

Quantum Computing: a future perspective for scientific computing

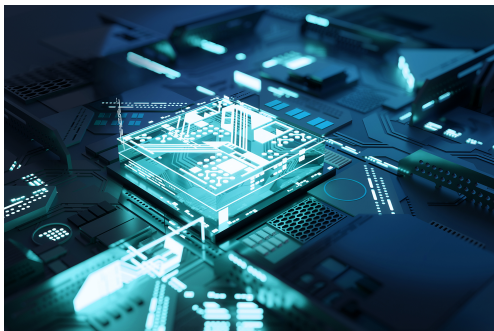
CQTA

Karl Jansen

University of Padova, 10.4.2025



Overview



- > Introduction
- > Quantum Electrodynamics in 1+1 dimensions
- > Real time scattering experiment
- > Optimal Flight Gate Assignment
- > CQTA
- > Conclusion

Why quantum computing

- > **Quantum Biotechnology**, N. Mauranyapin, et.al, arXiv:2111.02021
- > *Emerging quantum computing algorithms for **quantum chemistry***, M. Motta, et.al., arXiv:2109.02873
- > **Quantum Theory Methods** *as a Possible Alternative for the Double-Blind Gold Standard of Evidence-Based Medicine: Outlining a New Research Program*, D.k Aerts, et.al., arXiv:1810.13342
- > **Quantum Battery** *with Ultracold Atoms: Bosons vs. Fermions*, Tanoy Kanti Konar, et.al., arXiv:2109.06816
- > *Hybrid Quantum-Classical Algorithms for **Loan Collection Optimization** with Loan Loss Provisions*, J. Tangpanitanon, et.al, arXiv:2110.15870
- > *A Quantum Natural Language Processing Approach to **Musical Intelligence*** E. Miranda, et.al., arXiv:2111.06741

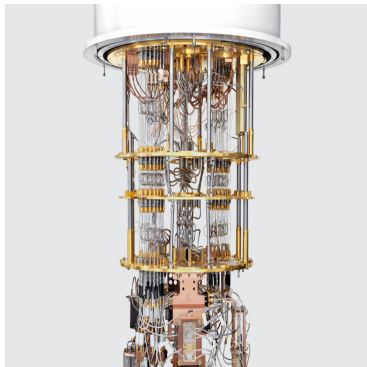
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- > *Developing a Framework for **Sonifying** Variational Quantum Algorithms: Implications for **Music Composition***, Paulo Vitor Itaboraí, Peter Thomas, Arianna Crippa, Karl Jansen, Tim Schwägerl, María Aguado Yáñez, arXiv: 2409.07104

Quantum computer: from the outside



Quantum computer: from the inside



- Shielded to 50,000 times less than Earth's magnetic field
- In a high vacuum: pressure is 10 billion times lower than atmospheric pressure
- Cooled 180 times colder than interstellar space (0.015 Kelvin)
 - prevent quantum noise
- IBMQ: 433 qubits 2022, >1000 qubits 2023, >4000 qubits 2024
 - 10K to 100K error corrected, parallelized
- Google promise: 1.000.000 qubits 2030, 1000 qubits error corrected

How to quantum compute

- > python programming language
 - company provides quantum libraries
- > very convenient setup
 - simulator runs on your local machine
 - hardware usable through quantum cloud service
 - build on reservation system
- > documentation, tutorials and examples available on website, e.g. IBM's textbook: <https://qiskit.org/textbook/preface.html>
 - you can start now!



Schwinger model: 2-dimensional Quantum Electrodynamics

Quantization via Feynman path integral

$$\mathcal{Z} = \int \mathcal{D}A_\mu \mathcal{D}\Psi \mathcal{D}\bar{\Psi} e^{-S_{\text{gauge}} - S_{\text{ferm}}}$$

Fermion action

$$S_{\text{ferm}} = \int d^2x \bar{\Psi}(x) [D_\mu + m] \Psi(x)$$

gauge covariant derivative

$$D_\mu \Psi(x) \equiv (\partial_\mu - ig_0 A_\mu(x)) \Psi(x)$$

with A_μ gauge potential, g_0 bare coupling

$$S_{\text{gauge}} = \int d^2x F_{\mu\nu} F_{\mu\nu}, \quad F_{\mu\nu}(x) = \partial_\mu A_\nu(x) - \partial_\nu A_\mu(x)$$

Lattice Schwinger model

introduce a 2-dimensional lattice with lattice spacing a

fields $\Psi(x)$, $\bar{\Psi}(x)$ on the lattice sites $x = (t, \mathbf{x})$ integers

discretized fermion action

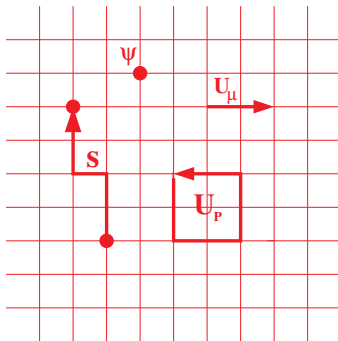
$$S \rightarrow a^2 \sum_x \bar{\Psi} [\gamma_\mu \partial_\mu - r \underbrace{\partial_\mu^2}_{\nabla_\mu^* \nabla_\mu} + m] \Psi(x)$$

$$\partial_\mu = \frac{1}{2} [\nabla_\mu^* + \nabla_\mu]$$

discrete derivatives

$$\nabla_\mu \Psi(x) = \frac{1}{a} [\Psi(x + a\hat{\mu}) - \Psi(x)] , \quad \nabla_\mu^* \Psi(x) = \frac{1}{a} [\Psi(x) - \Psi(x - a\hat{\mu})]$$

second order derivative \rightarrow remove doubler \leftarrow break chiral symmetry



Implementing gauge invariance

Wilson's fundamental observation: introduce Paralleltransporter connecting the points x and $y = x + a\hat{\mu}$:

$$U(x, \mu) = e^{iaA_\mu(x)} \in U(1)$$

\Rightarrow lattice derivatives

$$\begin{aligned}\nabla_\mu \Psi(x) &= \frac{1}{a} [U(x, \mu) \Psi(x + \mu) - \Psi(x)] \\ \nabla_\mu^* \Psi(x) &= \frac{1}{a} [\Psi(x) - U^{-1}(x - \mu, \mu) \Psi(x - \mu)]\end{aligned}$$

action gauge invariant under

$$\begin{aligned}\Psi(x) &\rightarrow g(x) \Psi(x), \quad \bar{\Psi}(x) \rightarrow \bar{\Psi}(x) g^*(x), \\ U(x, \mu) &\rightarrow g(x) U(x, \mu) g^*(x + \mu)\end{aligned}$$

