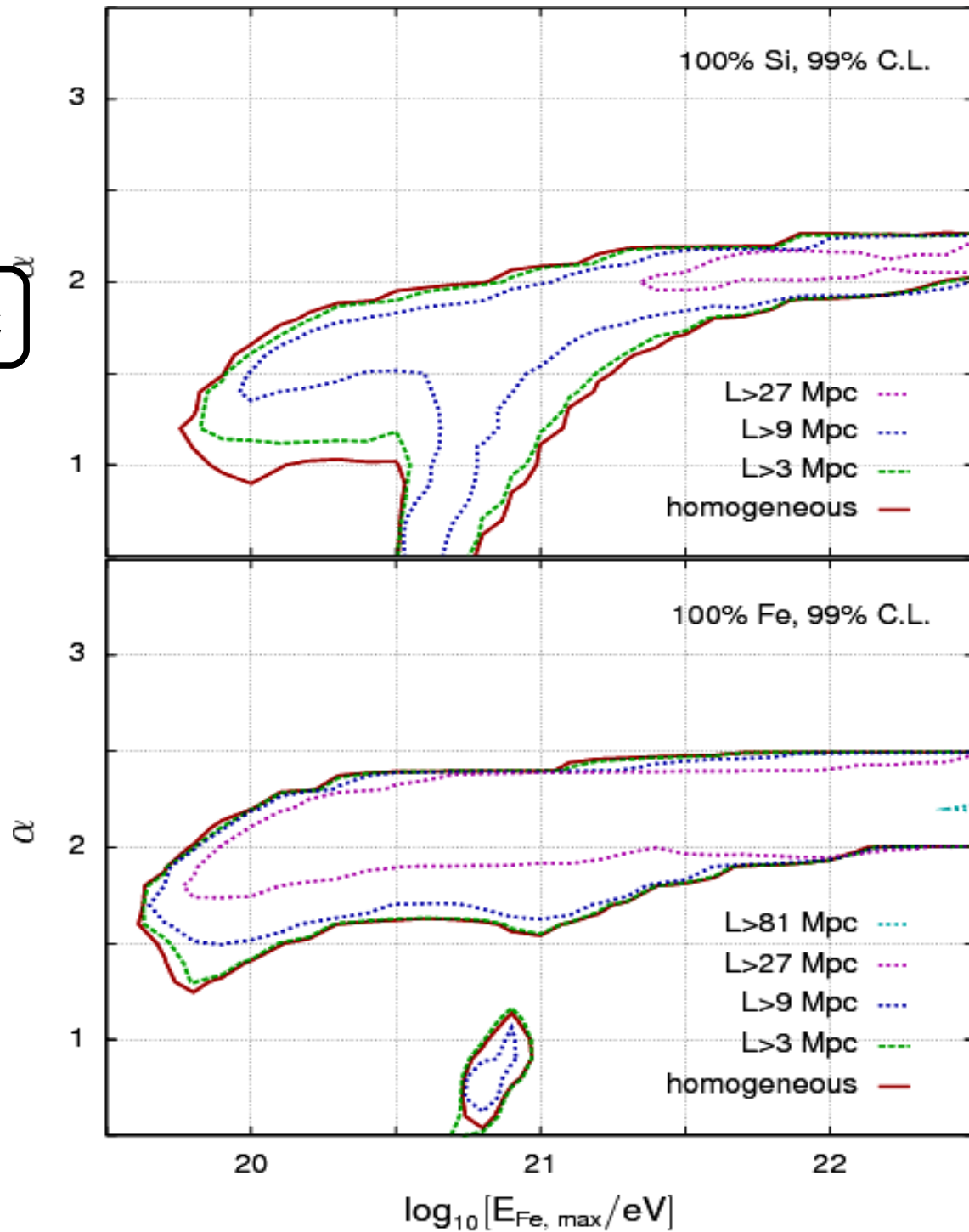
# How Far is the Nearest Source?

If  $E_{\text{max}} < 10^{22} \text{ eV}$

**Silicon-  $D < 60 \text{ Mpc}$**

**Iron-  $D < 80 \text{ Mpc}$**

astro-ph/1107.2055



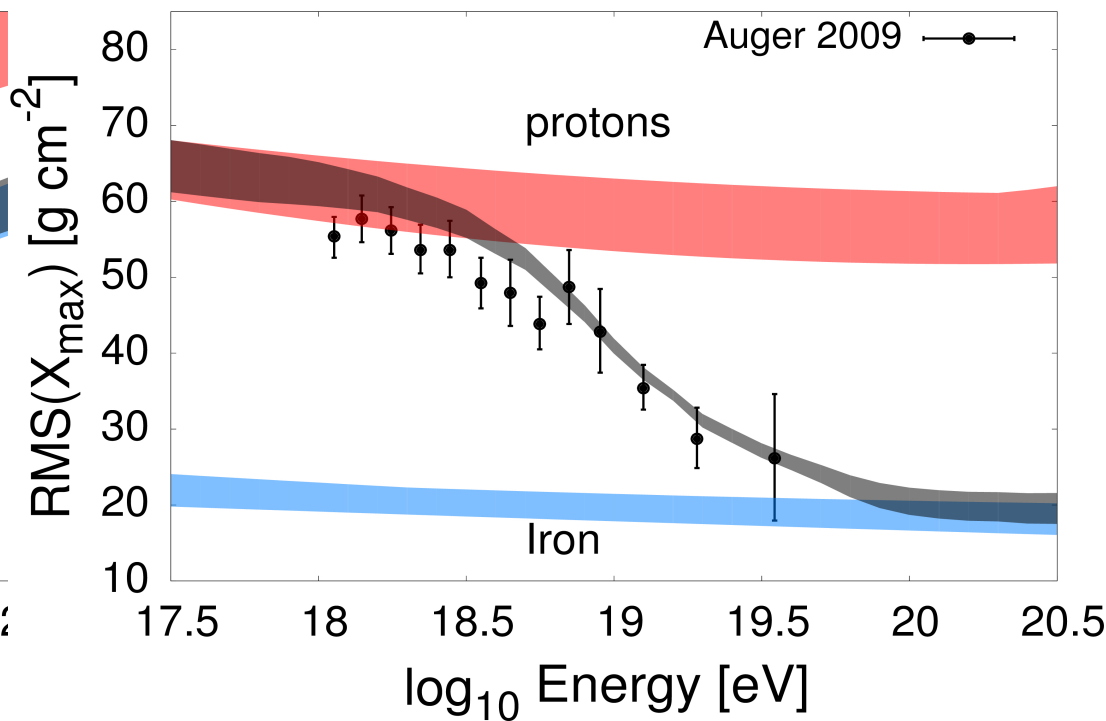
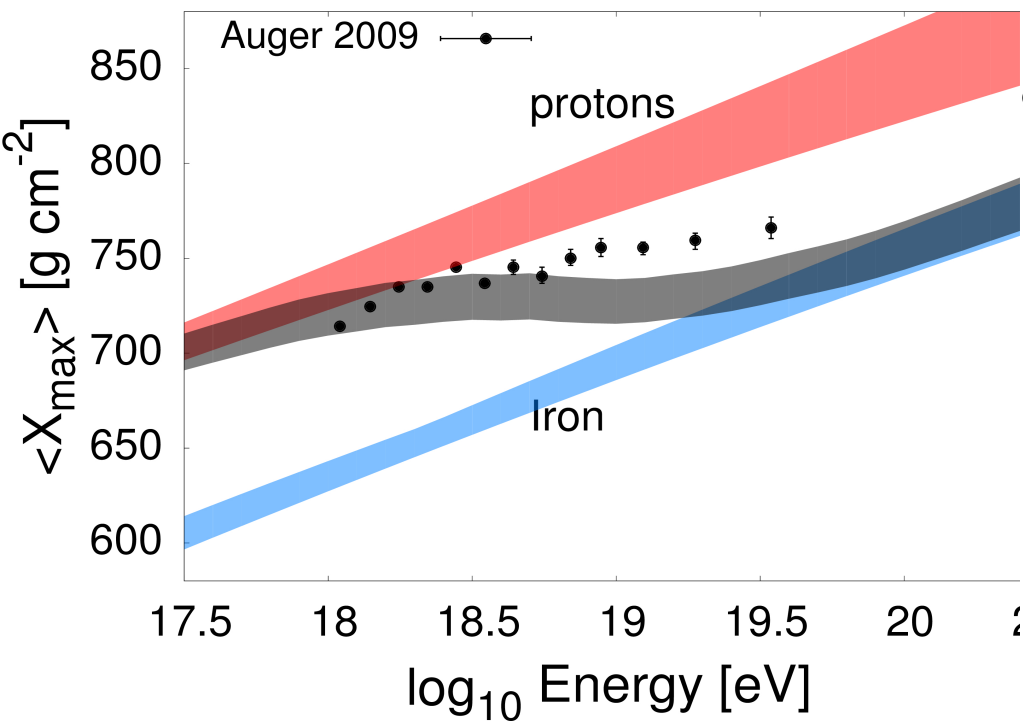
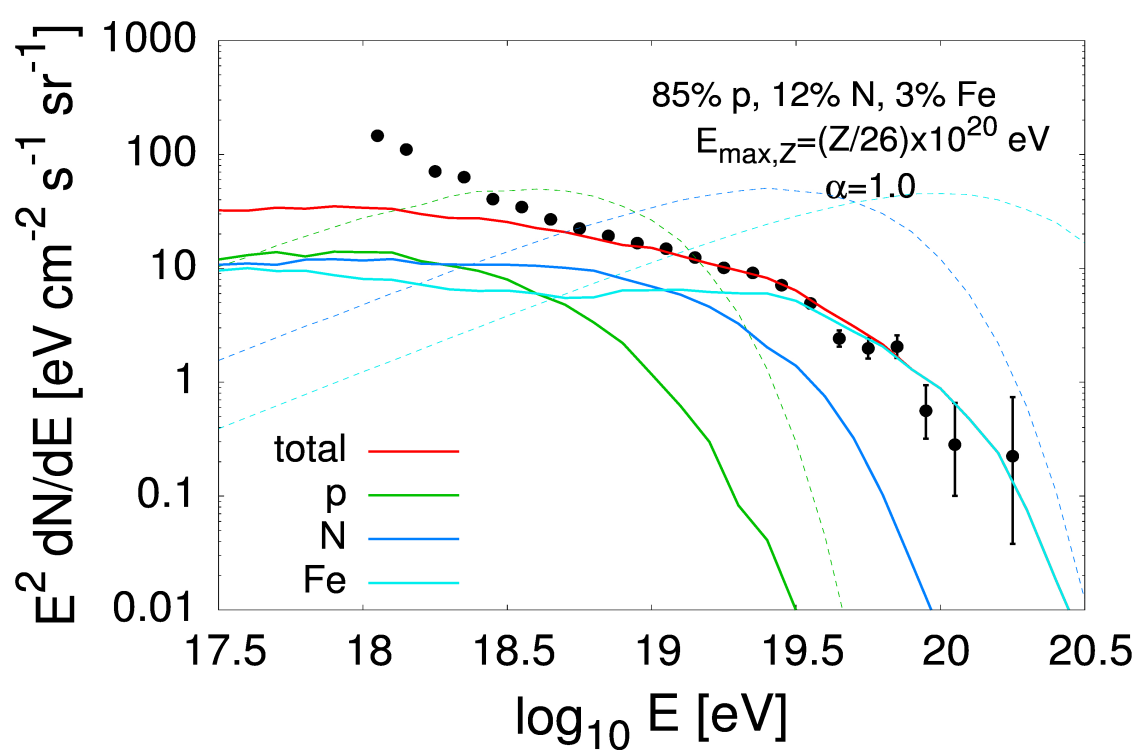
# Mixed Composition

