



Analytic Integration Methods in QFT: an Introduction

KMPB-Ukraine Workshop, Berlin & Kiev, March 4, 2025

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DESY AND TU DORTMUND

Recent survey books:

- Computer Algebra in Quantum Field Theory: Integration, Summation and Special Functions, Eds. JB, C. Schneider, Springer, 2013
- Elliptic Integrals, Elliptic Functions, and Modular Forms in Quantum Field Theory, Eds. JB, P. Paule, C. Schneider, 2019
- Anti-Differentiation and the Calculation of Feynman Amplitudes, Eds. JB, C. Schneider, Springer, 2021

Outline



- Introduction
- Principle Computation Steps
- Important Function Spaces
- The Different Calculation Methods
- Conclusions
- Backup: Aspects of elliptic solutions

Loops and Legs:

Feynman diagrams describe elementary scattering processes between bosons and fermions in Quantum Field Theory (QFT). Here we will thoroughly refer to renormalizable QFTs or effective field theories.

Where are these techniques important?

- 1 Perturbation Theory of the Standard Model and its renormalizable extensions.
- 2 String amplitude calculations
- 3 Perturbative calculations in Gravity
- 4 non-relativistic field theories in vacuum and at finite temperature and/or density

We will calculate **Feynman diagrams**. These are skeletons according to Feynman rules, connecting vertices with propagators.

They possess external lines: The Legs.

They possess internal closed lines: The Loops.

The machine, for which we perform the calculations:



LHC, Geneva/CH

Why are these calculations important ?

- 1 Precision extraction of coupling constants: $\alpha_s(M_Z)$ @1%
- 2 Do couplings unite at high scales and in which field theories?
- 3 Precision measurements of m_c, m_b, m_t at LHC and a future ILC
- 4 Precision understanding of the **Higgs and top sector** (at the LHC, ILC and possibly other machines)
- 5 Unravel the mathematical structure of microscopic processes analytically: **get further with the Stueckelberg-Feynman programme** as far as you can.

⇒ Genetic Code of the Micro Cosmos

The calculation methods



- 1 Generalized Hypergeometric Functions and Extensions \implies P. Paule, M. Kalmykov
- 2 Mellin-Barnes techniques
- 3 PSLQ: zero-dimensional integrals \implies D. Broadhurst
- 4 Guessing: one-dimensional integrals \implies this talk
- 5 Hyperlogarithms \implies E. Panzer
- 6 Solution of master-integrals using difference and differential equations \implies C. Schneider, J.-A. Weil, J. Henn
- 7 Risch algorithms \implies C. Raab
- 8 Summation techniques: construction of difference rings and fields \implies C. Schneider
- 9 Holonomic Integration \implies C. Koutschan
- 10 (multivalued) Almquist-Zeilberger algorithm \implies J. Ablinger
- 11 Expansion by regions \implies V. Smirnov
- 12 Elliptic integrals and related topics \implies S. Weinzierl, J. Brödel
- 13 Cutting techniques \implies D. Kreimer

Nested Sums and Iterative Integrals Function Spaces



Nested sums:

$$S_{1\dots m}(N) = \sum_{k_1=1}^N s(k_1) \sum_{k_2=1}^{k_1} s(k_2) \dots \sum_{k_m=1}^{k_{m-1}} s(k_m)$$

Iterated integrals:

$$F_{1\dots l}(x) = \int_0^x dy_1 f_1(y_1) \int_0^{y_1} dy_2 f_2(y_2) \dots \int_0^{y_{l-1}} dy_l f_l(y_l)$$

Mellin transform:

$$\sum c_\alpha S_\alpha(N) = \int_0^1 dx x^{N-1} F(x)$$
$$F_1(x) \cdot F_2(x) \stackrel{\alpha}{=} (F_1 F_2)(x) = F_{1,2}(x) + F_{2,1}(x) \quad \text{Shuffling}$$

$$S_k(N) \cdot S_l(N) = S_{k,l}(N) + S_{l,k}(N) - S_{k \wedge l}(N), \quad k \wedge l = \text{sign}(k)\text{sign}(l)(|k| + |l|), \quad k, l \in \mathbb{Z} \setminus \{0\} \quad \text{Stuffing}$$

For further details, see: J. Ablinger, JB, C. Schneider, 1304.7071, 1310.5645, JB, C. Schneider, 1809.02889