The Schwinger model: the action

- > Wilson's fundamental observation: introduce parallel transporter connecting the points x and $y = x + a\hat{\mu}$:

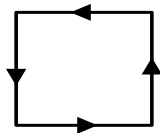
$$U(x, \mu) = e^{iaA_\mu(x)} \in U(1)$$

- > lattice derivative: $\nabla_\mu \Psi(x) = \frac{1}{a} [U(x, \mu)\Psi(x + \mu) - \Psi(x)]$

- > plaquette action

$$U_p = U(x, \mu)U(x + \mu, \nu)U^\dagger(x + \nu, \mu)U^\dagger(x, \nu)$$

$$\rightarrow F_{\mu\nu}F^{\mu\nu}(x) \quad \text{for} \quad a \rightarrow 0$$



$$S = a^2 \sum_x \left\{ \beta (= \frac{1}{g_0^2}) [1 - \text{Re}(U_{(x,p)})] + \bar{\Psi}(x) \left[m + \frac{1}{2} \{ \gamma_\mu (\nabla_\mu + \nabla_\mu^*) - a \nabla_\mu^* \nabla_\mu \} \right] \Psi \right\}$$

Physical observables

expectation value of physical observables \mathcal{O}

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int_{\text{fields}} \mathcal{O} e^{-S}$$



lattice discretization



01011100011100011110011



Going to Quantum Chromodynamics

change	factor
2-d (40×40) \rightarrow 4-d (40^4)	1600
gauge field $U(x, \mu) \in U(1) \rightarrow U(x, \mu) \in SU(3)$	8
Pauli matrices $\sigma_\mu \rightarrow$ Gell-Mann-matrices γ_μ	4
spinors become 12-component complex vectors	24

> total factor: 1228800

> theory needs renormalization

Present day Lattice Gauge Theory calculations

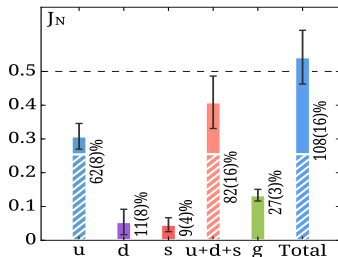
- > Large activity of theoretical lattice simulations of the standard model of particle interaction
→ see FLAG review, Y. Aoki et.al, Eur.Phys.J.C 82 (2022) 10, 869
- > discretize space and time
- > use euclidean metric
- > Feynman path integral → statistical mechanical system
- > Lattice simulations of QCD, Higgs-Yukawa sector, Supersymmetry via **Markov Chain Monte Carlo Methods**

A Towards resolving the spin puzzle of the nucleon

(C. Alexandrou, M. Constantinou, K. Hadjiyiannakou,

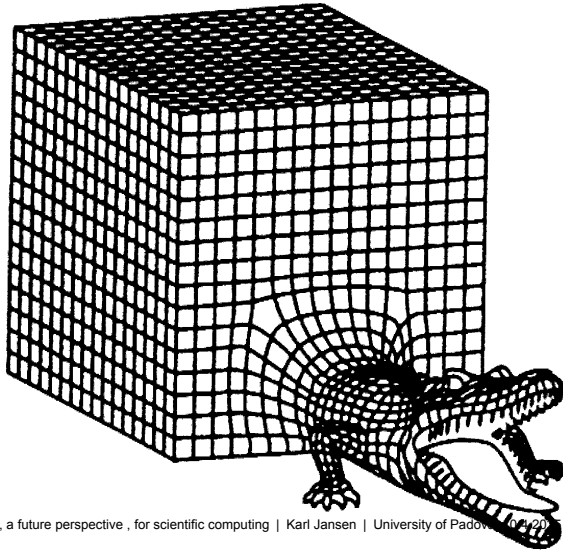
C. Kallidonis, G. Koutsou, A. Vaquero Avilés-Casco, C. Wiese, K. Jansen)

- > old puzzle: quarks provide only surprisingly small contribution to spin
 - remained unsolved for decades
- > lattice gauge theory advances
 - very demanding, dedicated effort
 - including four lightest quarks **and gluon** → obtain full spin decomposition



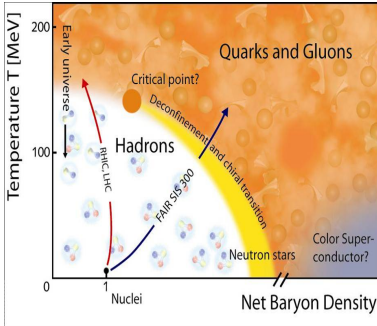
- stripped segments: valence quarks
- solid segments: sea quarks and gluons
- find large gluon contribution

There are dangerous lattice animals



Understanding the early universe

- > Markov Chain Monte Carlo: only zero baryon density accessible
 - understanding of phase transitions?
 - early universe
 - heavy ion experiments
 - exotic regions of PD
- > do not understand origin of today's universe
- > topological terms
- > real time evolution
- > large autocorrelations towards continuum limit
- solution: Hamiltonian formulation



A problem with Hamiltonian approach



- determine wave function $|\Psi\rangle$

$$|\Psi\rangle = \sum_{i_1, i_2, \dots, i_N} C_{i_1, i_2, \dots, i_N} |i_1 i_2 \dots i_N\rangle$$

C_{i_1, i_2, \dots, i_N} coefficient matrix with 2^N entries

→ problem scales exponentially

Two solutions for Hamiltonian approach

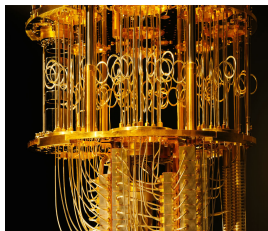
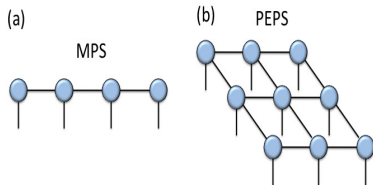
> Tensor Networks

- gapped, local Hamiltonians
 - area law
 - very fast (exponential) convergence
- costly: entanglement
 - phase transitions, (long) real time evolution



> Quantum Computing

- conceptually clean path
- noise, low number of qubits, ...