Archetype 1: Harder  
Low Cutoff

astro-ph/1401.0199

For  $E_{\text{max}} = 10^{20}$  eV

85%p, 12%N, 3%Fe



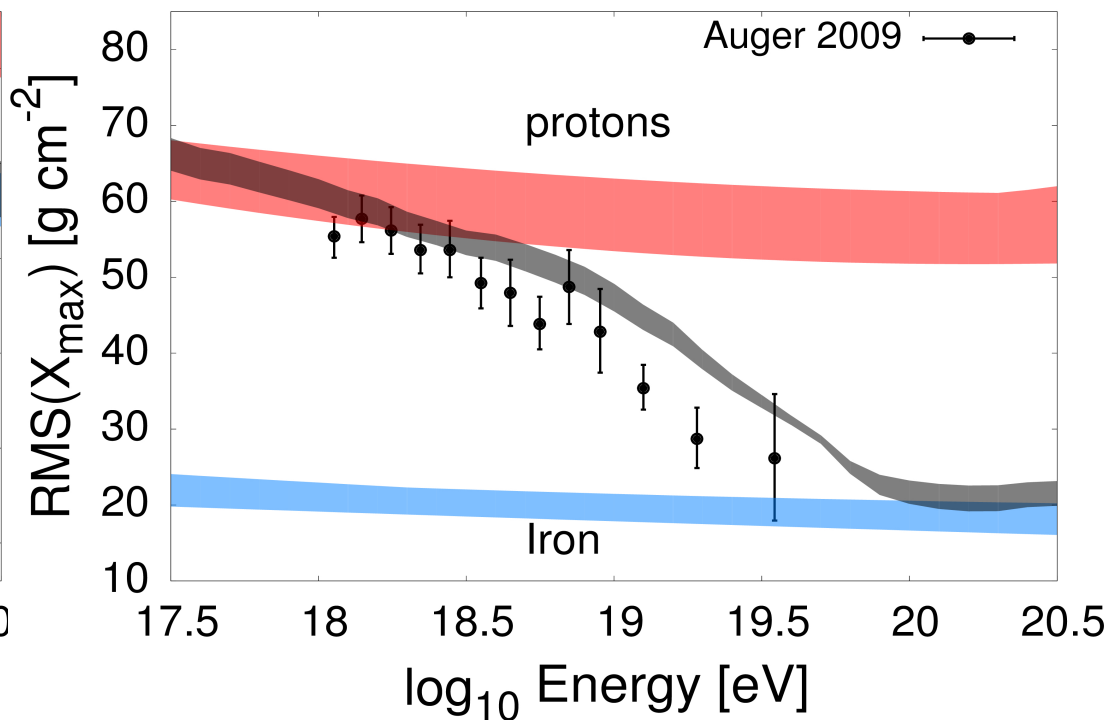
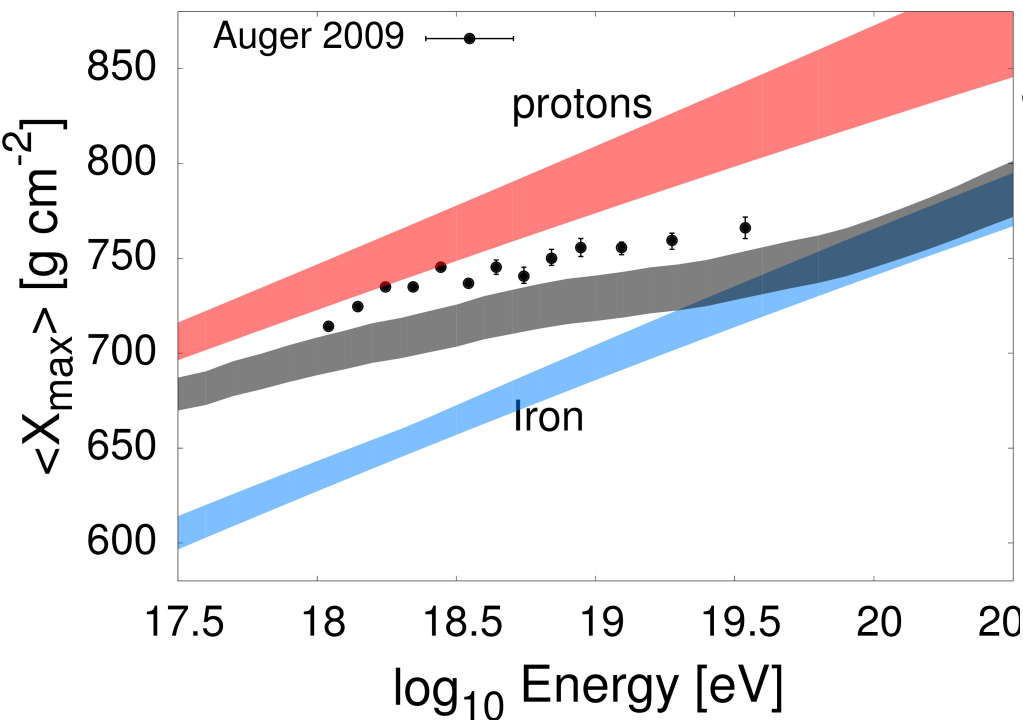
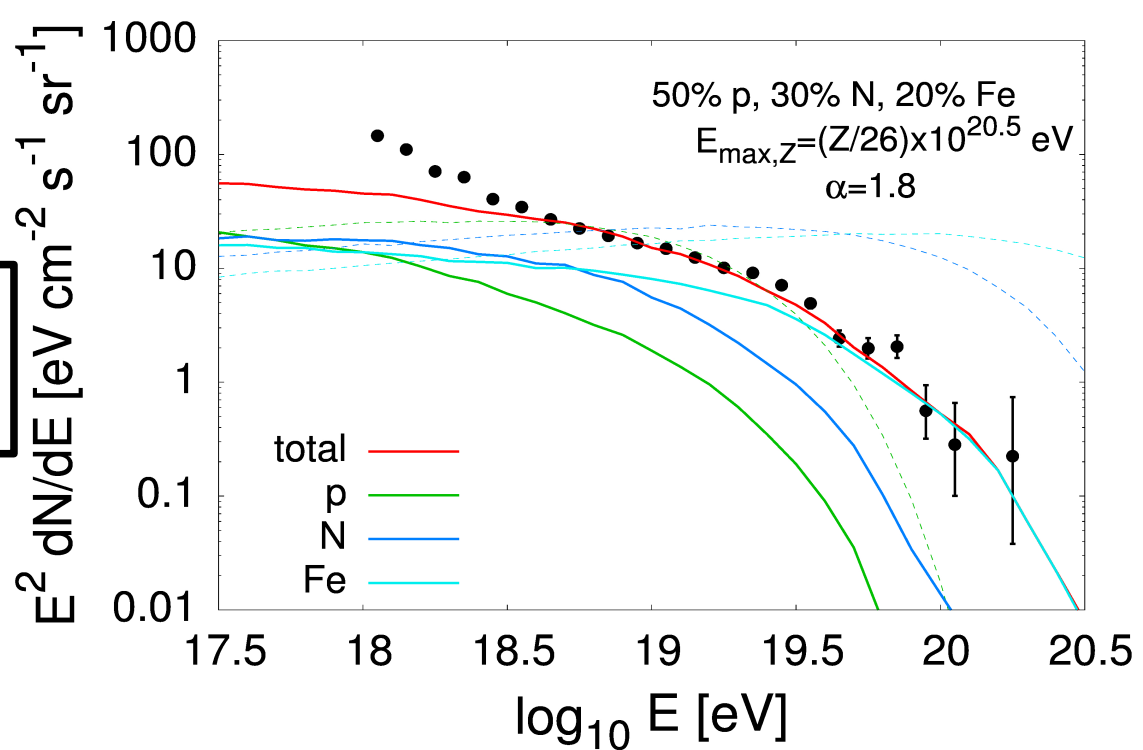
# Mixed Composition

Archetype 2: Softer  
Higher Cutoff

astro-ph/1401.0199

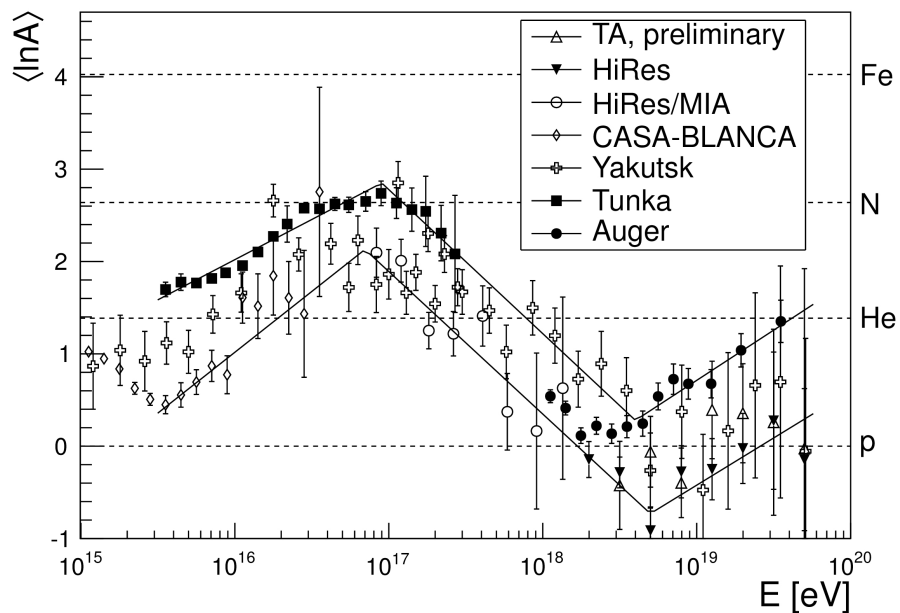
For  $E_{\text{max}} = 10^{20.5}$  eV

**50%p, 30%N, 20%Fe**

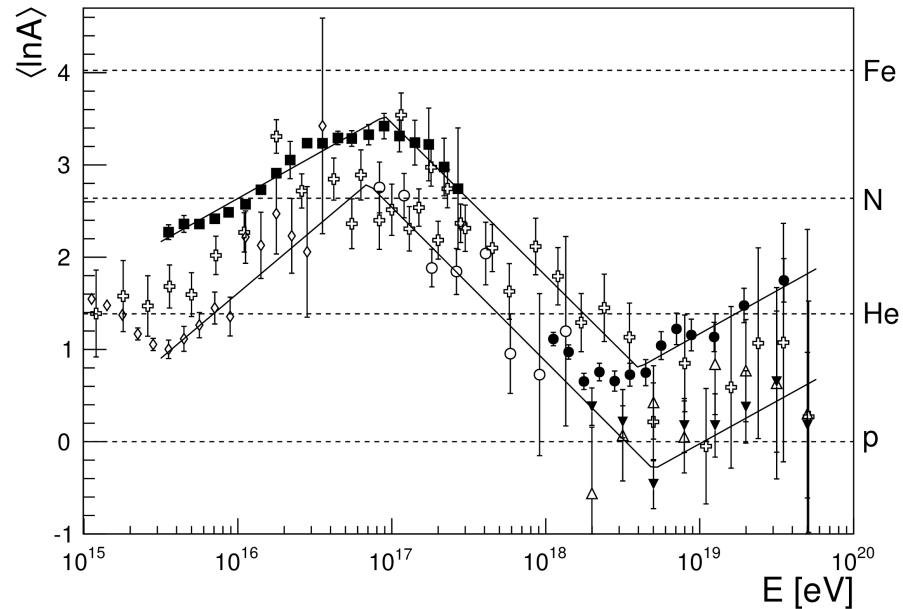


# QGSJET

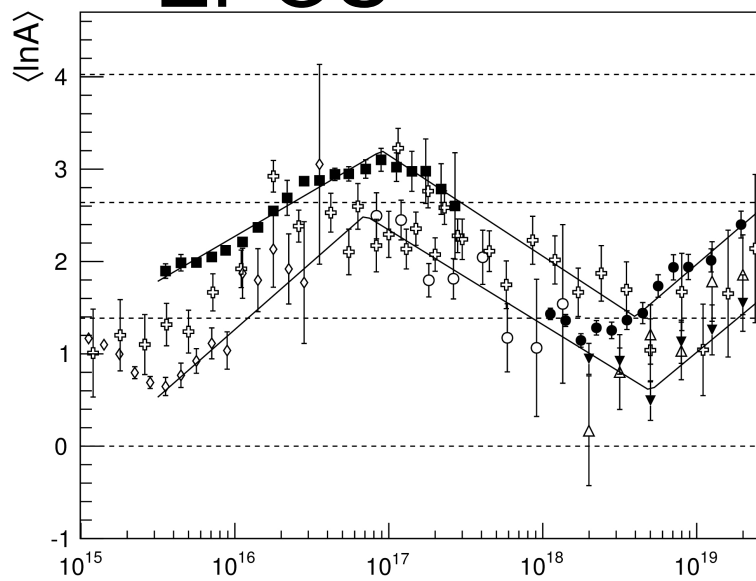
astro-ph/1201.0018



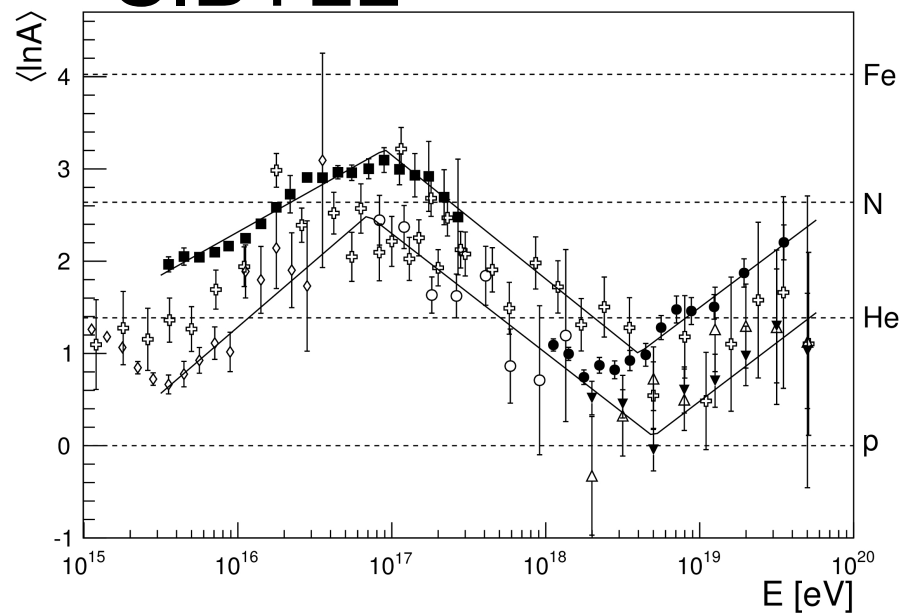
# QGSJETII



# EPOS



# SIBYLL





# Summary So Far....

The dominance of nuclei at the high energies provides useful new information about the proximity of UHECR sources

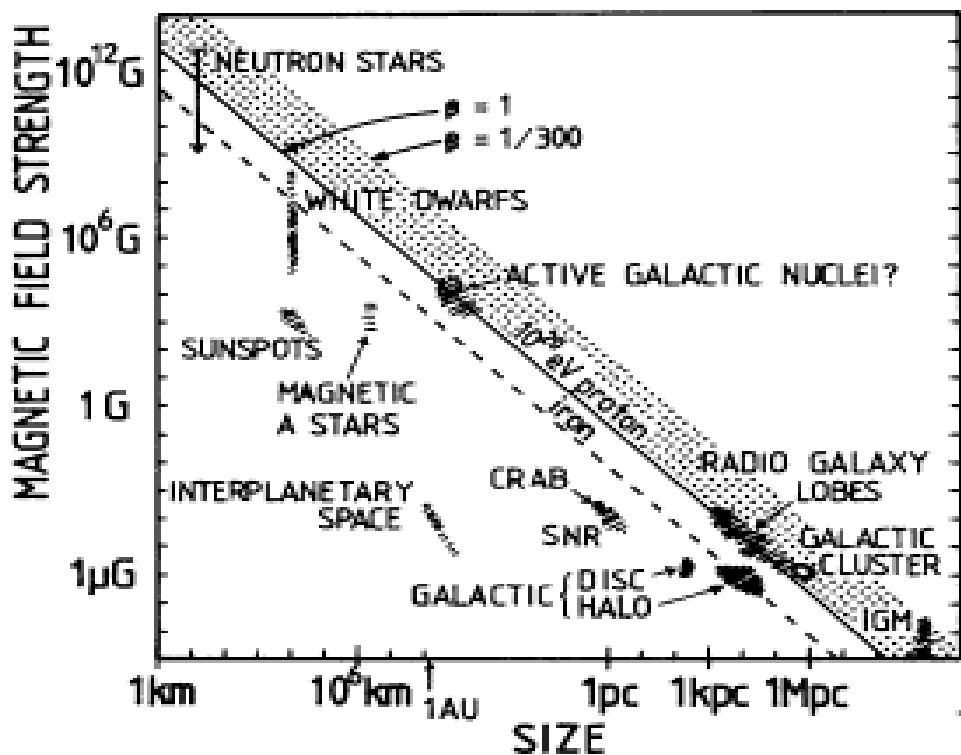
Analytic calculations can be used to describe with reasonable accuracy the spectrum and the composition results

Agreement with both the spectral and composition information require that local sources exist which produce hard spectra

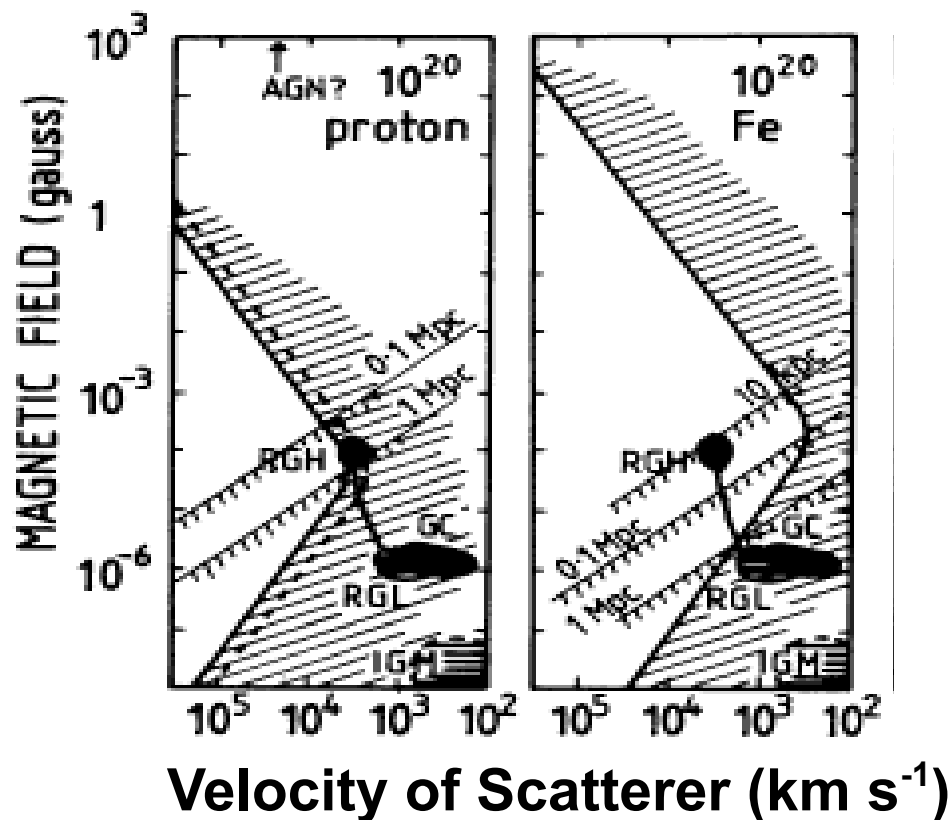
These requirements remain when the more general case of a mixed composition is considered