Function Spaces



Sums

Harmonic Sums

$$\sum_{k=1}^N \frac{1}{k} \sum_{l=1}^k \frac{(-1)^l}{l^3}$$

gen. Harmonic Sums

$$\sum_{k=1}^N \frac{(1/2)^k}{k} \sum_{l=1}^k \frac{(-1)^l}{l^3}$$

Cycl. Harmonic Sums

$$\sum_{k=1}^N \frac{1}{(2k+1)} \sum_{l=1}^k \frac{(-1)^l}{l^3}$$

Binomial Sums

$$\sum_{k=1}^N \frac{1}{k^2} \binom{2k}{k} (-1)^k$$

Integrals

Harmonic Polylogarithms

$$\int_0^x \frac{dy}{y} \int_0^y \frac{dz}{1+z}$$

gen. Harmonic Polylogarithms

$$\int_0^x \frac{dy}{y} \int_0^y \frac{dz}{z-3}$$

Cycl. Harmonic Polylogarithms

$$\int_0^x \frac{dy}{1+y^2} \int_0^y \frac{dz}{1-z+z^2}$$

root-valued iterated integrals

$$\int_0^x \frac{dy}{y} \int_0^y \frac{dz}{z\sqrt{1+z}}$$

iterated integrals containing elliptic structures

$$\int_0^z dx \frac{\ln(x)}{1+x} {}_2F_1 \left[\frac{4}{3}, \frac{5}{3}; \frac{x^2(x^2-9)^2}{(x^2+3)^3} \right]$$

Special Numbers

multiple zeta values

$$\int_0^1 dx \frac{\text{Li}_3(x)}{1+x} = -2\text{Li}_4(1/2) + \dots$$

gen. multiple zeta values

$$\int_0^1 dx \frac{\ln(x+2)}{x-3/2} = \text{Li}_2(1/3) + \dots$$

cycl. multiple zeta values

$$\mathbf{c} = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^2}$$

associated numbers

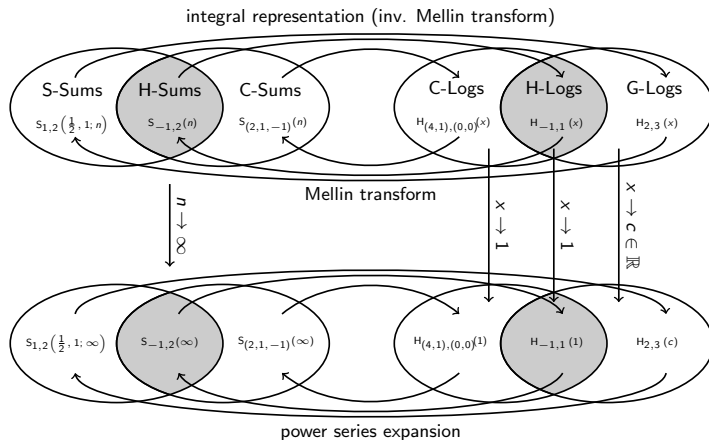
$$H_{8,w_3} = 2\text{arccot}(\sqrt{7})^2$$

associated numbers

$$\int_0^1 dx {}_2F_1 \left[\frac{4}{3}, \frac{5}{3}; \frac{x^2(x^2-9)^2}{(x^2+3)^3} \right]$$

shuffle, stuffle, and various structural relations \implies algebras

Function Spaces



square-root valued letters \iff nested binomial sums $\binom{2i}{i}$

non-iterative iterative integrals \implies Iterate on ${}_2F_1$'s: rat. argument
 (special cases: complete elliptic integrals) [arXiv:1706.01299]



Spill-Off: New Function Classes and Algebras



- **1998:** Harmonic Sums [Vermaseren; JB]
- **1999:** Harmonic Polylogarithms [Remiddi, Vermaseren]
- **2000, 2003, 2009:** Analytic continuation of harmonic sums, systematic algebraic reduction; structural relations [JB]
- **2001:** Generalized Harmonic Sums [Moch, Uwer, Weinzierl]
- **2004:** Infinite harmonic (inverse) binomial sums [Davydychev, Kalmykov; Weinzierl]
- **2011:** (generalized) Cyclotomic Harmonic Sums, polylogarithms and numbers [Ablinger, JB, Schneider]
- **2013:** Systematic Theory of Generalized Harmonic Sums, polylogarithms and numbers [Ablinger, JB, Schneider]
- **2014:** Finite nested Generalized Cyclotomic Harmonic Sums with (inverse) Binomial Weights [Ablinger, JB, Raab, Schneider]
- **2014-:** Elliptic integrals with (involved) rational arguments.

Particle Physics Generates **NEW** Mathematics.

Classes of generalized hypergeometric functions



At lower number of legs and/or loops Feynman integrals happen to be represented by these functions. After suitable mappings these functions have compact representations in infinite (multiple) absolutely convergent sums.