> Here: focus on quantum computing Quantum Field Theories

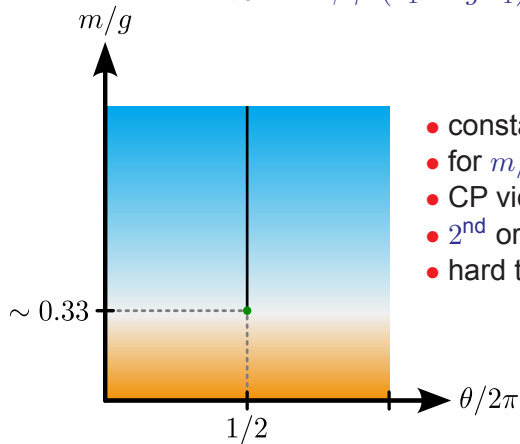
Schwinger Model

(Schwinger 1962)

- > existence of bound states (mass gap)
 - > asymptotic free ($g_0 \rightarrow 0$ for distance between charges going to zero)
 - > exactly solvable for zero fermion mass (Coleman)
 - > super-renormalizable
 - > with topological term: interesting CP-violating phase transition
- ⇒ valuable test laboratory for QCD

Schwinger model in the continuum and phase diagram

$$\mathcal{H} = -i\bar{\psi}\gamma^1(\partial_1 - igA_1)\psi + m\bar{\psi}\psi + \frac{1}{2}\left(\dot{A}_1 + \frac{g\theta}{2\pi}\right)^2$$



- constant external electric field: θ
- for $m/g > 0.33$, 1st order phase transition
- CP violating
- 2nd order endpoint at $m/g = 0.33$
- hard to explore with Monte Carlo methods

Schwinger model on the lattice: Wilson fermions

(Takis Angelides, Arianna Crippa, Lena Funcke, Karl Jansen,
Stefan Kühn, Pranay Naredi, Ivano Tavernelli, Derek Wang, arxiv:2312.12831)

> Wilson Hamiltonian

$$H_W = \sum_{n=0}^{N-2} \left(\bar{\phi}_n \left(\frac{1 + i\gamma^1}{2a} \right) U_n \phi_{n+1} + \text{h.c.} \right) \\ + \sum_{n=0}^{N-1} \left(m_{\text{lat}} + \frac{1}{a} \right) \bar{\phi}_n \phi_n + \sum_{n=0}^{N-2} \frac{ag^2}{2} (L_n + l_0)^2.$$

> mass m_{lat} ; coupling g ; lattice spacing a ; electric field $l_0 = \frac{\theta}{2\pi}$

> Link operator $U_\mu = e^{igA_\mu}$, A_μ gauge potential

Pauli representation through Jordan-Wigner transformation

- > Jordan-Wigner transformation

$$\phi_{n,\alpha} \rightarrow \chi_{2n - \lfloor \frac{\alpha}{2} \rfloor + 1}$$

$$\chi_n = \prod_{k < n} (iZ_k) \sigma_n^-$$

- > (dimensionless) Wilson Hamiltonian, $x = 1/(ag)^2$
→ open boundary conditions: eliminate gauge fields

$$W_W = x \sum_{n=0}^{N-2} (X_{2n+2} X_{2n+3} + Y_{2n+2} Y_{2n+3}) + \left(\frac{m_{\text{lat}}}{g} \sqrt{x} + x \right) \sum_{n=0}^{N-1} (X_{2n+1} X_{2n+2} + Y_{2n+1} Y_{2n+2}) + \sum_{n=0}^{N-2} (l_0 + \sum_{k=0}^n Q_k)^2$$

Pauli representation through Jordan-Wigner transformation

- > electric field density operator

$$L_W = \sum_{k=0}^{N-1} Q_k = \sum_{k=0}^N \phi_n^\dagger \phi_n$$

→ JW-transformation:
$$L_W = l_0 + \frac{1}{2} \sum_{k=0}^{\lceil N/2 \rceil - 1} (Z_{2k} + Z_{2k+1})$$

- > particle number operator

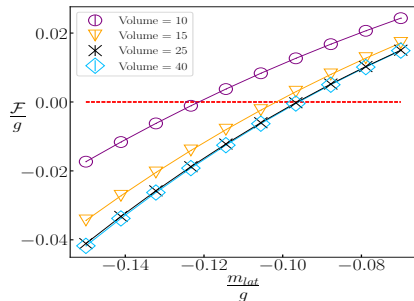
$$P_W = N + \frac{1}{2} \sum_{n=0}^{N-1} \phi \phi$$

→ JW-transformation:
$$P_W = N + \frac{1}{2} \sum_{n=0}^{N-1} (X_{2n+1} X_{2n+2} + Y_{2n+1} Y_{2n+2})$$

Determining the mass shift: a MPS calculation

- > electric field density (EFD) in mass perturbation theory

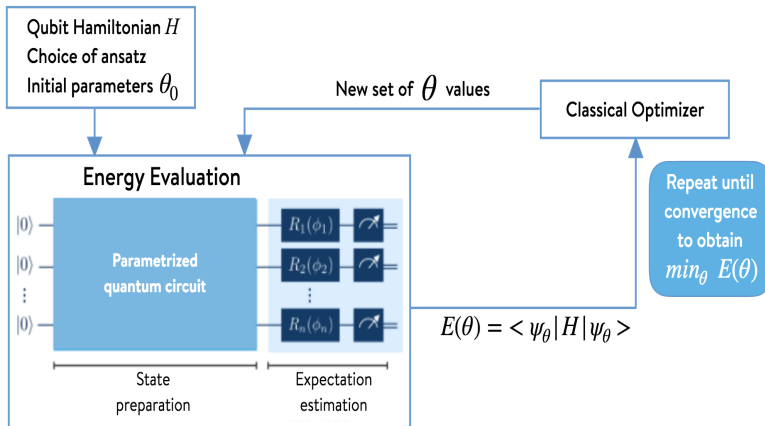
$$\frac{\mathcal{F}}{g} = \frac{e^\gamma}{\sqrt{\pi}} \left(\frac{m}{g} \right) \sin \theta - 8.9139 \frac{e^{2\gamma}}{4\pi} \left(\frac{m}{g} \right)^2 \sin(2\theta) \Rightarrow \text{for } m = 0 \text{ EFD vanishes}$$