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# The Need For Hard Spectra Sources of Nearby Heavy Cosmic Rays



Hillas 1984



# Compactness of UHECR Sources: Proton/Nuclei Synchrotron Losses

For a  $10^{20}$  eV cosmic ray:

$$t_{\text{acc}} = \eta \frac{r_{\text{larmor}}}{c\beta^2}$$

$$t_{\text{esc}} = \frac{R^2}{c\eta r_{\text{larmor}}}$$

$$t_{\text{loss}}^{\text{sync}} = \frac{9}{8\pi Z^2 \alpha} \frac{h}{E} \left( \frac{B}{B_{\text{crit}}} \right)^2$$

$$t_{\text{loss}}^{\text{photo-dis}} = 1 \text{ Myr}$$

$$\beta Z \left( \frac{R}{\text{kpc}} \right) \left( \frac{B}{0.1 \text{ mG}} \right) > 1$$

# Compactness of UHECR Sources: Proton/Nuclei Synchrotron Losses

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$$t_{\text{loss}}^{\text{photo-dis}} = 1 \text{ Myr}$$

$$\left( \frac{\beta}{Z} \right)^2 \left( \frac{30 \text{ mG}}{B} \right) > 1$$

# Compactness of UHECR Sources: Proton/Nuclei Synchrotron Losses

For a  $10^{20}$  eV cosmic ray:

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$$t_{\text{loss}}^{\text{photo-dis}} = 1 \text{ Myr}$$

$$Z\beta^2 \left( \frac{B}{0.1 \text{ mG}} \right) > 1$$

# Compactness of UHECR Sources: Proton/Nuclei Synchrotron Losses

-Need for high velocity shocks  
( $\beta \sim 1$ )

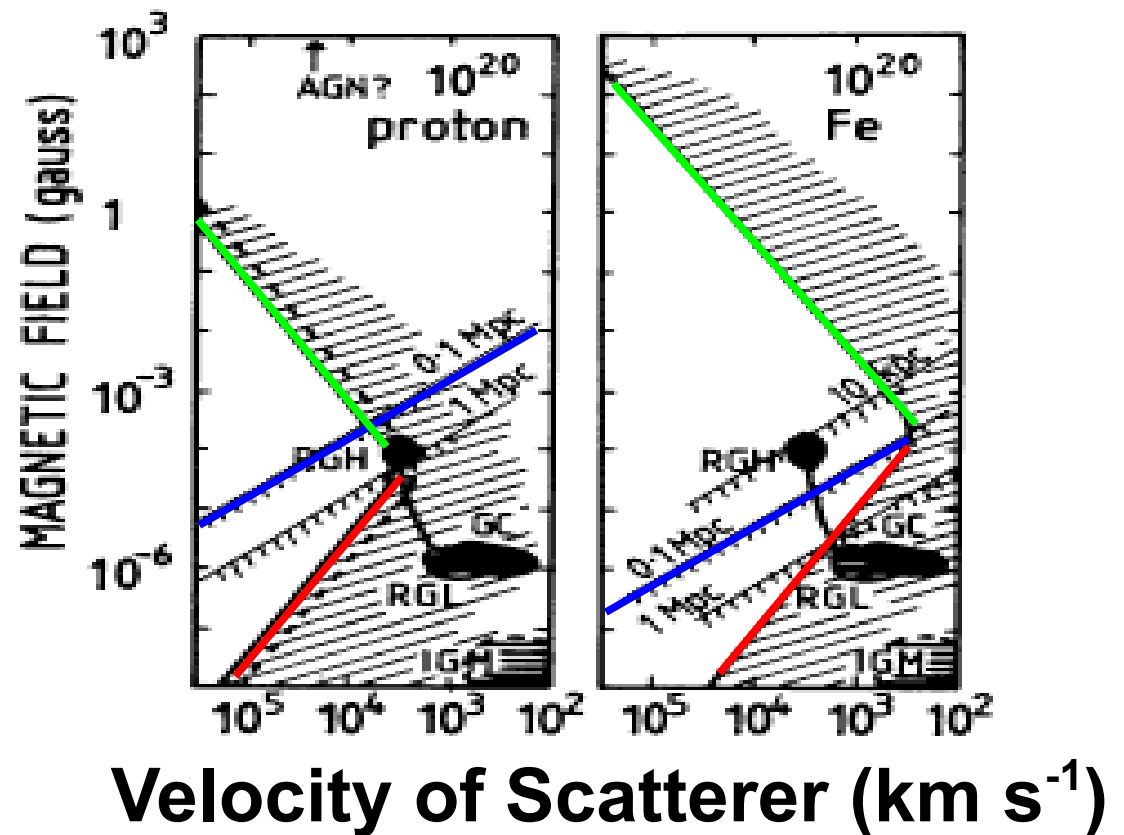
-Need for close to optimal  
acceleration conditions  
( $\eta \sim 1$  - Bohm Limit)

$$\beta Z \left( \frac{R}{\text{kpc}} \right) \left( \frac{B}{0.1 \text{ mG}} \right) > 1$$

$$\left( \frac{\beta}{Z} \right)^2 \left( \frac{30 \text{ mG}}{B} \right) > 1$$

$$Z \beta^2 \left( \frac{B}{0.1 \text{ mG}} \right) > 1$$

$\eta \sim 10$  (solid lines)



# Compactness of UHECR Sources: Nuclei Photo-disintegration

$$f = \frac{t_{\text{esc}}}{t_{\text{int.}}^{\text{CR}\gamma}}$$

$$t_{\text{int.}}^{\text{CR}\gamma} \approx \frac{1}{n_{\gamma} \sigma_{\text{CR}\gamma} c}$$

$$n_{\gamma} = \frac{L_{\gamma}}{c 4\pi R^2 \epsilon_{\gamma}}$$

$$t_{\text{esc}} \approx \frac{R^2}{2D} = \frac{3R^2}{2R_{\text{Larmor}}}$$

$$f^{\text{CR}\gamma} = \frac{3L_{\gamma} \sigma_{\text{CR}\gamma} Z e B}{8\pi \epsilon_{\gamma} E_{\text{CR}}}$$



# Compactness of UHECR Sources: Nuclei Photo-disintegration

$$f^{\text{CR}\gamma} = \frac{3L_\gamma \sigma_{\text{CR}\gamma} Z e B}{8\pi \epsilon_\gamma E_{\text{CR}}} = \frac{s_1}{s_2}$$

Photo-disintegration threshold:

$$2E_{\text{CR}} \epsilon_\gamma > A m_p c^2 E_{\text{bind.}}, \text{ where } m_p c^2 E_{\text{bind.}} = 10^{16} \text{ eV}^2$$

Since,

$$L_\gamma [10^{43} \text{ erg s}^{-1}] = 2 \times 10^{44} \text{ eV cm}^{-1}$$
$$\sigma_{\text{CR}\gamma} [A \text{ mb}] = A \times 10^{-27} \text{ cm}^2$$
$$eB [1 \text{ mG}] = 0.3 \text{ eV cm}^{-1}$$

$$\frac{L_\gamma \sigma_{\text{CR}\gamma} e B}{A} = 6 \times 10^{16} \text{ eV}^2$$

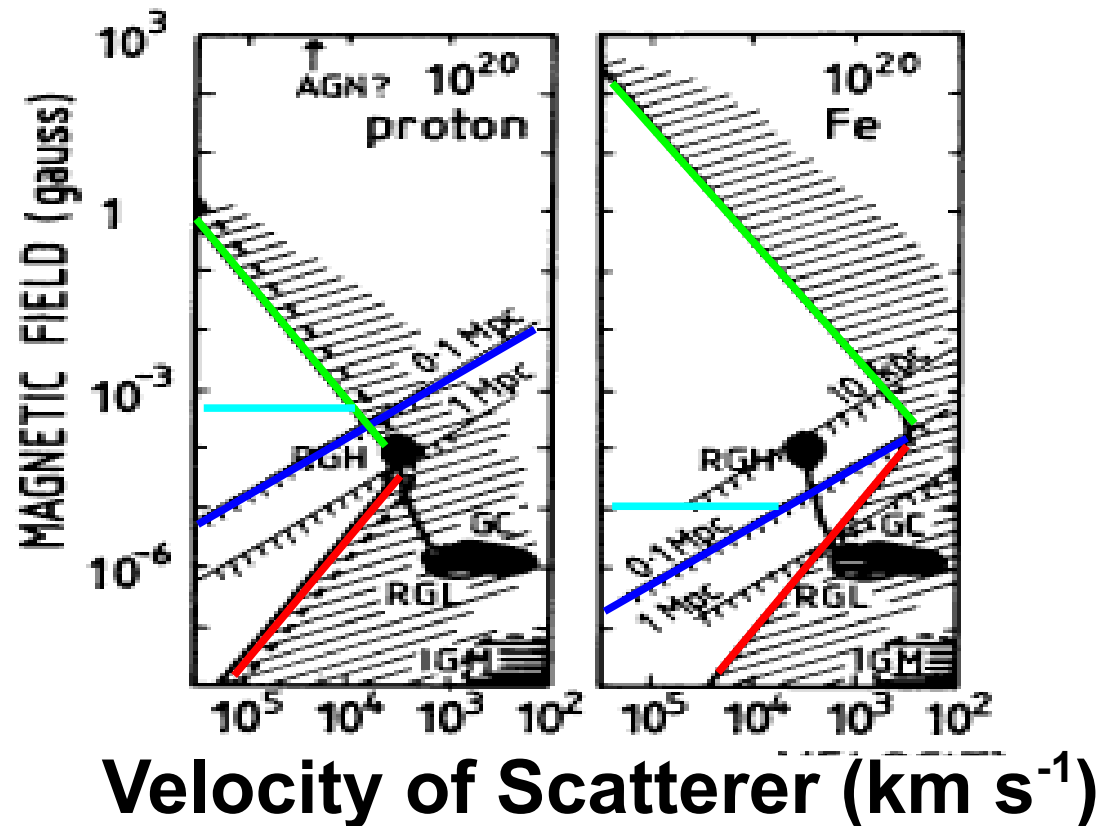




# Compactness of UHECR Sources: Nuclei Photo-disintegration

Summary-

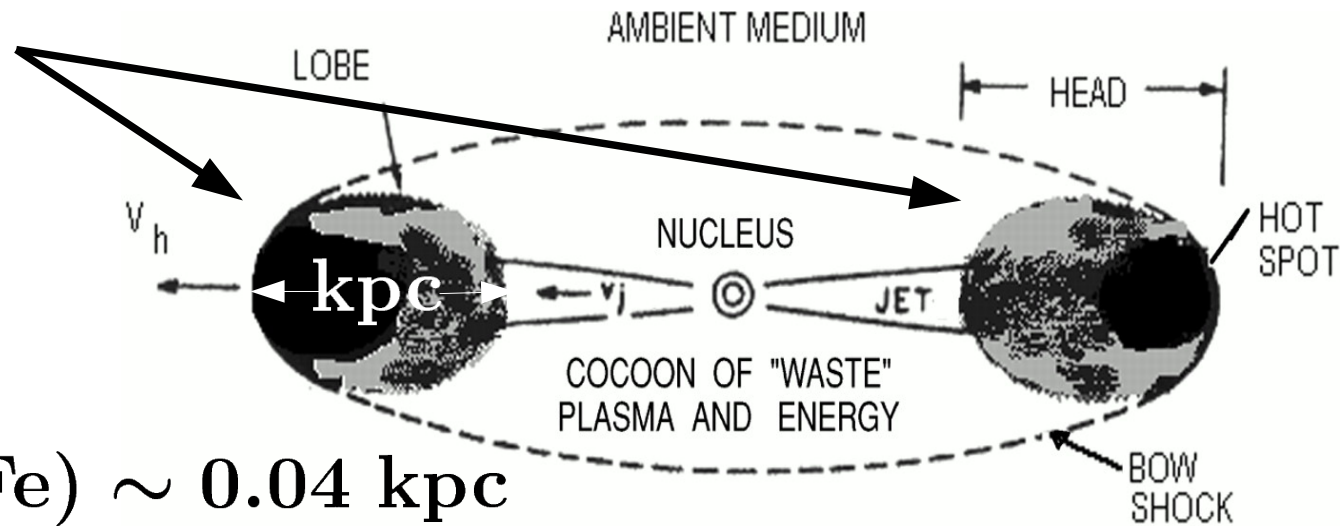
$$f^{\text{CR}\gamma} = 50 \left( \frac{Z}{26} \right) \left( \frac{L_\gamma}{10^{43} \text{ erg s}^{-1}} \right) \left( \frac{B}{1 \text{ mG}} \right)$$



# Example Candidate UHECR Source (a Nuclei Friendly Environment)

Stochastic Acceleration  
in Radio Lobes:

Diagram taken from Ferrari -1998



$$B_{\text{source}} \sim 10^{-4} \text{ G}$$

$$\rightarrow R_{\text{Larmor}}(10^{20} \text{ eV Fe}) \sim 0.04 \text{ kpc}$$

$$t_{\text{acc}} < 10^6 \text{ yrs for } \beta_{\text{scat.}} > 10^{-2}$$

(INTERSTELLAR OR INTERGALACTIC GAS)

