* the calculation of disconnected diagrams requires the evaluation of fermion bubbles

$$= B_{\Gamma}(t) = \sum_{\vec{x}} \Gamma_{\alpha\beta} S_{\beta\alpha}^{aa}(x,x)$$
 $\Gamma = \text{Dirac matrix}$

* all-to-all propagator: $\sum_{z} D_{\alpha\beta}^{ab}(x,z) S_{\beta\gamma}^{bc}(z,y) = \delta_{\alpha\gamma} \delta_{ac} \delta_{x,y}$

* stochastic method: $S(x,y) = \lim_{N_s \to \infty} \frac{1}{N_s} \sum_{r=1}^{N_s} \Phi^r(x) \left[\eta^r(y) \right]^{\dagger}$ (drop color and spin labels)

$$\eta^{r} = \text{stochastic source } (Z_{2} - \text{noise})$$

$$\begin{cases}
\left[\eta^{r}(x)\right]^{\dagger} \eta^{r}(x) = 1 \\
\lim_{N_{s} \to \infty} \frac{1}{N_{s}} \sum_{r=1}^{N_{s}} \left[\eta^{r}(x)\right]^{\dagger} \eta^{r}(y) = \delta_{x,y}
\end{cases}$$

$$\sum_{z} D(x,z) \Phi^{r}(z) = \eta^{r}(x)$$

$$B_{\Gamma}(t) = \lim_{N_s \to \infty} \frac{1}{N_s} \sum_{r=1}^{N_s} \sum_{\bar{x}} \left[\eta^r(x) \right]^{\dagger} \Gamma \Phi^r(x)$$

$$\frac{1}{N_s} \sum_{r=1}^{N_s} \sum_{\vec{x}} \left[\boldsymbol{\eta}^r(x) \right]^{\dagger} \Gamma \boldsymbol{\Phi}^r(x) = \frac{1}{N_s} \sum_{r=1}^{N_s} \sum_{\vec{x}} \left[\boldsymbol{\eta}^r(x) \right]^{\dagger} \Gamma \boldsymbol{S}(x,y) \boldsymbol{\eta}^r(y)$$

$$\frac{1}{N_s} \sum_{r=1}^{N_s} \left[\boldsymbol{\eta}^r(x) \right]^{\dagger} \boldsymbol{\eta}^r(y) \neq \delta_{x,y}$$

$$B_{\Gamma}(t)$$
 + gauge-variant noise (non-local bubble)

* recently the TrinLat collaboration has proposed a new technique: the dilution method (hep-lat/0505023)

- in few words: exact treatment of the all-to-all problem in a subset of variables only

* time dilution: L_t stochastic sources defined for each time slice

$$\eta^{r}(x) = \eta^{r}(\vec{x}) \delta_{t,t_{r}} \qquad r = 1, 2, ... L_{t}$$

* color-spin dilution: 12 stochastic sources defined for each color-spin combination

$$\left[\eta_{\alpha}^{a}(x)\right]^{r} = \eta^{r}(x)\,\delta_{\alpha,\alpha_{r}}\delta_{a,a_{r}} \qquad \alpha_{r} = 1,2,3,4, \ \alpha_{r} = 1,2,3$$

* we want to compare **three methods:** <u>time dilution</u>, <u>color-spin dilution</u> and no dilution (<u>direct method</u>)

* quenched gauge configurations at V T = 16^3 32 and β = 6.0 twisted-mass Wilson fermions at κ_c = 0.157409 (χ LF collaboration)

* disconnected diagrams for π^0 and η ' correlators:

$$\pi^{0}: \text{ connected} + \left\langle \sum_{t} \left\{ B_{\gamma_{5}}^{u}(t) B_{\gamma_{5}}^{u}(t+\delta t) + B_{\gamma_{5}}^{d}(t) B_{\gamma_{5}}^{d}(t+\delta t) - B_{\gamma_{5}}^{u}(t) B_{\gamma_{5}}^{d}(t+\delta t) - B_{\gamma_{5}}^{d}(t) B_{\gamma_{5}}^{u}(t+\delta t) \right\} \right\rangle$$

$$\eta': \text{ connected} + \left\langle \sum_{t} \left\{ B_{\gamma_{5}}^{u}(t) B_{\gamma_{5}}^{u}(t+\delta t) + B_{\gamma_{5}}^{d}(t) B_{\gamma_{5}}^{d}(t+\delta t) + B_{\gamma_{5}}^{d}(t) B_{\gamma_{5}}^{d}(t+\delta t) + B_{\gamma_{5}}^{d}(t) B_{\gamma_{5}}^{u}(t+\delta t) \right\} \right\rangle$$