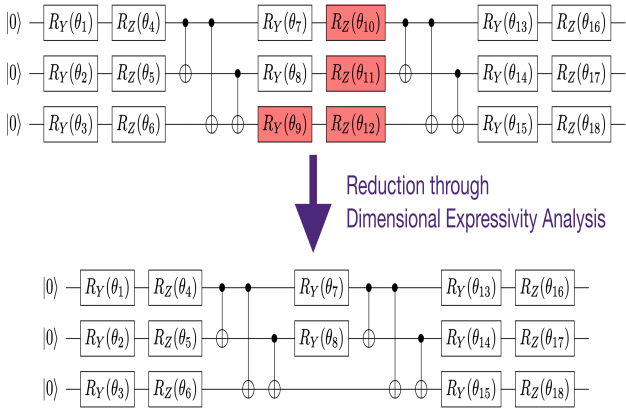
Wilson

Variational Quantum Eigensolver (VQE)

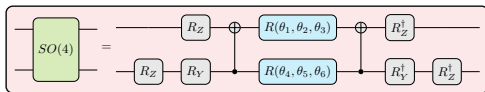
- > a hybrid quantum/classical variational approach



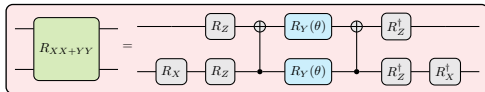
Example for a quantum circuit



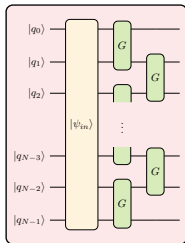
The Ansatz



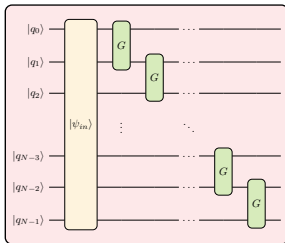
(a)



(b)



(c)



(d)

- decomposition of $SO(4)$ and R_{XX+YY} gates

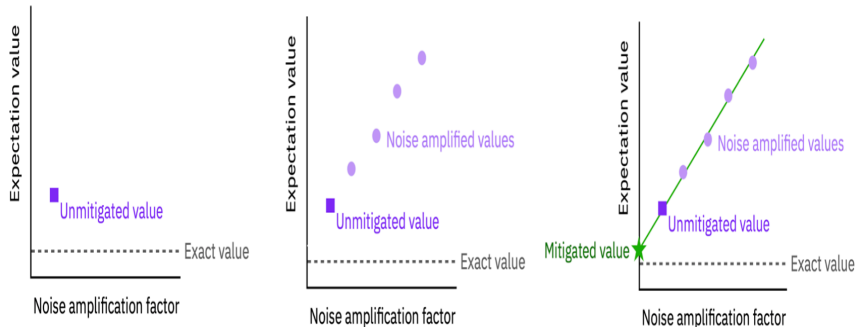
- brick and ladder ansatz

Mitigating quantum computing results

- > zero noise extrapolation (ZNE) in theory

$$\langle \psi | O | \psi \rangle = \langle 0 | U^\dagger O U | 0 \rangle$$

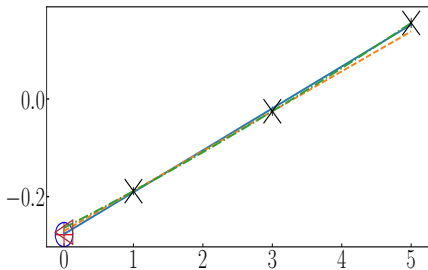
$$|\psi\rangle = U|0\rangle = UU^\dagger U|0\rangle = UU^\dagger UU^\dagger U|0\rangle$$



ZNE in practise

(c): Staggered $\frac{\mathcal{F}}{g}$

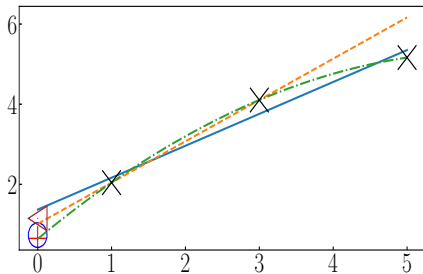
$$N = 6, m_{lat}/g = 0, l_0 = 0.65$$



when it works

(f): Wilson PN

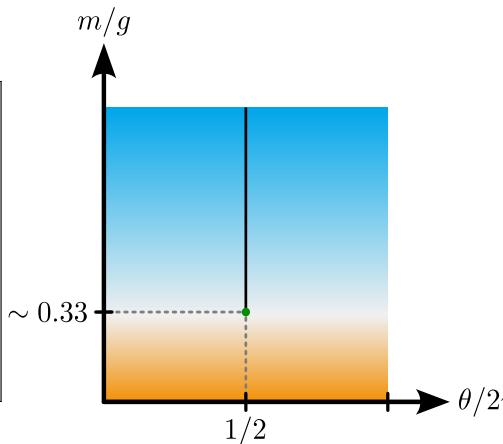
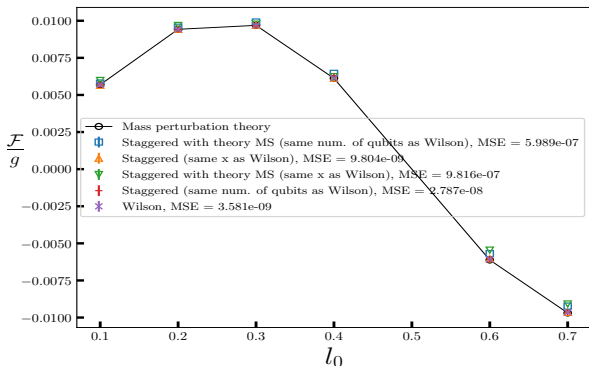
$$N = 6, m_{lat}/g = 0, l_0 = 0.475$$



when not

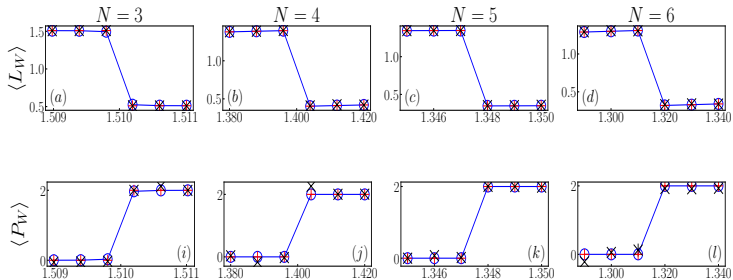
Results: small mass

> results for $m_r/g = 0.01$ (remember: $x = 1/(ag)^2$)