**General PROBLEM for Non-Compact  
Accelerators- ACCELERATION TIME**



# Example Candidate UHECR Source (a Nuclei Friendly Environment)

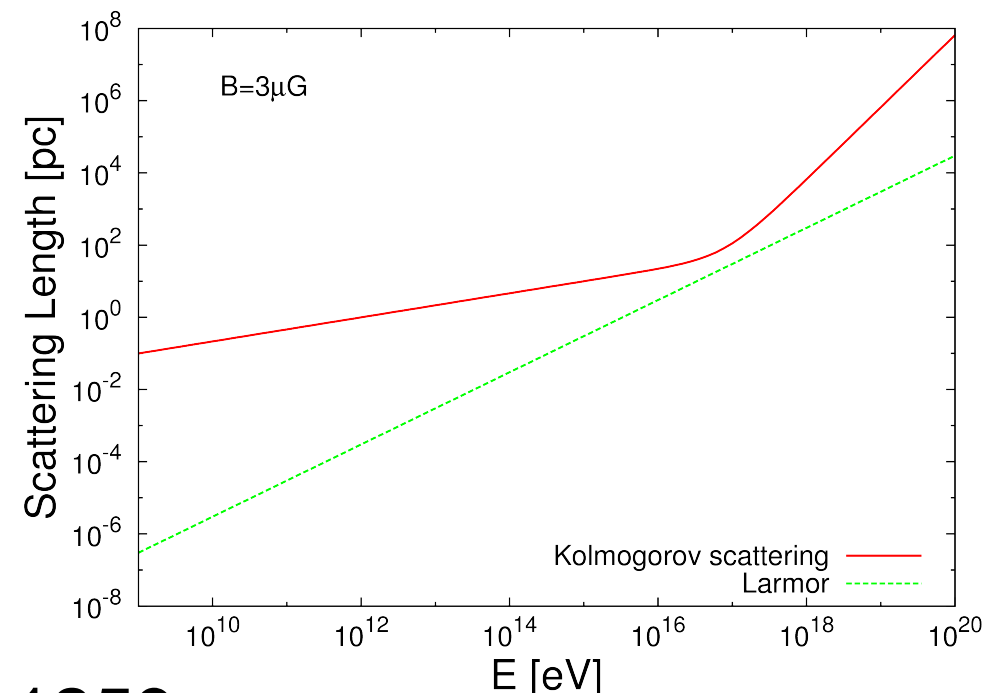
$$t_{\text{acc}} = \eta \frac{r_{\text{larmor}}}{c\beta^2}$$

$\eta$  dependent on normalisation and slope of magnetic turbulence power-spectrum

$$P(k) \propto k^{-\alpha} \quad \alpha = \frac{5}{3}$$

Important parameters:

$$\frac{\delta B^2}{B_0^2}, \quad \lambda_{\text{max}}$$



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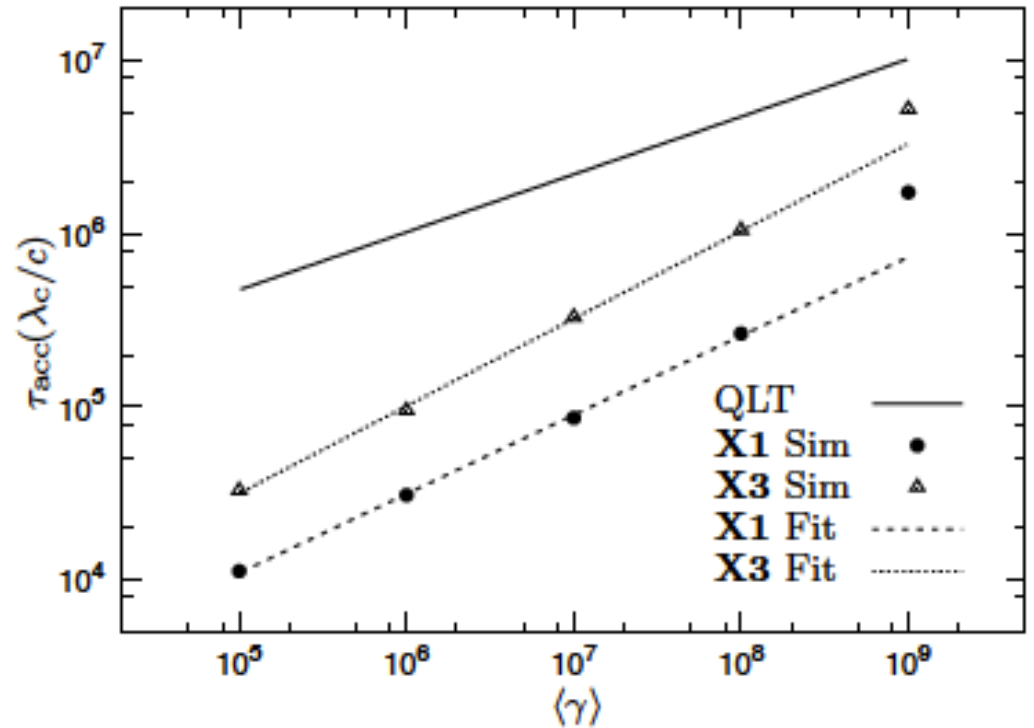
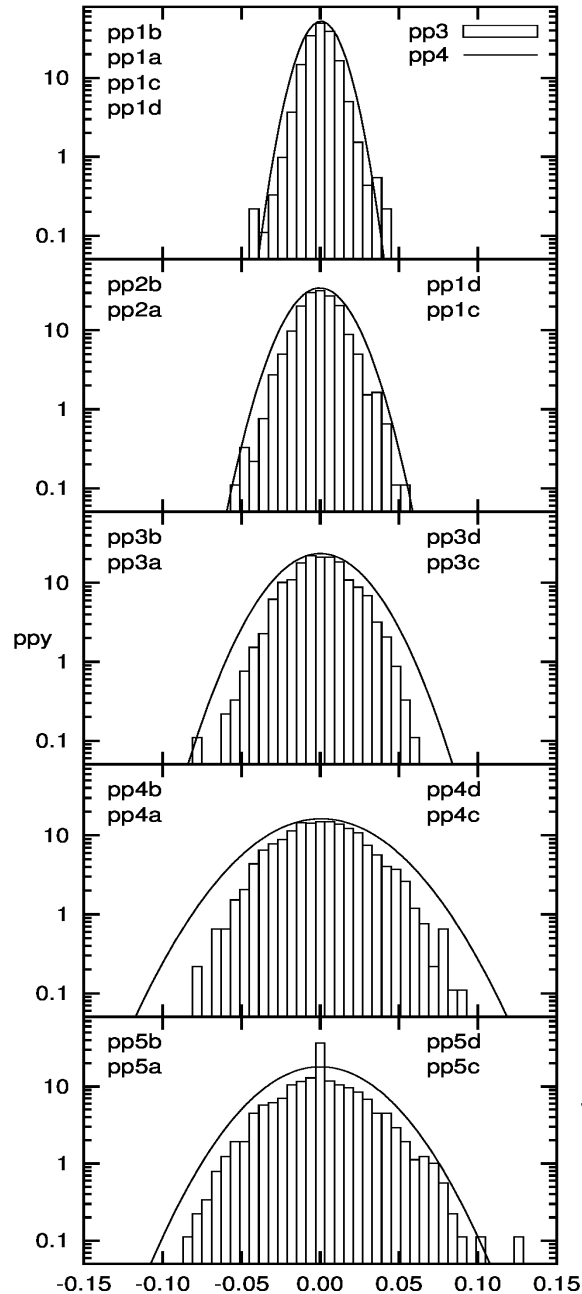
astro-ph/0903.1259

# Example Candidate UHECR Source

(a Nuclei Friendly Environment)

astro-ph/0903.1259

$$t_{\text{acc}} = \eta \frac{r_{\text{larmor}}}{c\beta^2}$$



Yes, but requires:

$$\beta_A > 0.1$$

where 
$$\beta_A = \frac{1}{c} \frac{B}{\sqrt{4\pi m_p n_p}}$$

# Which Nearby AGN?

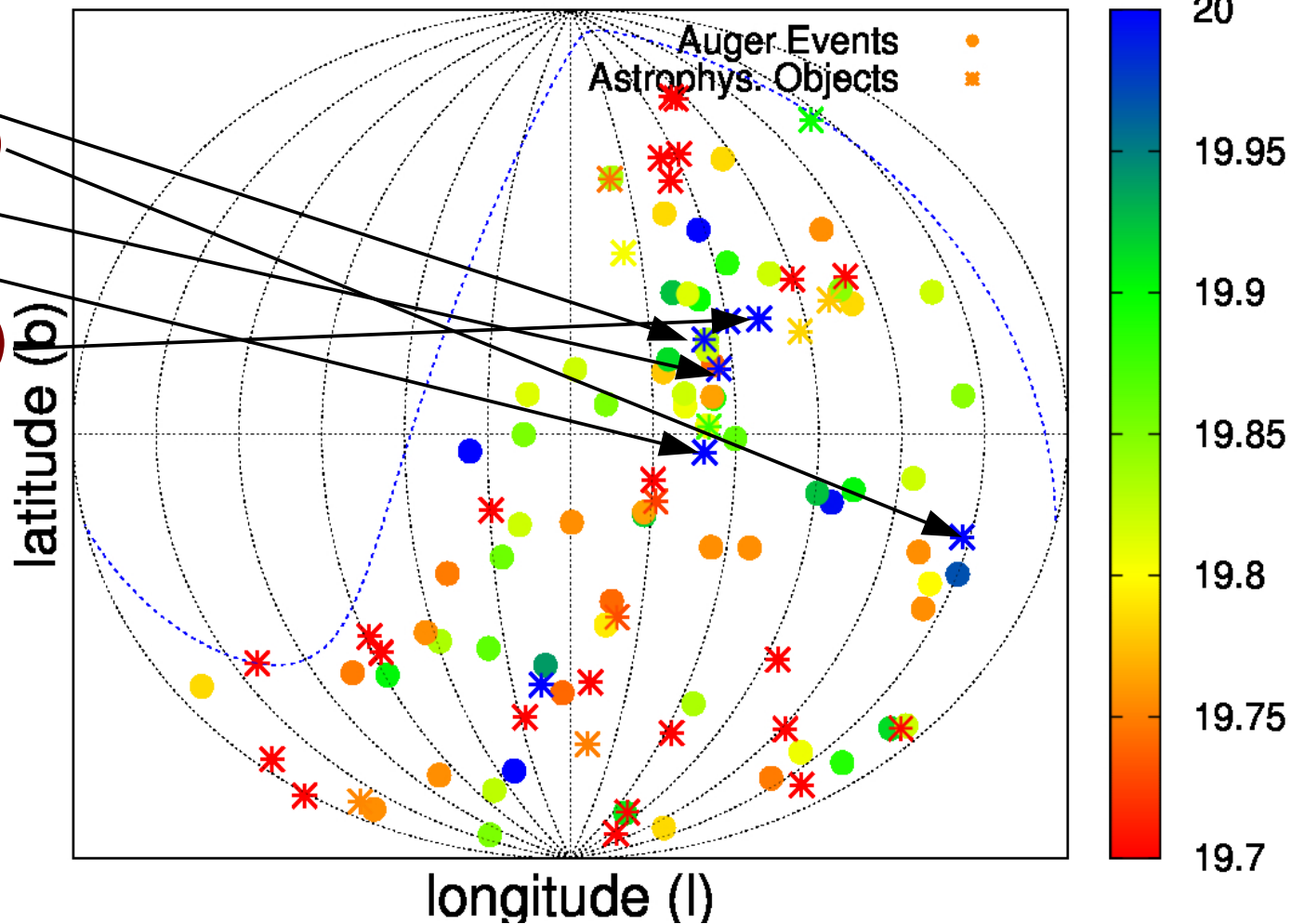
$(L_X > 10^{42} \text{ erg s}^{-1}, D < 60 \text{ Mpc})$

## TOP 5 in $L/D^2$ :

- 1) Cen A (4 Mpc)
- 2) NGC 2110 (30 Mpc)
- 3) NGC 4945 (4 Mpc)
- 4) Circinus (4 Mpc)
- 5) NGC 3783 (40 Mpc)

(from 1<sup>st</sup> Integral or Swift 58 month catalogues)

**TOTAL: 41**



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$$s = \log_{10}(L_X / z^2) - 27.5$$



# Which Nearby AGN?

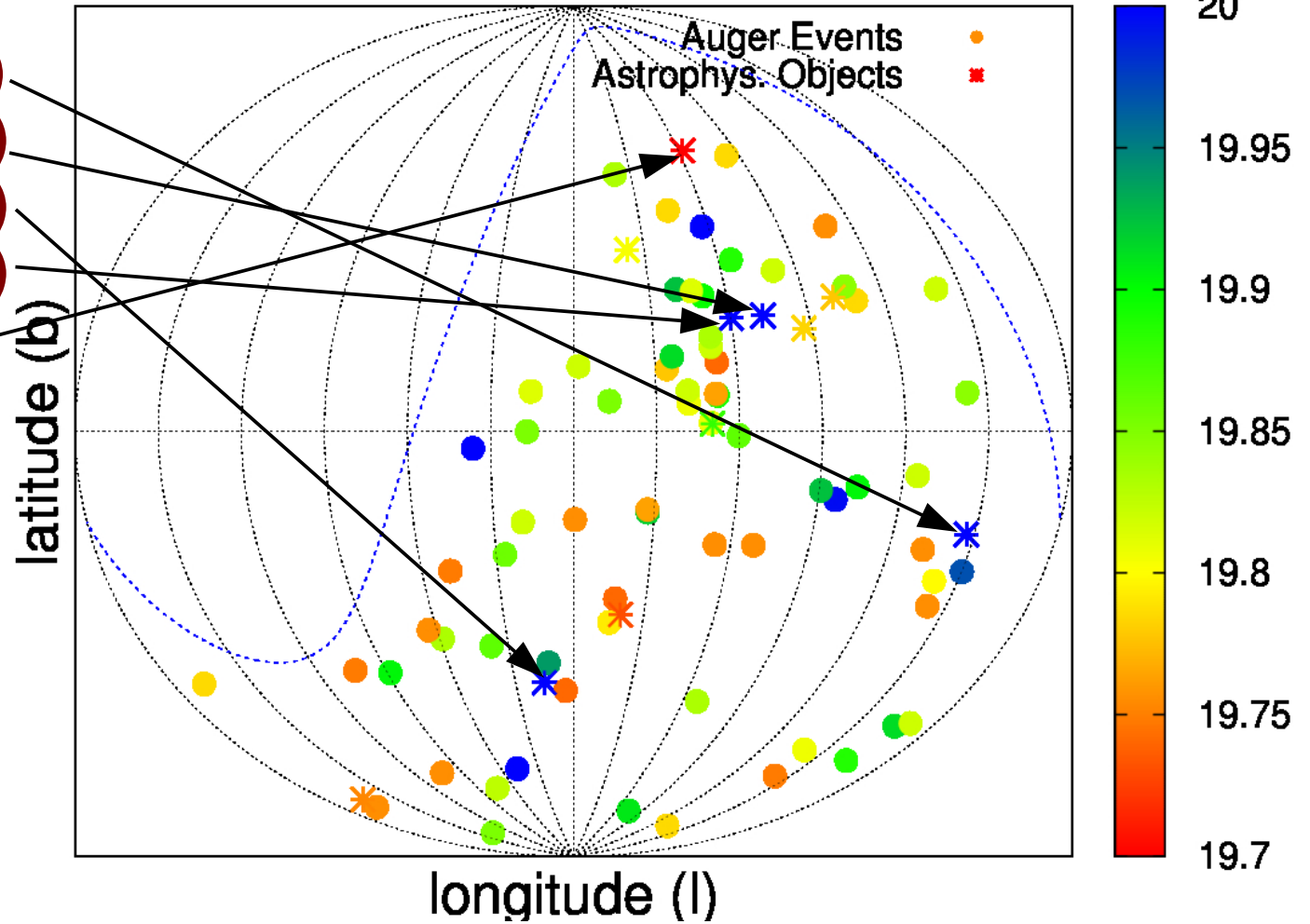
$(L_X > 10^{43} \text{ erg s}^{-1}, D < 60 \text{ Mpc})$