* in tmLQCD,
$$\pi^0$$
 mixes with identity: $B_{\gamma_5}^{\tau_3}(t) = \hat{B}_{\gamma_5}^{\tau_3}(t) + (\pm)^{\tau_3} \frac{1}{a^3} \rho_{\gamma_5}(\mu)$ subtraction of $\sum_{t} \langle B_{\gamma_5}^{\tau_3}(t) \rangle \langle B_{\gamma_5}^{\tau_3'}(t + \delta t) \rangle$

* η ' is free from such a mixing

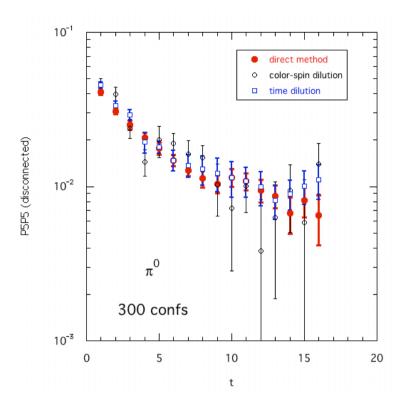
computational costs (per gauge conf.)

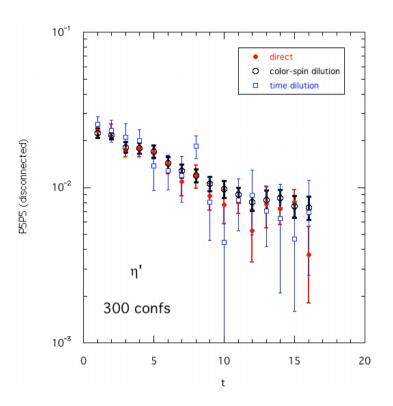
2*2*32 = 128 inversions (time dilution)

2*2*12 = 48 inversions (direct method)

2*12 = 24 inversions (color-spin dilution)

a
$$\mu = 0.0038$$
 [a $M_{\pi +} = 0.131$ (2)]





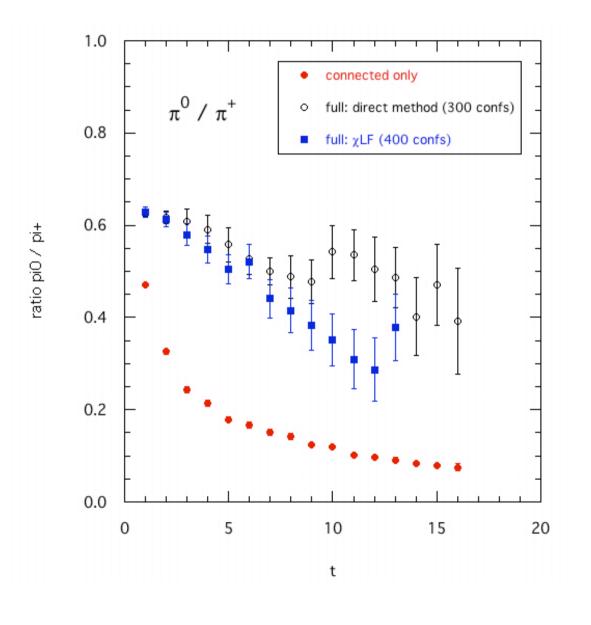
best: direct method

worst: color-spin dilution

best: color-spin dilution

worst: time dilution

next method: a hybrid one composed by spin dilution + direct method



χLF ('05): direct method + noise variance reduction trick

$$a * m(\pi^0) = 0.18 \pm 0.03$$
 (direct)
 $a * m(\pi^0) = 0.19 \pm 0.02$ (χ LF)

$$a * \Delta m (\pi^0 - \pi^+) = 0.05 \pm 0.03$$
 (direct)
 $a * \Delta m (\pi^0 - \pi^+) = 0.07 \pm 0.02$ (χ LF)