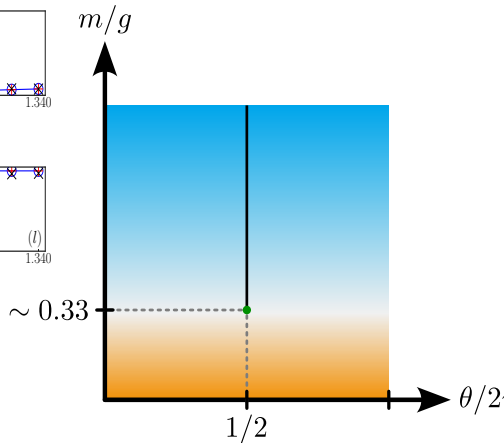


• blue circles: exact diagonalization, red pluses: exact simulations, black crosses: quantum hardware

Results: large mass $m_r/g = 10$



● including hardware results



The Hamiltonian of 2+1 dimensional QED \hat{H}_{QED}

> Electric field operator: $\frac{g^2}{2} \sum_{\vec{n}} \left(\hat{E}_{\vec{n},x}^2 + \hat{E}_{\vec{n},y}^2 \right)$

> Plaquette operator: $-\frac{1}{2a^2 g^2} \sum_{\vec{n}} \left(\hat{P}_{\vec{n}} + \hat{P}_{\vec{n}}^\dagger \right)$

> mass term $+m \sum_{\vec{n}} (-1)^{n_x+n_y} \hat{\phi}_{\vec{n}}^\dagger \hat{\phi}_{\vec{n}}$

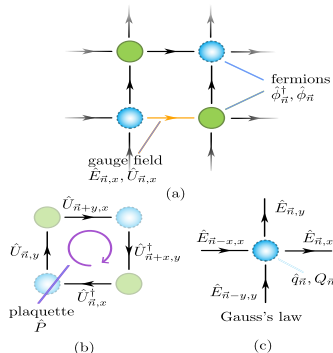
> kinetic term $\hat{U}_{\vec{n},x} = e^{iag\hat{A}_{\vec{n},x}}$

$$\frac{i}{2a} \sum_{\vec{n}} \left(\hat{\phi}_{\vec{n}}^\dagger \hat{U}_{\vec{n},x}^\dagger \hat{\phi}_{\vec{n}+x} - \text{h.c.} \right) \\ - \frac{(-1)^{n_x+n_y}}{2a} \sum_{\vec{n}} \left(\hat{\phi}_{\vec{n}}^\dagger \hat{U}_{\vec{n},y}^\dagger \hat{\phi}_{\vec{n}+y} + \text{h.c.} \right)$$

> investigated the running coupling with matching to Markov Chain Monte Carlo

> looked at electric flux configurations for Coulomb, confinement and string breaking regimes of static potential

> developed lattice Chern-Simons Hamiltonian

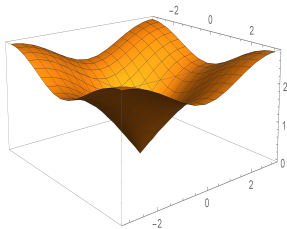


Chern-Simons term in 2+1 dimensional QED

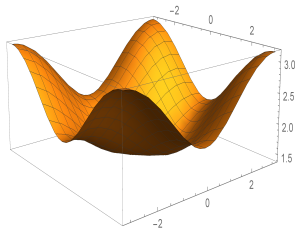
(C. Peng, C. Diamantini, L. Funcke, A. Hassan, K. Jansen, Stefan Kühn, D. Luo, P. Naredi, arxiv:2407.20225)

$$\hat{H} = \sum_{x \in \text{sites}} \frac{e^2}{2a^2} \left[\left(\hat{p}_{x;1} - \frac{ka^2}{4\pi} \hat{A}_{x-\hat{2};2} \right)^2 + \left(\hat{p}_{x;2} + \frac{ka^2}{4\pi} \hat{A}_{x-\hat{1};1} \right)^2 \right] + \frac{1}{2e^2} \left(\square \hat{A}_{x;1,2} \right)^2$$

> energy bands



pure Maxwell theory
massless photon

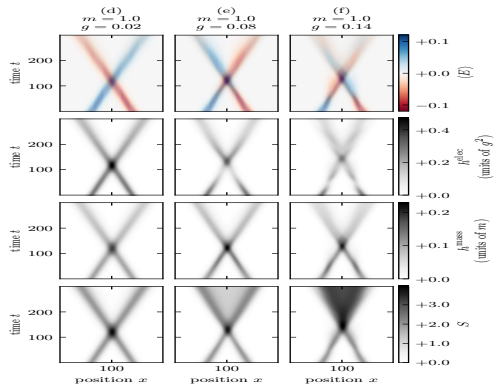


adding Chern-Simons term
topological mass generation

> opens door to investigate e.g. fermion/boson dualities, fractional quantum Hall effect,

Scattering from MPS: real mesons 1+1 dimensional Schwinger model

Marco Rigobello, Simone Notarnicola, Giuseppe Magnifico, Simone Montangelo,
Phys. Rev. D 104, 114501



- meson wave packets
- generate entanglement
- strong dependence on coupling
- rich phenomenology after collision

Scattering on a quantum computer

(Yahui Chai, Arianna Crippa, Karl Jansen, Stefan Kühn, Ivano Tavernelli, Francesco Tacchino, arxiv:2312.02272)

> Continuum Lagrangian of Thirring model

$$\mathcal{L} = i\bar{\psi}\gamma^\mu\partial_\mu\psi - m\bar{\psi}(x)\psi(x) - \frac{\lambda}{2}(\bar{\psi}\gamma_\mu\psi)(\bar{\psi}\gamma^\mu\psi)$$

> Hamiltonian lattice version

$$H = \sum_{n=0}^{N-1} \left\{ \frac{i}{2a} \left(\xi_{n+1}^\dagger \xi_n - \xi_n^\dagger \xi_{n+1} \right) + (-1)^n m \xi_n^\dagger \xi_n \right\} + \sum_{n=0}^{N-1} \frac{g(\lambda)}{a} \xi_n^\dagger \xi_n \xi_{n+1}^\dagger \xi_{n+1}$$