## TOP 5 in $L/D^2$ :

- 1) NGC 2110 (30 Mpc)
- 2) NGC 3783 (40 Mpc)
- 3) NGC 7172 (38 Mpc)
- 4) NGC 4507 (52 Mpc)
- 5) 4U 1344 (53 Mpc)

(from 1<sup>st</sup> Integral or Swift 58 month catalogues)

**TOTAL: 11**



$$s = \log_{10}(L_X / z^2) - 27.5$$

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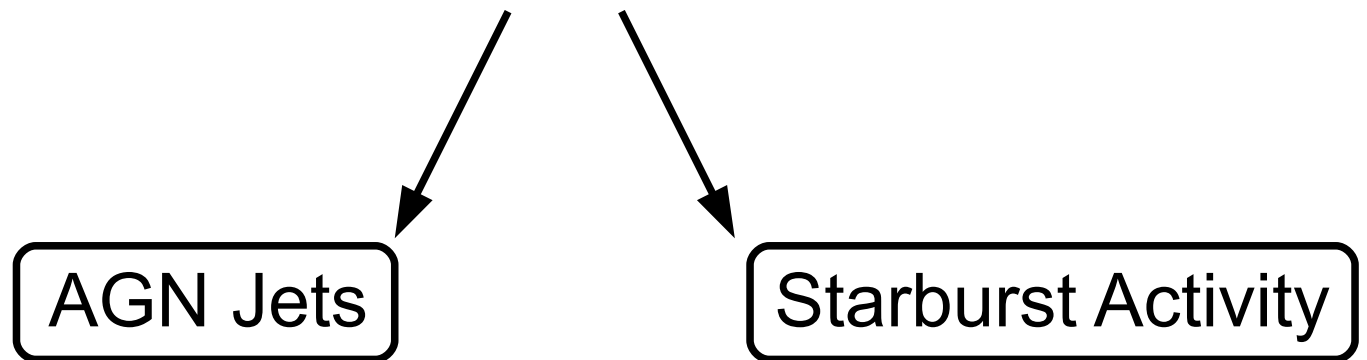


# VHE gamma-ray Emission Local AGN?

( $L_x > 10^{42}$  erg s $^{-1}$ ,  $D < 60$  Mpc)

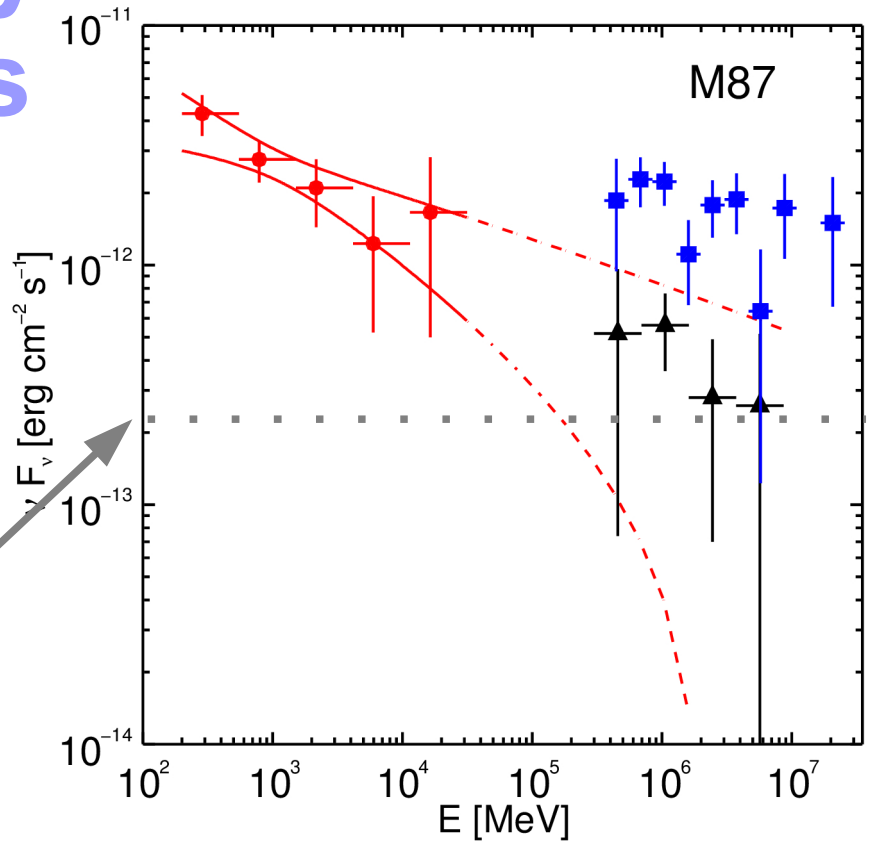
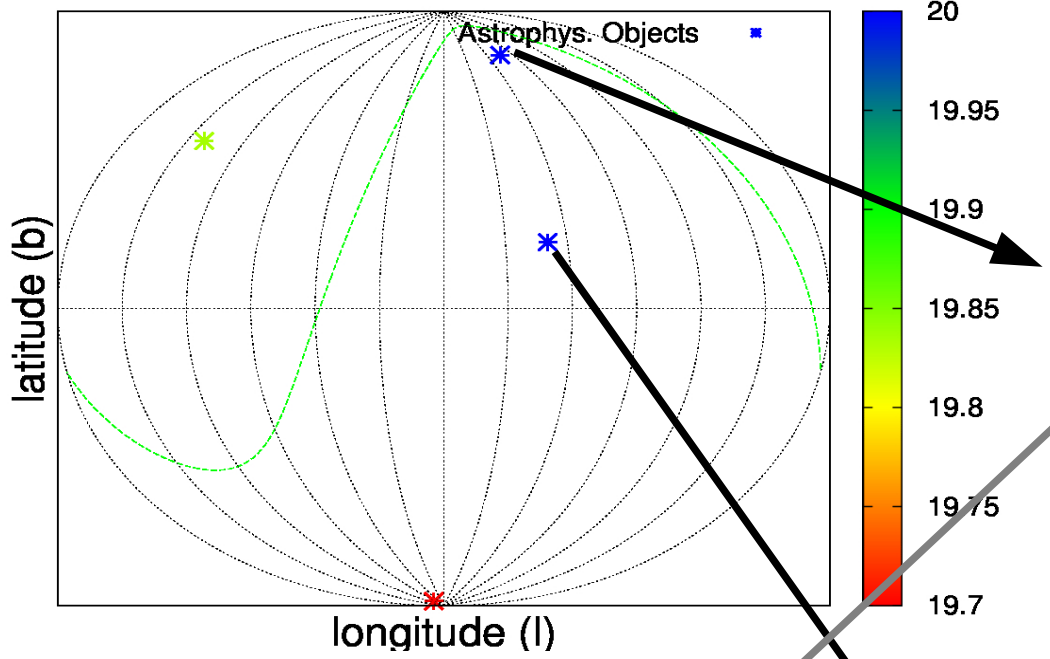
<b>TOP 5 in <math>L/D^2</math>:</b>	GeV Emission	TeV Emission
<b>1) Cen A (4 Mpc)</b>	✓	✓
<b>2) NGC 2110 (30 Mpc)</b>		
<b>3) NGC 4945 (4 Mpc)</b>	✓	
<b>4) Circinus (4 Mpc)</b>	✓	
<b>5) NGC 3783 (40 Mpc)</b>		

(from 1<sup>st</sup> Integral or  
Swift 58 month  
catalogues)