> Spin representation → Jordan Wigner

Spin representation

> Jordan-Wigner transformation

$$\xi_n^\dagger = \prod_{l < n} \sigma_l^z \sigma_n^-, \quad \xi_n = \prod_{l < n} \sigma_l^z \sigma_n^+$$

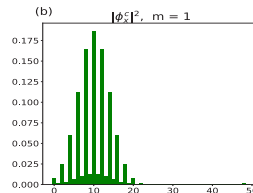
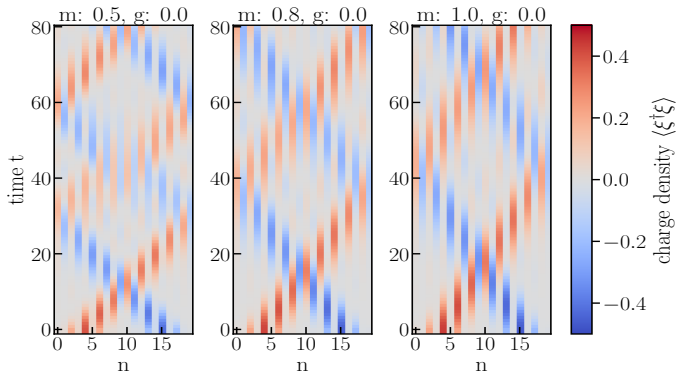
$$\sigma_l^\pm = (\sigma_l^x \pm i\sigma_l^y) / 2$$

> Hamiltonian

$$\begin{aligned} H = & \frac{i}{2a} \sum_{n=0}^{N-2} (\sigma_{n+1}^- \sigma_n^+ - \sigma_n^- \sigma_{n+1}^+) \\ & + \frac{i}{2a} (\sigma_0^- \sigma_1^z \cdots \sigma_{N-2}^z \sigma_{N-1}^+ - \sigma_{N-1}^- \sigma_{N-2}^z \cdots \sigma_1^z \sigma_0^+) \\ & + \frac{m}{2} \sum_{n=0}^{N-1} (-1)^n (\mathbb{1} - \sigma_n^z) + \frac{g}{4a} \sum_{n=0}^{N-1} (\mathbb{1} - \sigma_n^z) (\mathbb{1} - \sigma_{n+1}^z) \end{aligned}$$

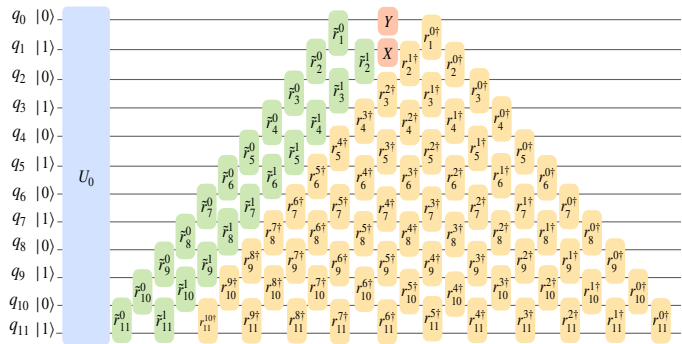
Gaussian wave packets

- > Gaussian wave packets $\phi_k^{c(d)} = \frac{1}{\mathcal{N}_k^{c(d)}} e^{-ik\mu_n^{c(d)}} e^{-(k-\mu_k^{c(d)})^2/4\sigma_k^2}$
- > time evolution: Givens rotation
- > time evolution for free fermions: charge distribution

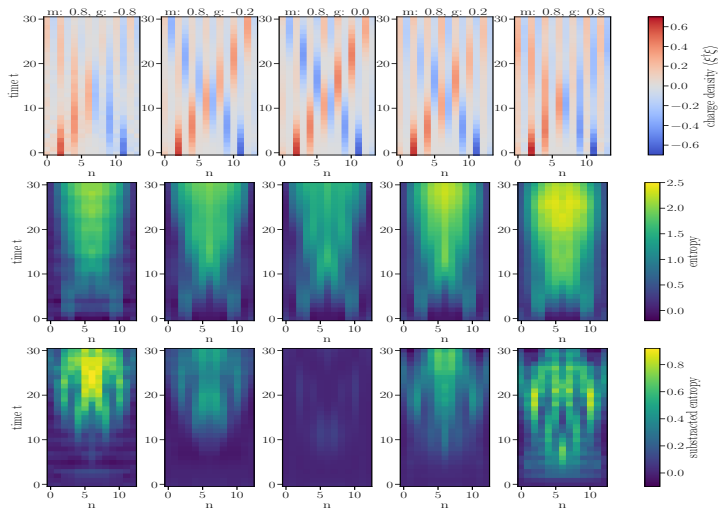


Quantum circuit

- > blue box: vacuum preparation
- > green and yellow boxes: wave packet preparation and time evolution

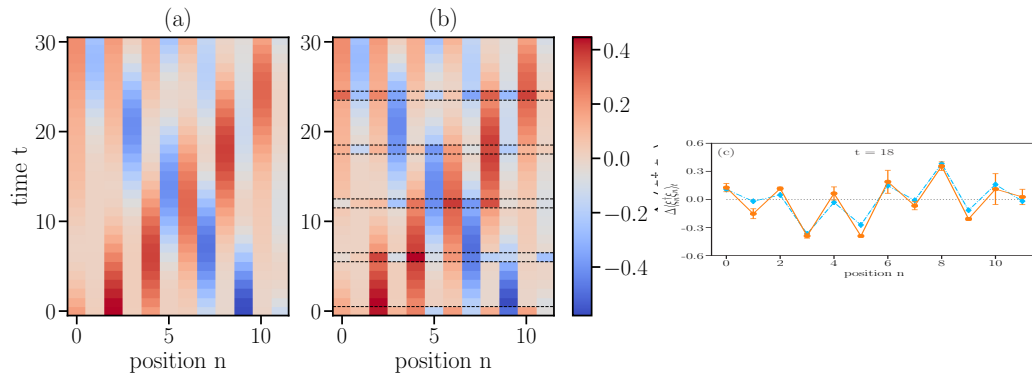


Interacting case



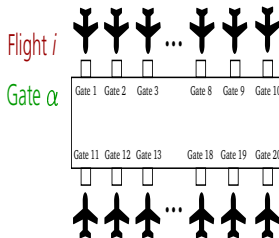
Hardware runs

> Ideal versus hardware



Quantum computing the flight gate assignment problem

- > A classical optimization problem: flight gate assignment
(Y. Chai, L. Funcke, T. Hartung, S. Kühn, T. Stollenwerk, P. Stornati, K. Jansen, arXiv:2302.11595)
- > Find shortest path between connecting flights
- > Different incoming and outgoing flights need to be assigned to gates
→ find optimal assignment
- > Classical optimization problem
→ quantum advantage?