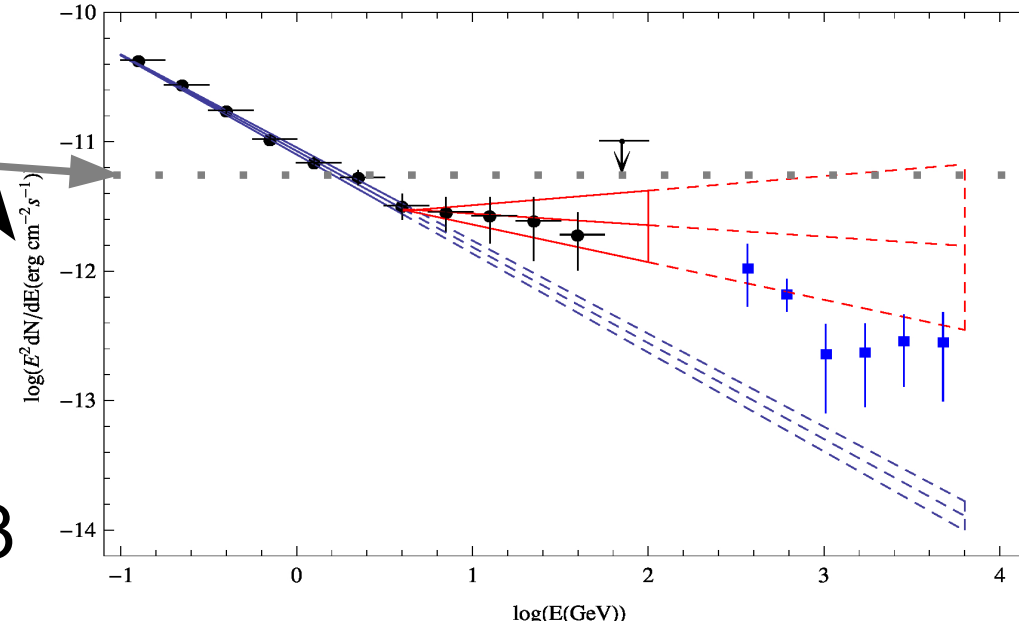


# Local AGN Detected by Cherenkov Telescopes

Abdo et al. 2009



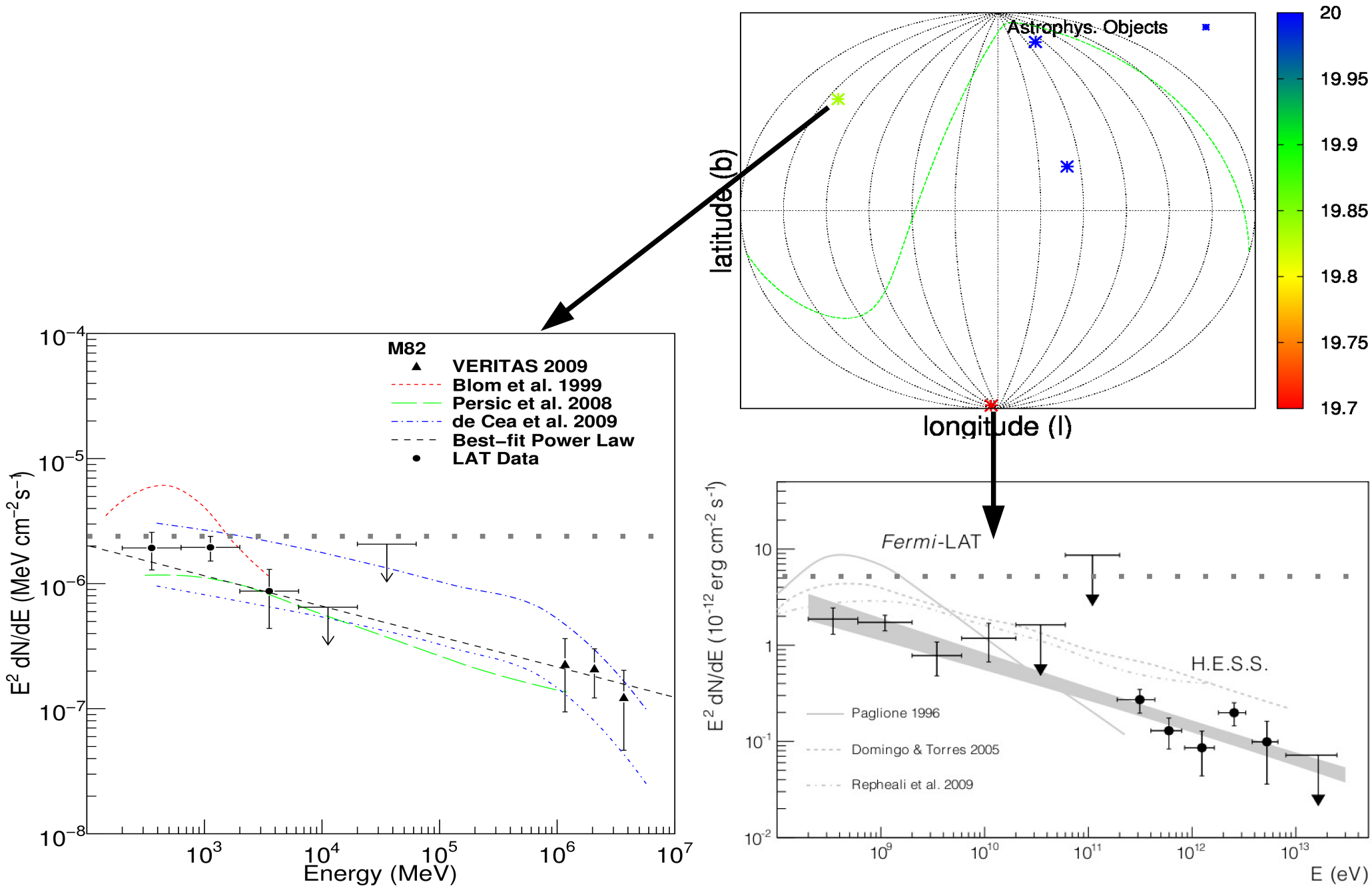
$10^{40} \text{ erg s}^{-1}$



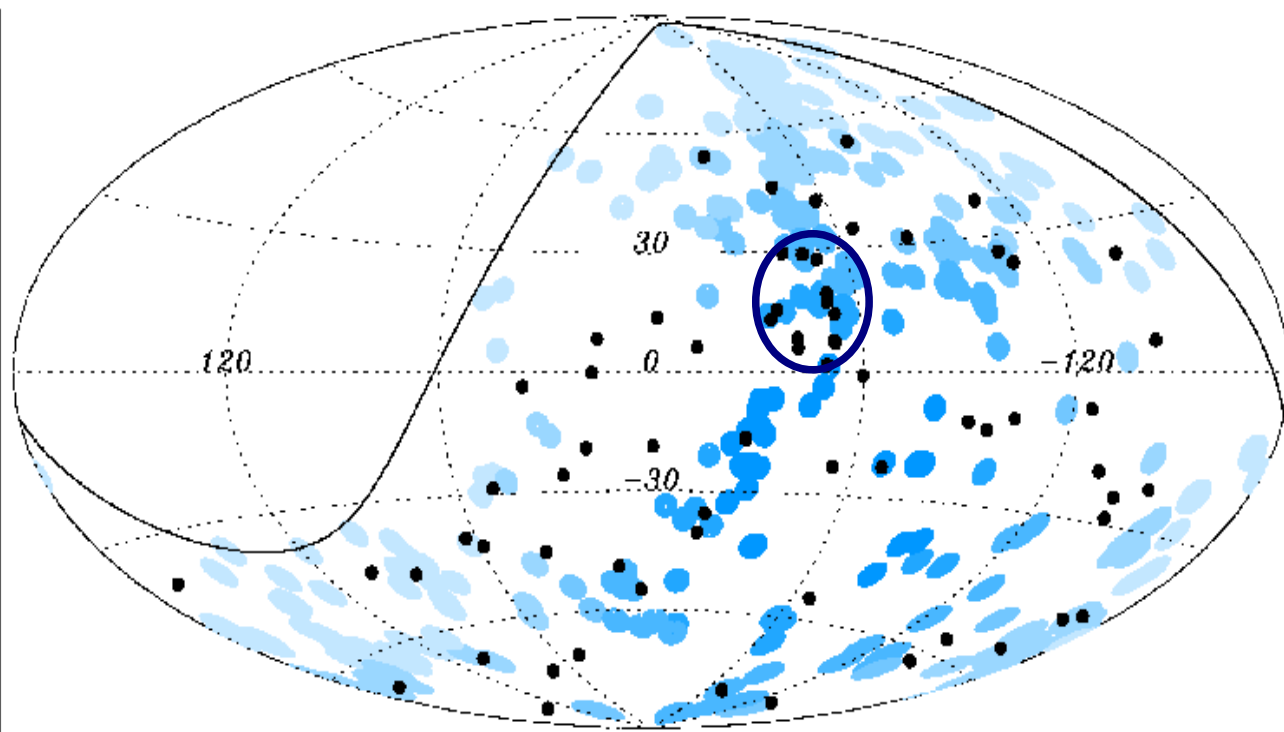
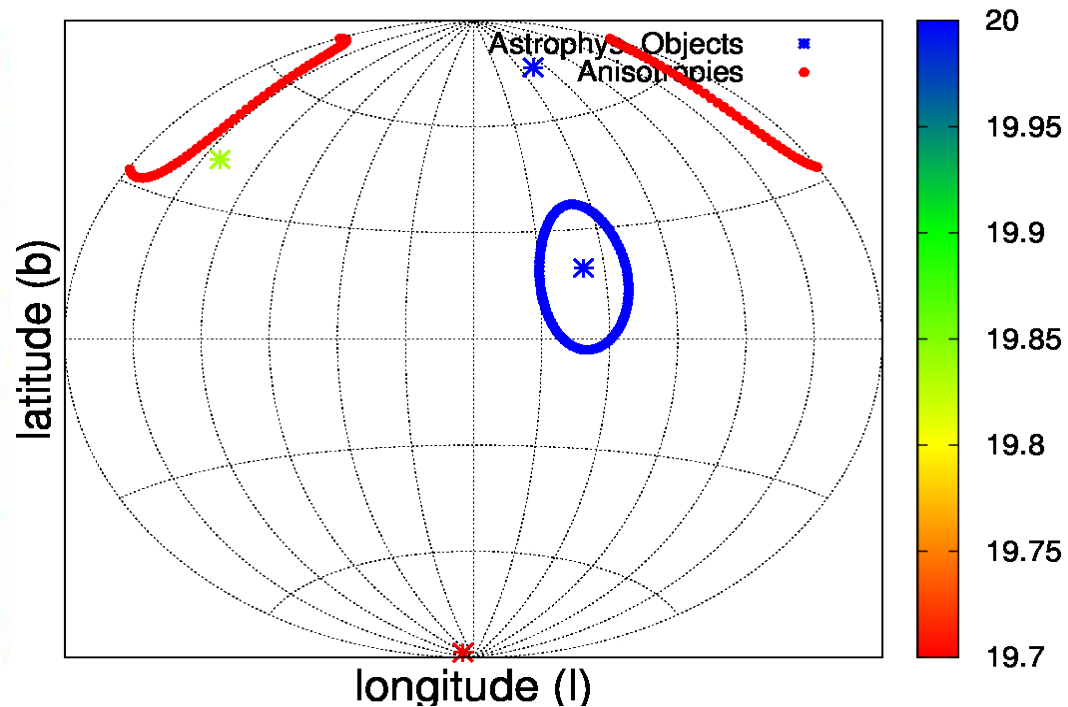
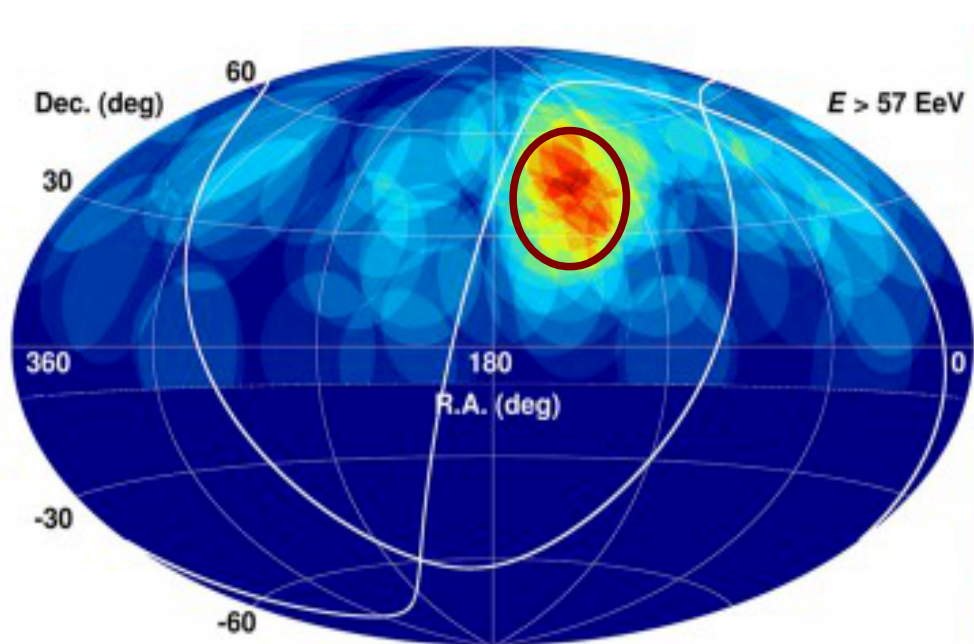
Sahakyan et al. 2013



# Local AGN/Starburst Galaxies Detected by Cherenkov Telescopes

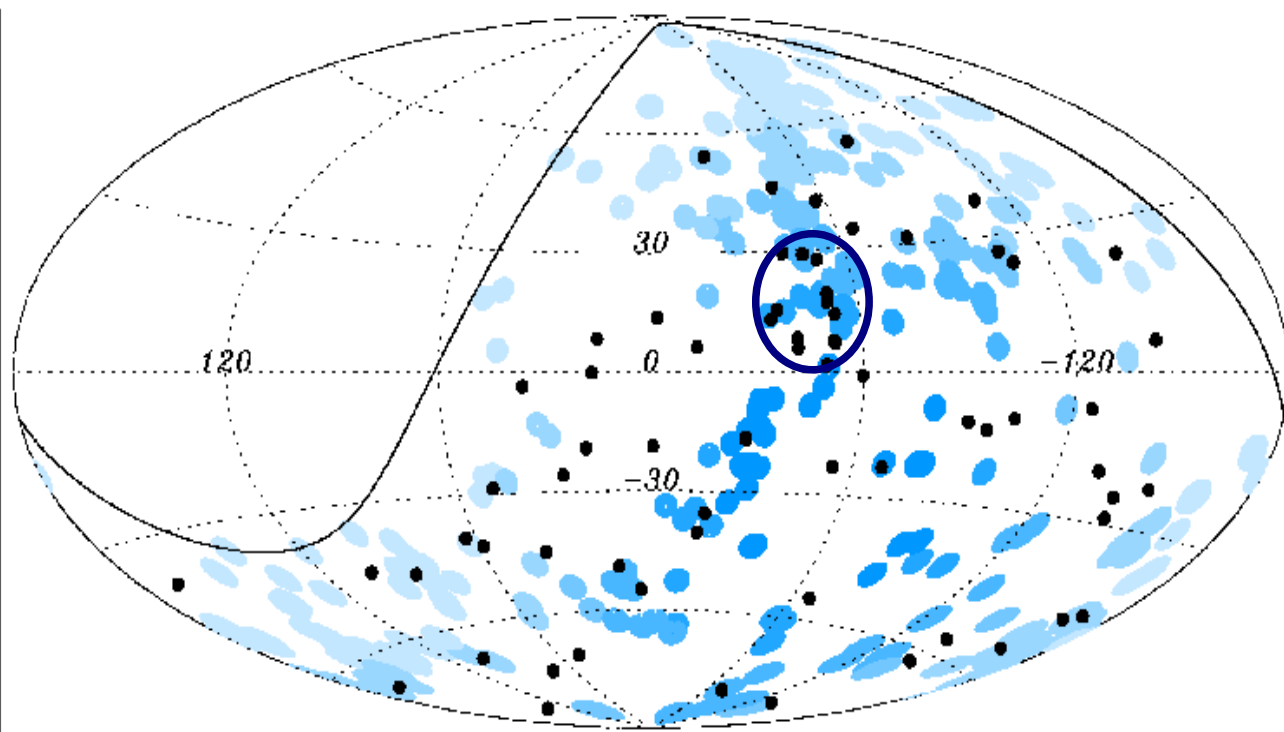
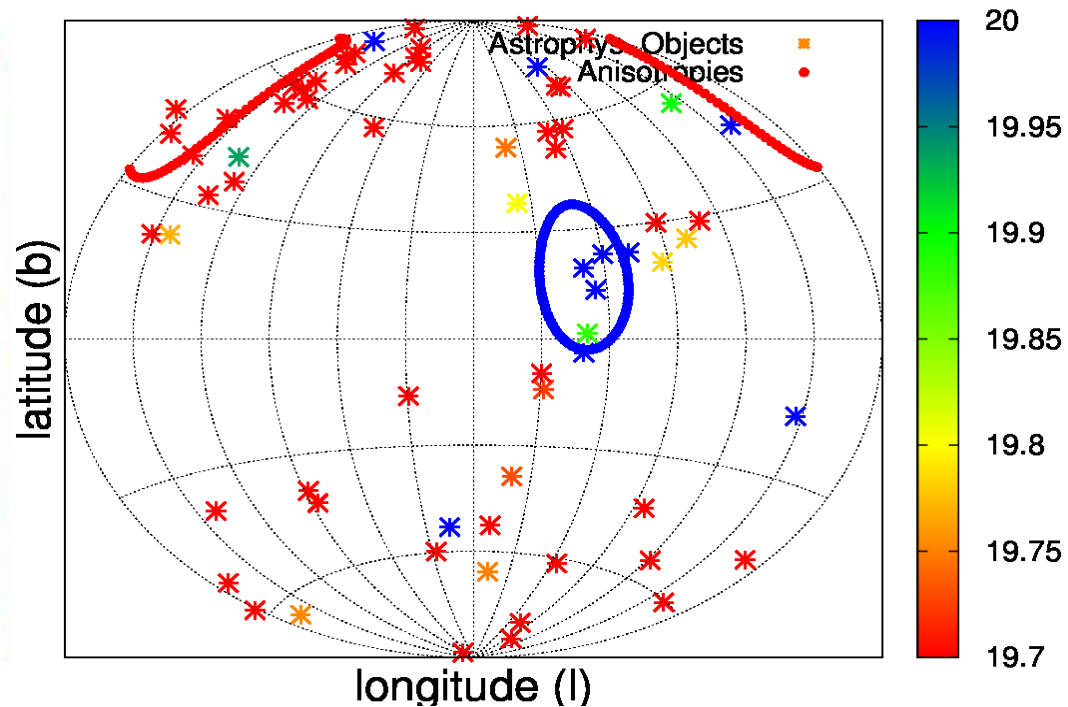
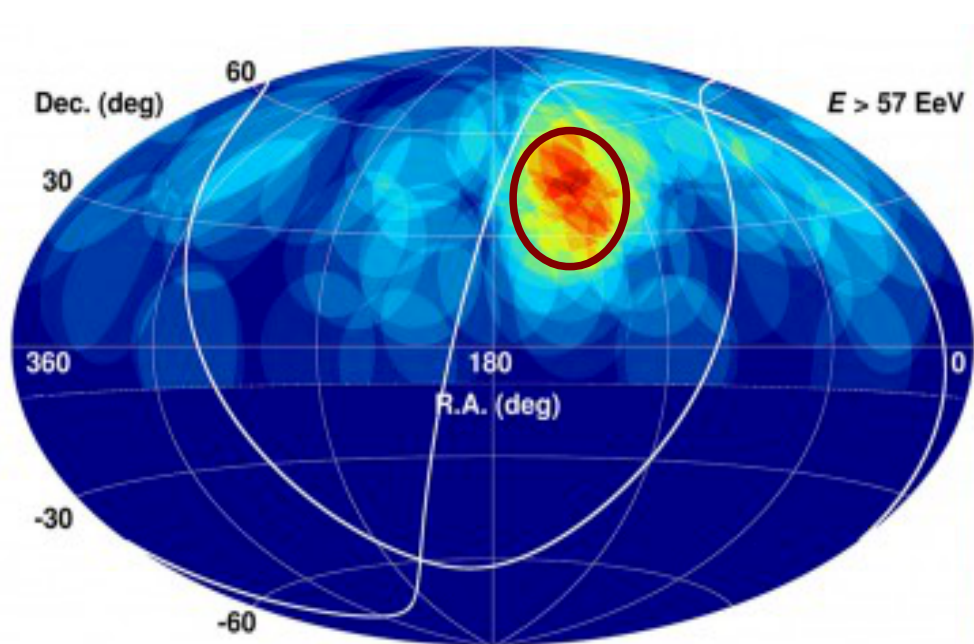


# Could There Be A Connection.....?



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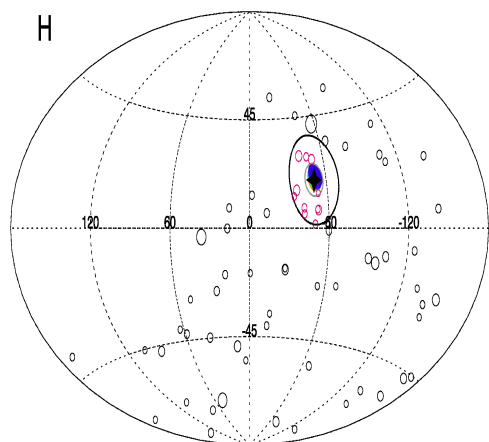
# Could There Be A Connection.....?