Quantum computing the flight gate assignment problem

- > binary variables encoding gates and flights

$$x_{i\alpha} = \begin{cases} 1, & \text{if flight } i \in F \text{ is assigned to gate } \alpha \in G \\ 0, & \text{otherwise} \end{cases}$$

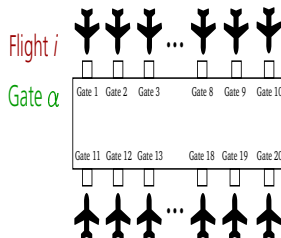
$x \in \{0, 1\}^{F \otimes G} \rightarrow x$ binary variable $\rightarrow x \in \{-1, 1\}$

eigenstate of third Pauli matrix σ_z

- > leads to mathematical description of Hamiltonian

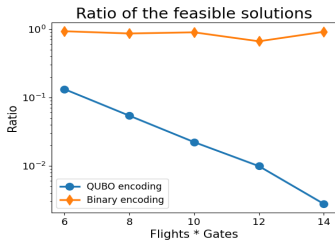
$$H = \sum_{j=1}^n Q_{jj} \sigma_j^z + \sum_{\substack{j,k=1 \\ j < k}}^n Q_{jk} \sigma_j^z \otimes \sigma_k^z$$

- > Task: find lowest energy \Leftrightarrow shortest path
- > Same mathematical description for problems in **traffic, logistics, particle tracking,**
...



Quantum computing the flight gate assignment problem

- > Started with QUBO implementation
- > Implementation of various improvements
 - using binary encoding
 - reformulation of Hamiltonian through projectors
 - Using Conditional Value at Risk (CVaR)

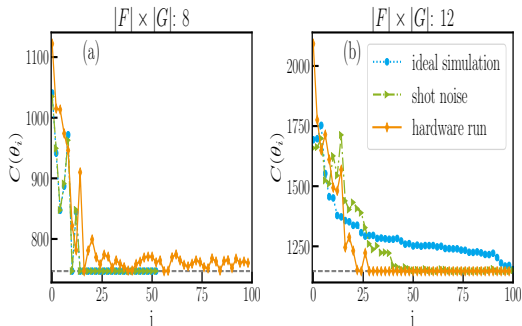


Feasible ratio

Quantum hardware runs of flight gate assignment problem

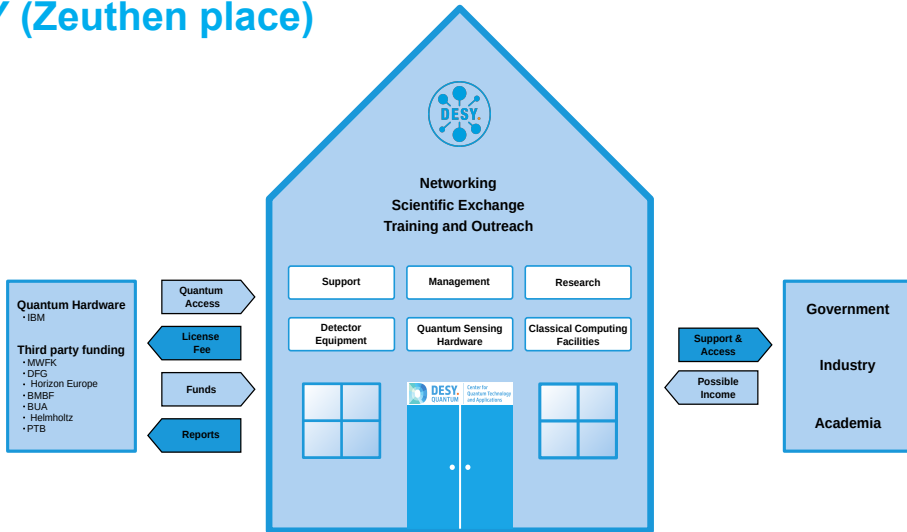
(Y. Chai, E. Epifanovsky, K. Jansen, A. Kaushik, S. Kühn, arxiv:2309.09686)

- > hardware runs on IonQ's Aria trapped ion quantum computer
- > circuit: efficientSU2
- > real VQE and inference runs

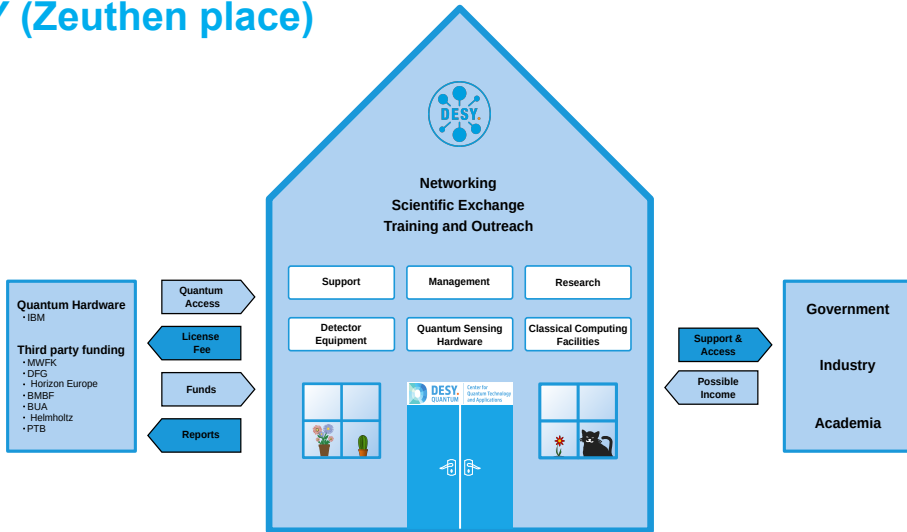


Convergence

Center for Quantum Technology and Applications at DESY (Zeuthen place)



Center for Quantum Technology and Applications at DESY (Zeuthen place)