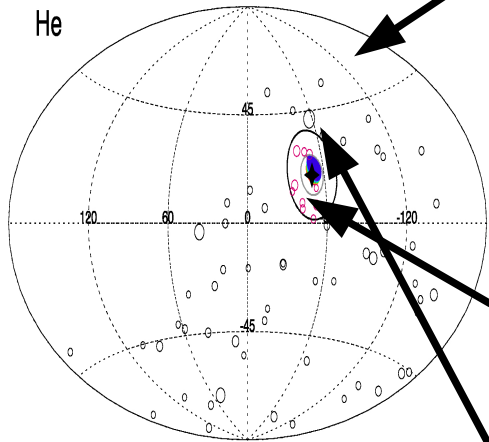
Andrew Taylor

# Anisotropy Signatures of UHECR From Cen A?

p only



He only

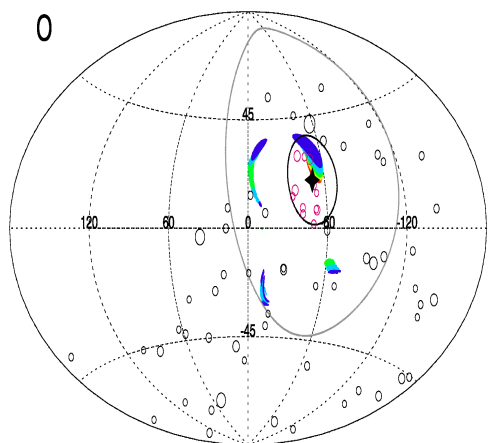


18°

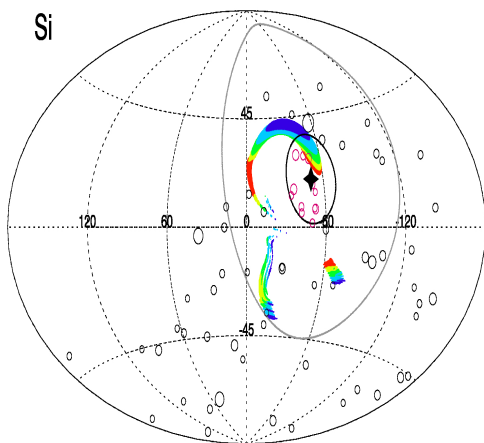
13 events

Cen A

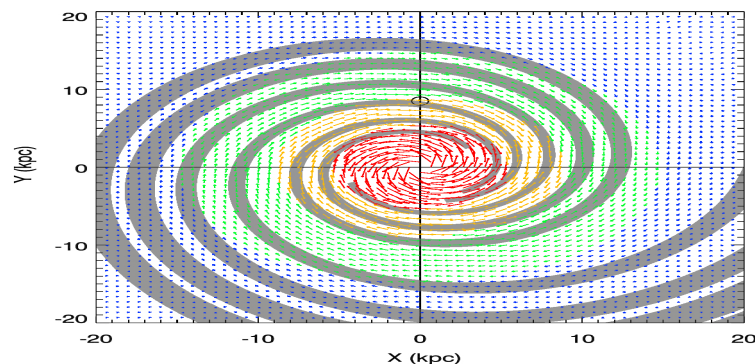
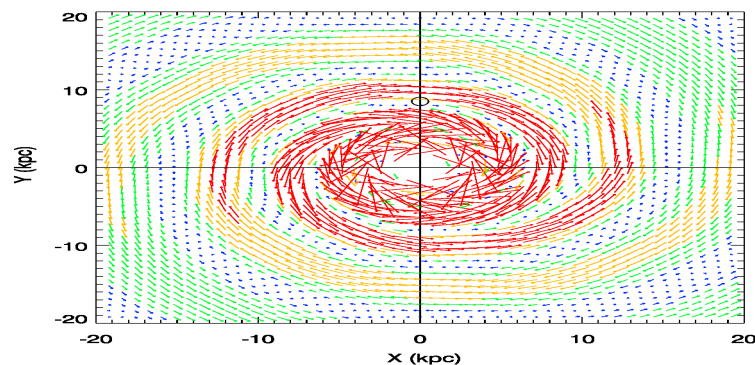
O only



Si only



astro-ph/1206.3907

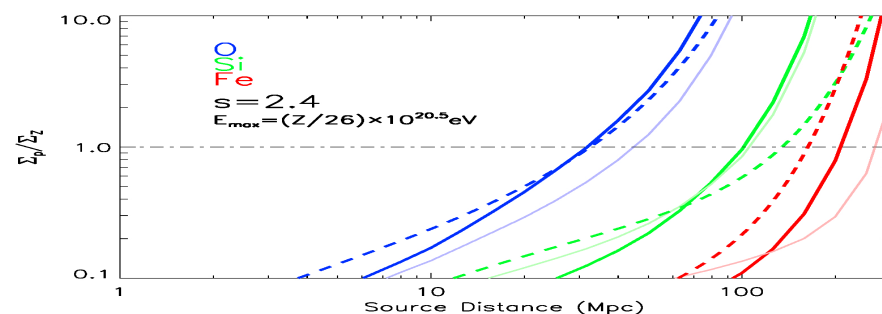
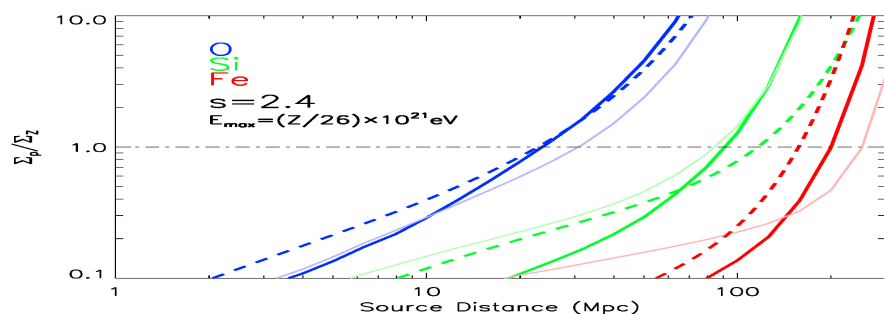
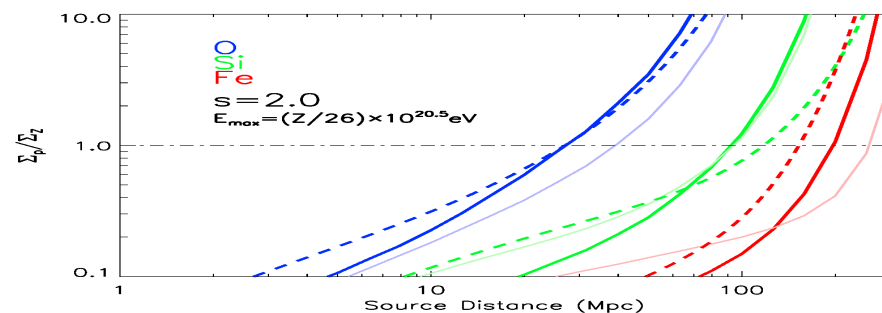
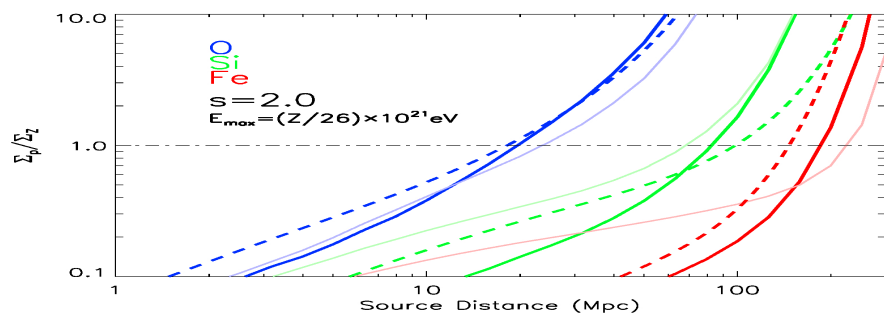




# Anisotropy Signatures of UHECR Sources?

$$55 \text{ EeV} < E_Z < 84 \text{ EeV}$$

$$55/Z \text{ EeV} < E_p < 84/Z \text{ EeV}$$



# Conclusion

The dominance of nuclei at high energies provides useful new information about the proximity and spectra of UHECR sources

Intermediate (not compact or too large scale) regions of AGN potentially can satisfy the source requirement criterion

The correlation of UHECR with AGN coupled with their composition information has great potential for providing “depth” information to a potential source

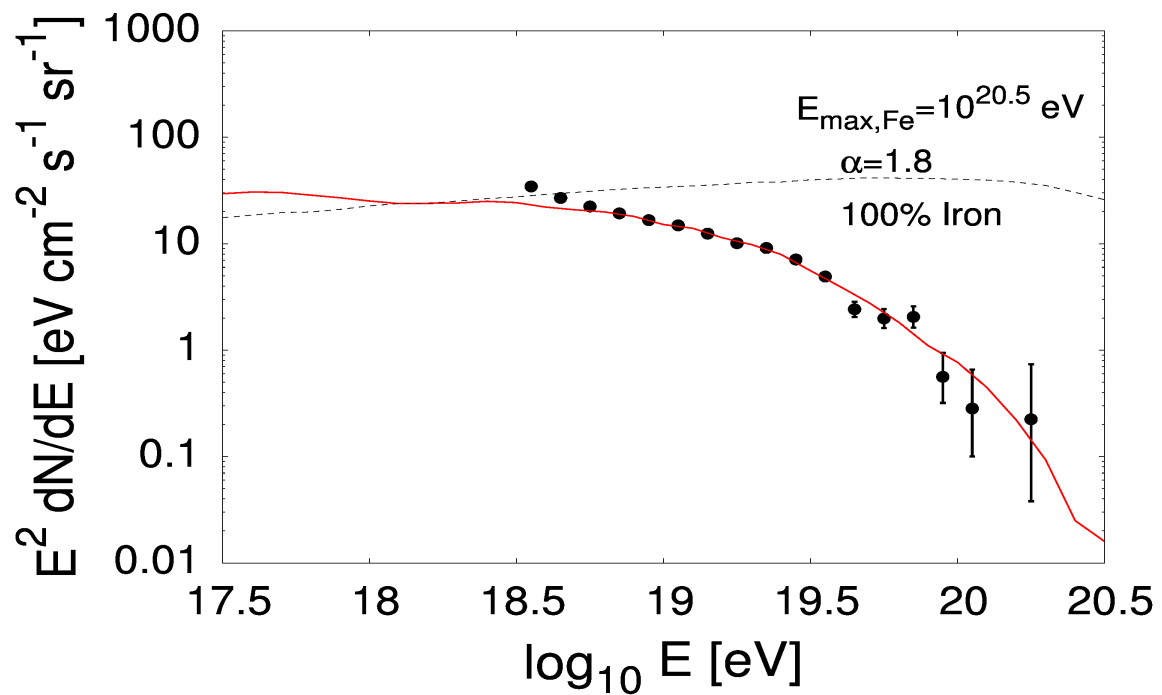
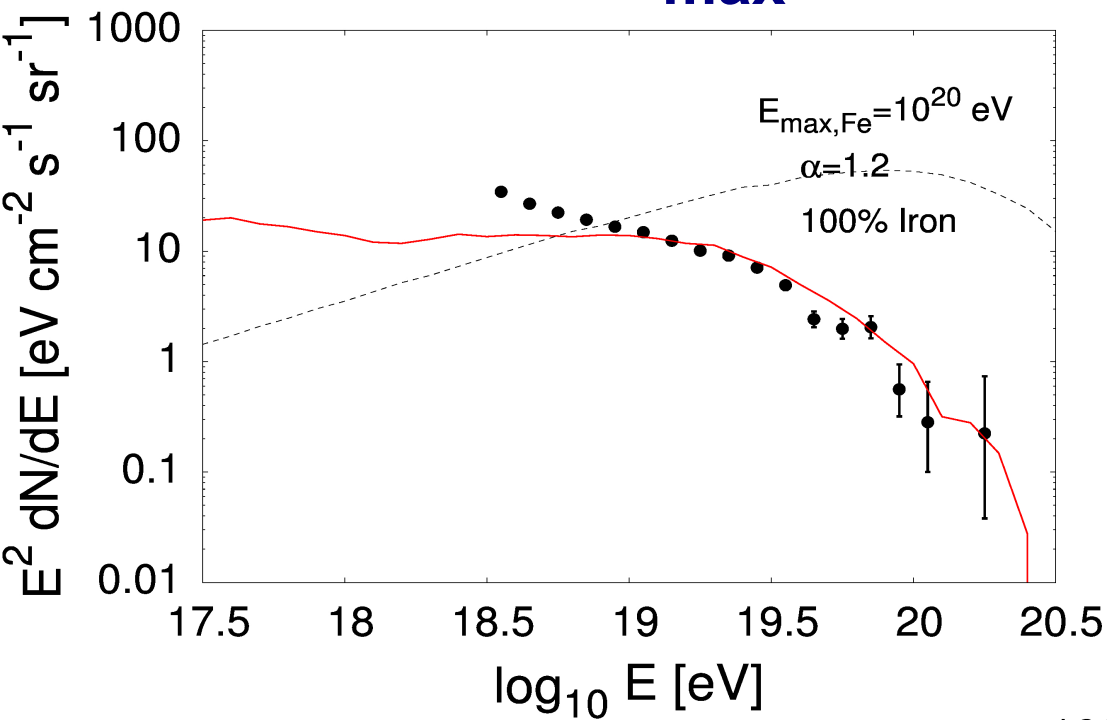


# Extra Slides

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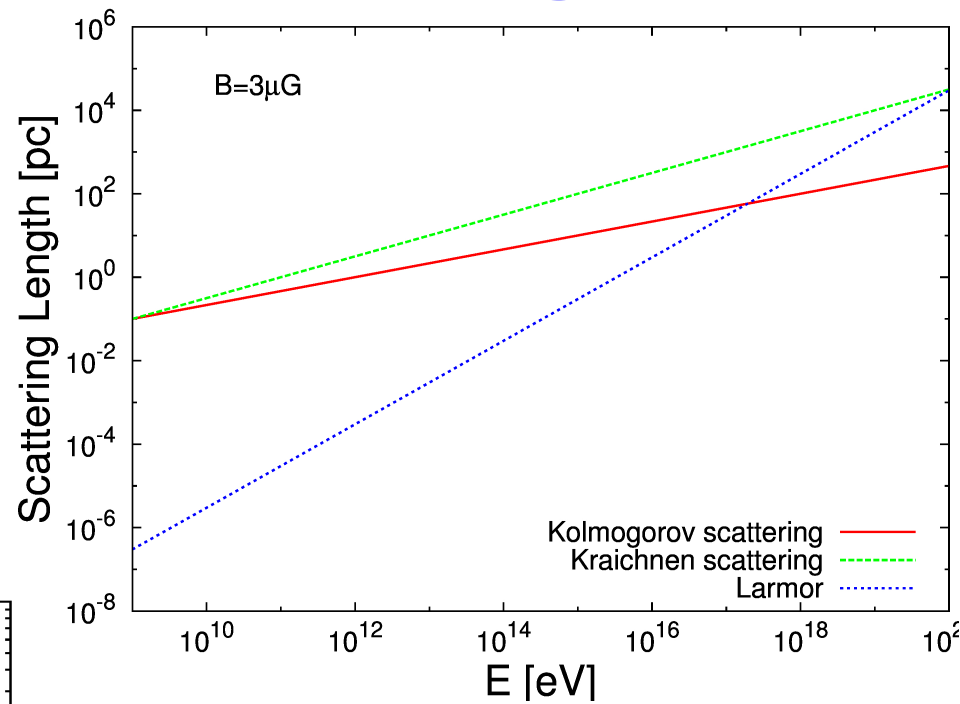
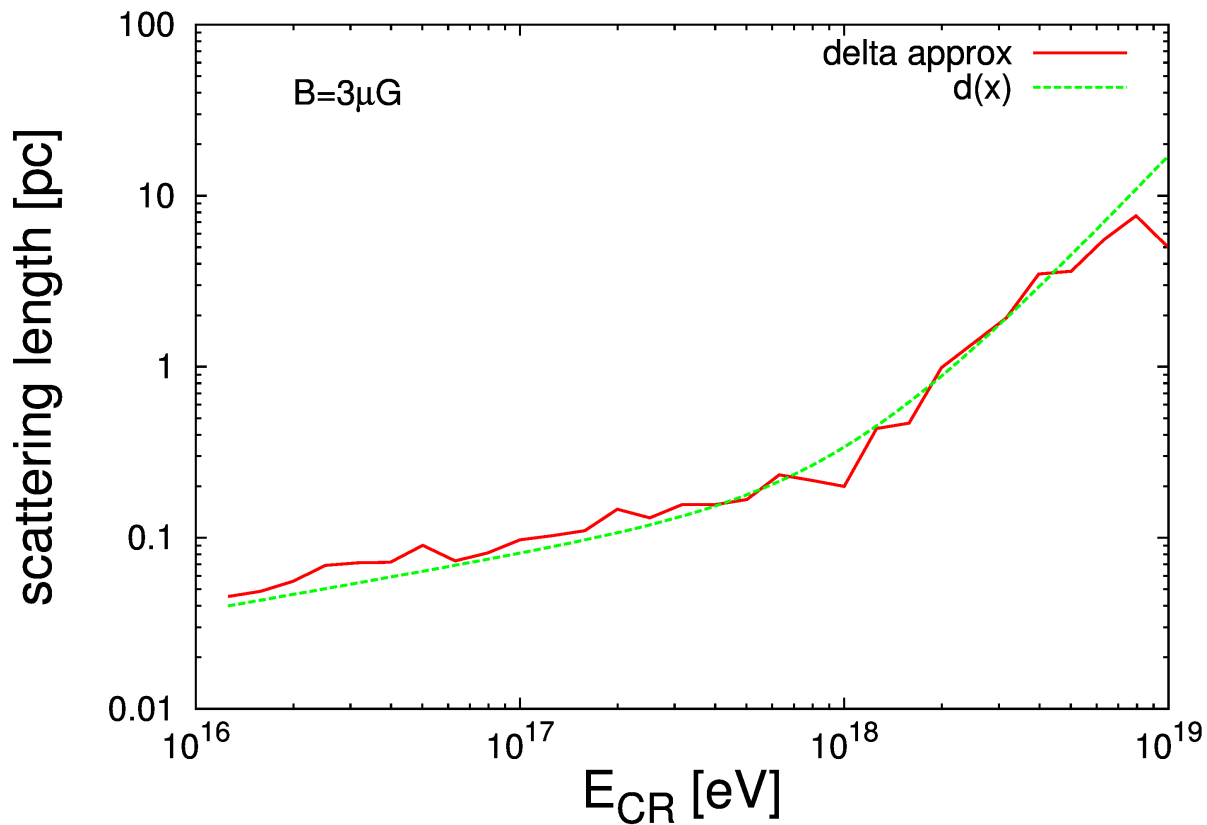