The CQTA group

> The group in Zeuthen in September 2021



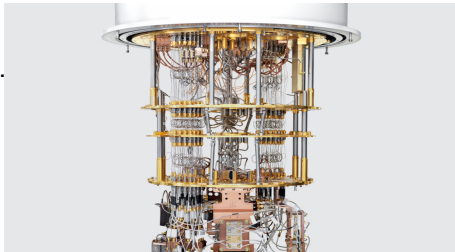
The CQTA group

> The present group in Zeuthen (missing 3 female members)



Center for Quantum Technology and Applications

- > **Quantum Field Theoretical models from condensed matter and high energy physics** → sign problem, real time phenomena
- > **Optimization/classification**
 - Particle track reconstruction/jet classification
 - Flight gate assignment
 - Gene/exon classification
- > **Quantum art**
 - Quantum music, Quantum painting, WS 9.7.-11.7.
- > **Others**
 - factoring, Feynman diagrams, matrix models, ...
- > **training**
- > **Algorithm development**
 - Expressivity
 - controllability
 - warm starts



ERA Chair QUEST (QUantum computing for Excellence in Science and Technology)

- > European Research Executive Agency funding (2.5 million Euro)
- > focus activities
 - Building up a quantum computing group at the Cyl
 - develop applications of uses case for industry, governmental agencies and academia
 - Act as hub for Eastern Mediterranean region
 - closely connected to Center for Quantum Technology and Applications (CQTA) at DESY



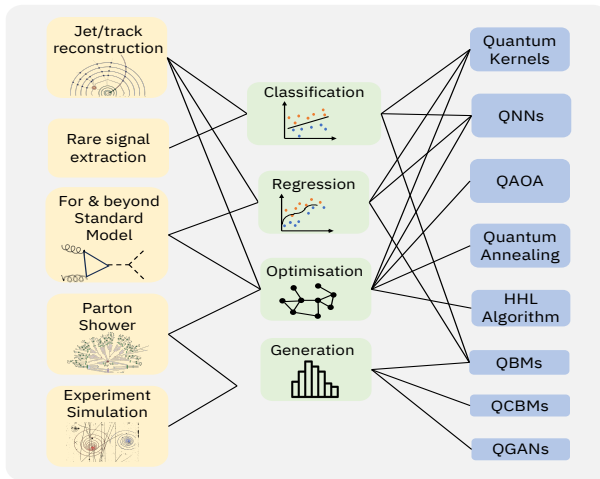
QC4HEP whitepaper, arXiv:2307.03236

Alberto Di Meglio,^{1,*} Karl Jansen,^{2,3,†} Ivano Tavernelli,^{4,‡} Constantia Alexandrou,^{5,3} Srinivasan Arunachalam,⁶
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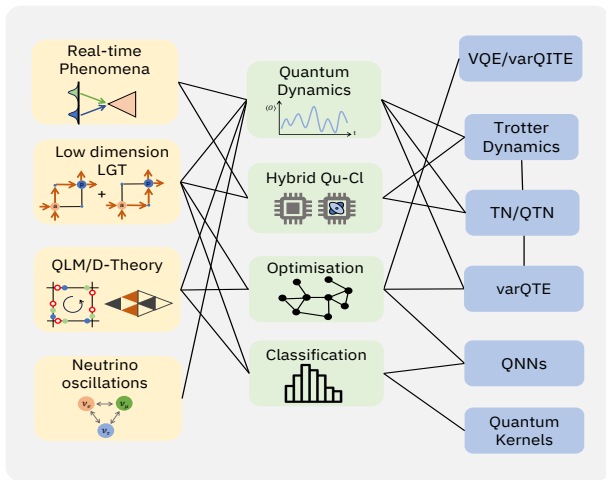
Abstract

Quantum computers offer an intriguing path for a paradigmatic change of computing in the natural sciences and beyond, with the potential for achieving a so-called quantum advantage, namely a significant (in some cases exponential) speed-up of numerical simulations. In particular, the high-energy physics community plays a pivotal role in accessing the power of quantum computing, since the field is a driving source for challenging computational problems. ...

QC4HEP: Experiment summary

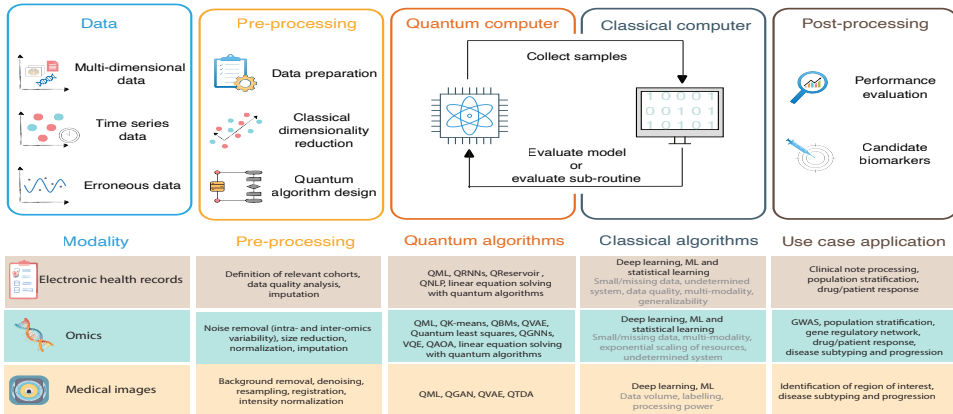


QC4HEP: Theory summary



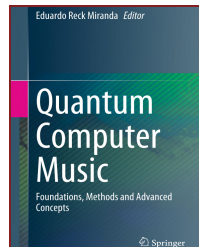
Quantum computing enhances biomarker discovery

(Frederik Flöther, Daniel Blankenberg, Maria Demidik, Karl Jansen, Raga Krishnakumar, Rajiv Krishnakumar, Numan Laanait, Laxmi Parida, Carl Saab, Filippo Utro, arXiv:2411.10511)



Summary and outlook

- > It took 40 years to start realizing Feynman's vision of using quantum computers
- > **Now:** first computations in high energy physics with $O(10)$ qubits on NISQ devices
 - experiment: particle tracking, Boltzmann machines, quantum neural networks, ...
 - theory: low-dimensional, abelian and non-abelian models in 1+1 and 2+1 dimensions, scattering, ...
- > **soon:** demonstrations, $O(100)$ qubits and circuit depth of $O(100)$
 - identify and evaluate applications for quantum computers
 - develop further quantum algorithms and methods
 - evaluate scaling with the number of qubits
 - quantum advantage? for what? when?
- > **future:** fault tolerant quantum computing




Thank you!

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