# $E_{\max}$ and $\alpha$ Relation

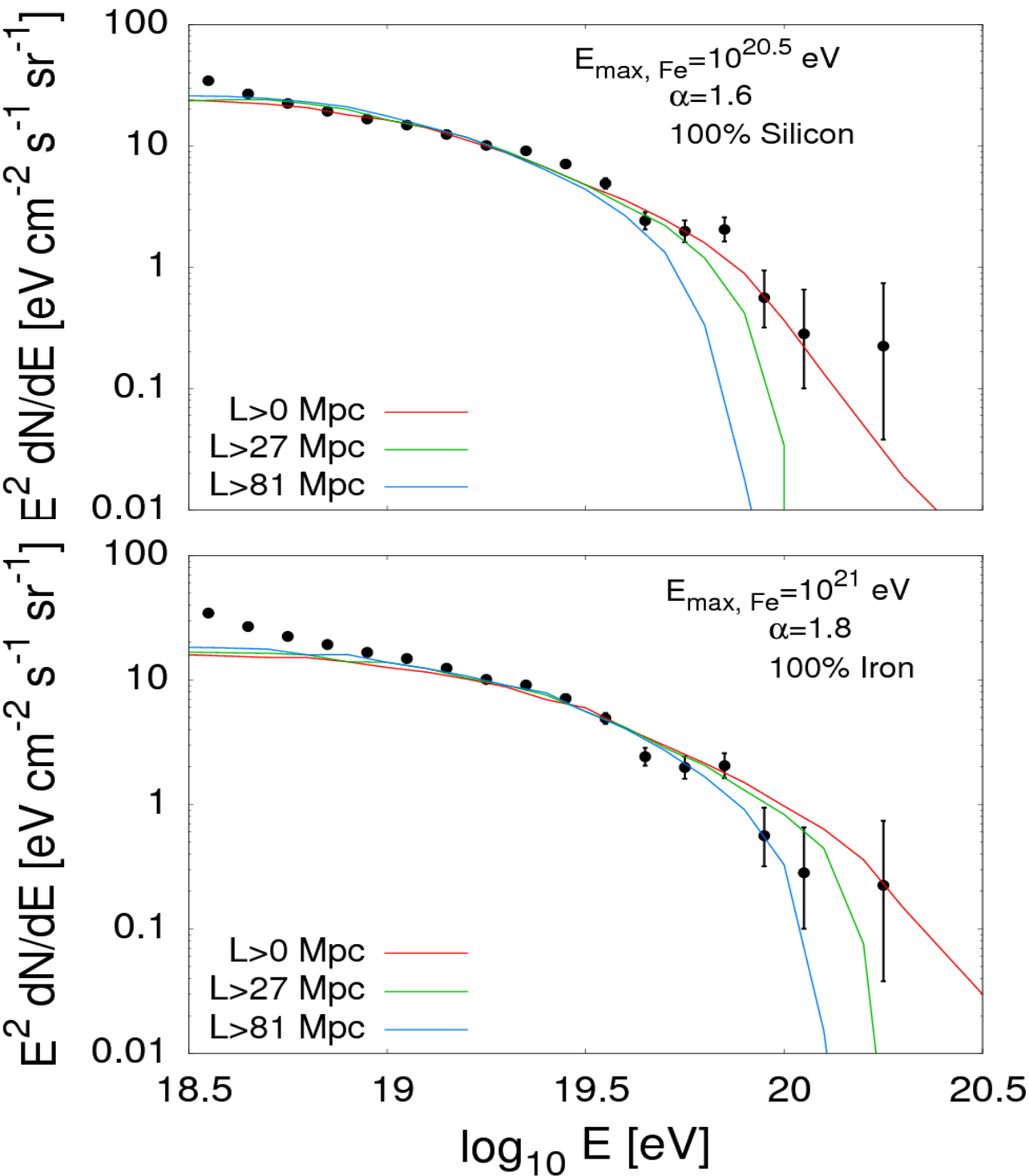




# Prospects for CR Astronomy?



# How Far is the Nearest Source?



**$L > 0 \text{ Mpc}$**

**$L > 27 \text{ Mpc}$**

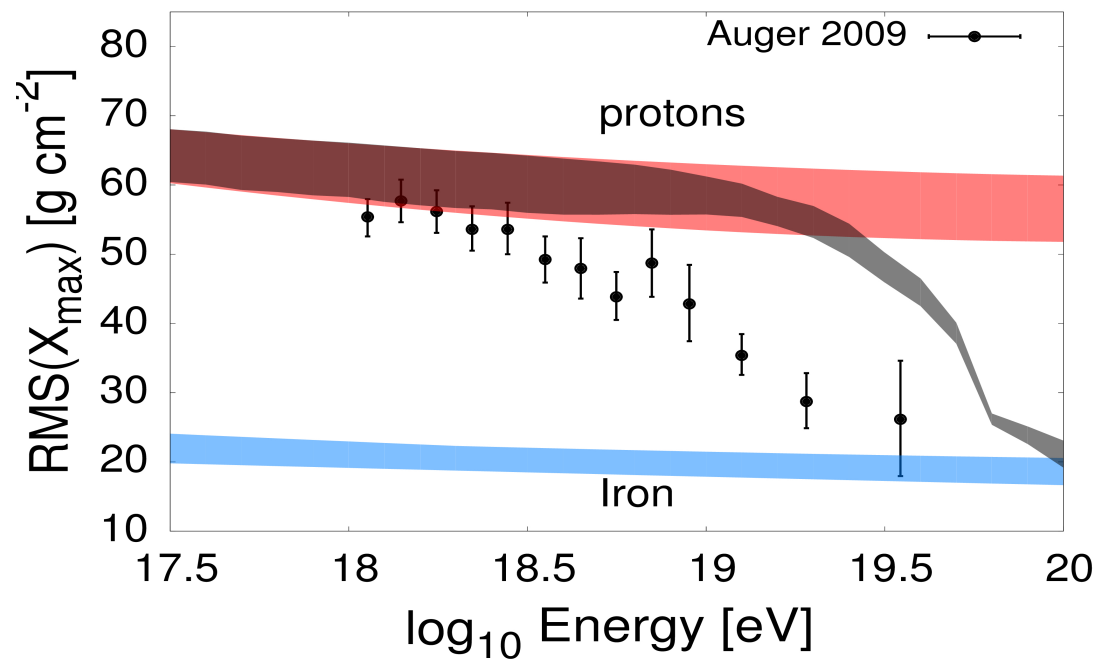
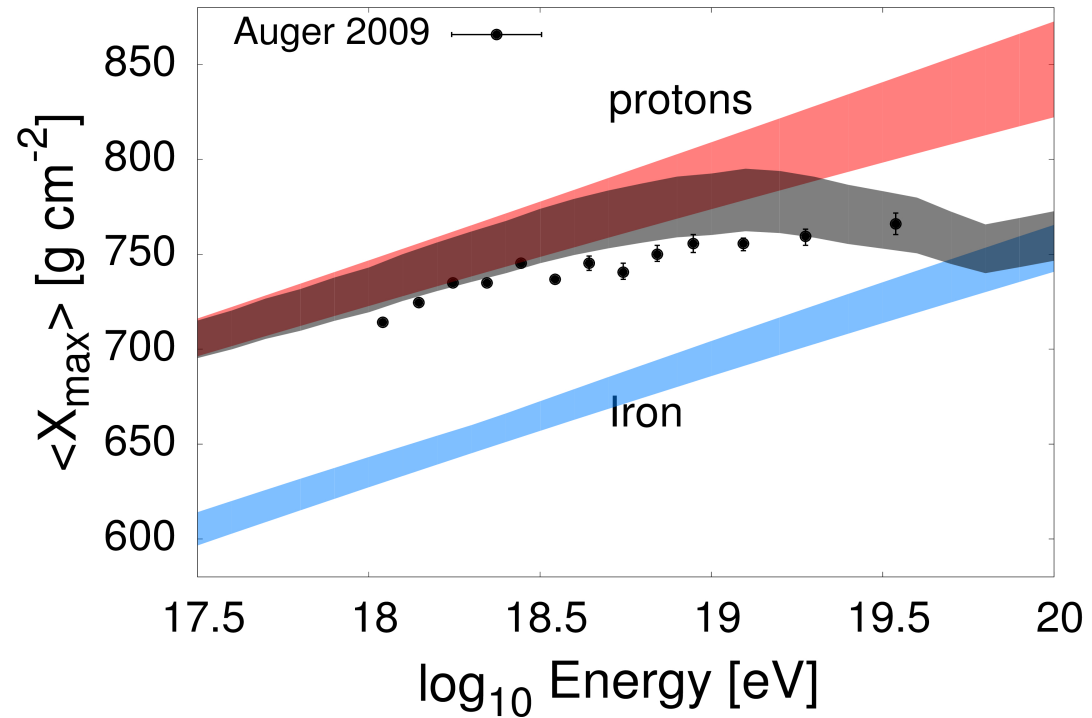
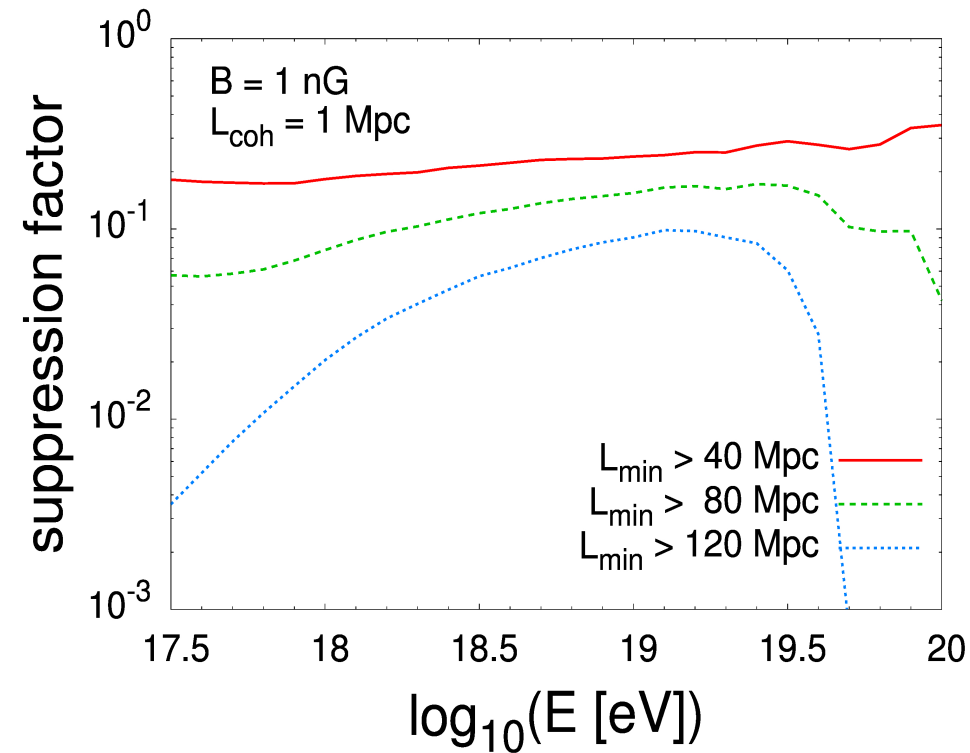
**$L > 81 \text{ Mpc}$**



# Extragalactic Magnetic Fields

$L_{\min} > 81 \text{ Mpc}$

$B = 1 \text{ nG}, L_{\text{coh}} = 1 \text{ Mpc}$



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# The Need for a Heavy Composition

$$\langle X_{\max} \rangle = \sum_{A=1}^{56} f_A X_{\max,A}$$

$$\sigma_{\text{tot}}^2 = \sum_{A=1}^{56} (f_A \sigma_A^2 + f_A (X_{\max,A} - \langle X_{\max} \rangle)^2)$$



# Interaction Rate + Attenuation Rate

## Interaction Rate

$$t_{\text{int.}}^{-1} = \frac{m_p^2}{2E_p^2} \int_0^\infty \frac{n(\epsilon'_\gamma)}{\epsilon'^2_\gamma} d\epsilon'_\gamma \int_0^{2\epsilon'_\gamma \frac{E_p}{m_p}} \epsilon_\gamma \frac{d\sigma}{d\epsilon_\gamma} d\epsilon_\gamma$$

## Attenuation Rate

$$t_{\text{att.}}^{-1} = \frac{m_p^2}{2E_p^2} \int_0^\infty \frac{n(\epsilon'_\gamma)}{\epsilon'^2_\gamma} d\epsilon'_\gamma \int_0^{2\epsilon'_\gamma \frac{E_p}{m_p}} K_{p\gamma} \epsilon_\gamma \frac{d\sigma}{d\epsilon_\gamma} d\epsilon_\gamma$$



# Extragalactic Magnetic Fields

$L > 27 \text{ Mpc}$

$$R_{\text{Larmor}} \approx 4 \left( \frac{E}{10^{20} \text{ eV}} \right) \left( \frac{\text{nG}}{B} \right) \left( \frac{26}{Z} \right) \text{ Mpc}$$