This allows for the Laurent-expansion in ϵ under the summation operator.

Important Examples:

- 1 $B(a, b)$
- 2 ${}_pF_q(a_i; b_j; x)$; always single sums
- 3 Appell functions; double sums
- 4 Kampé de Fériet functions, Horn functions and higher; more sums

The sums may be expanded and summed using algorithms like nestedsums, xsummer, HarmonicSums, Sigma, EvaluateMultiSums

Classes of generalized hypergeometric functions



Example:

Integrals of the following type emerge:

$$\begin{aligned}I_1(z) &= \int_0^1 dy y^\delta (1-y)^\eta \int_0^1 dx x^{\beta-1} (1-x)^{\gamma-\beta-1} (1-xyz)^{-\alpha} \\ &= B(\beta, \gamma - \beta) \int_0^1 dy y^\delta (1-y)^\eta {}_2F_1(\alpha, \beta; \gamma; yz) \\ &= B(\beta, \gamma - \beta) B(\delta, \eta - \delta) {}_3F_2(\delta, \alpha, \beta; \eta, \gamma; z)\end{aligned}$$

All ${}_pF_q$'s have single series representations. One series counts as one integral.

$${}_pF_q(a_1, \dots, a_p; b_1, \dots, b_q; z) = \sum_{k=0}^{\infty} \frac{(a_1)_k \dots (a_p)_k}{(b_1)_k \dots (b_q)_k} \frac{z^k}{k!}$$

Representations in terms of Mellin-Barnes Integrals



Used to resolve sums in denominators. In a way a counterpart to the binomial expansion of numerators.

$$\frac{1}{(A+B)^c} = \oint_{-i\infty+\gamma}^{i\infty+\gamma} d\sigma \frac{\Gamma(-\sigma)\Gamma(\sigma+c)}{\Gamma(c)} A^\sigma B^{-\sigma+c}$$

The contour integral covers the residues and closes either left or right.

One may use Beta-functions and (generalized) hypergeometric functions to perform the integrals further.

Important: One has to be able to undo the Mellin-Barnes integral(s).

- 1 Barnes Lemmas; map to Γ -functions, (in the most simple cases).
- 2 Residue theorem \implies leads to nested sums \implies use SIGMA.

The PSLQ Method



Seek an **Integer Relation** over a basis of **special numbers** out of a special class.

Example:

$$I = \int_0^1 dx \frac{\text{Li}_3(x)}{1+x}$$

The integral is of “**transcendentality**” $\tau = 4$.

The expected HPL(1) basis is spanned by:

$\ln^4(2)$, $\ln(2)\zeta_3$, $\ln^2(2)\zeta_2$, ζ_2^2 , $\text{Li}_4(1/2)$.

Calculate this integral numerically to high number of digits, e.g. 40 digits.

$$I \approx 0.3395454690873598695906678484608602061388$$

The PSLQ algorithm yields:

$$I = -\frac{1}{12} \ln^4(2) + \frac{\pi^4}{60} + \frac{3}{4} \ln(2)\zeta_3 + \frac{1}{12} \ln^2(2)\pi^2 - 2\text{Li}_4\left(\frac{1}{2}\right)$$
$$\zeta_{2k} = (-1)^{k-1} \frac{(2\pi)^{2k} B_{2k}}{2(2k)!}; \quad B_n \text{ [Bernoulli number]}$$



The method can be applied to all number classes, for which the bases are known.

- 1 multiple zeta values
- 2 classes of generalized harmonic numbers
- 3 cyclotomic multiple zeta values
- 4 elliptic numbers
- 5 etc. etc.

In various cases one can work with successive set extensions.

The method is important in determining the value of master integrals in the zero-scale case. This e.g. applies also to the master integrals appearing in the post-Newtonian expansion.

[Foffa et al. 1612.00482, 1902.10571, Blümlein et al. 1902.11180].

Similarly, constants of that kind form also initial values for differential and difference equations for single variate integrals.

Guessing Difference Equations



It is often easier to calculate Mellin moments for a quantity for fixed values of N than to derive the relation for general values of N in the first place. If the quantity under consideration is known to be **recurrent** than its difference equation is of finite order and degree.

$$\exists \sum_{k=0}^O P_k^{(l)}(N) F(k+N) = 0; \quad \max\{l\} - \text{degree}; \quad O - \text{order}$$

Example:

$$-(N+1)^3 F(N) - (3N^2 - 9N - 7)F(N+1) + (N+2)^3 F(N+2) = 0$$
$$F(1) = 1; \quad F(2) = \frac{1}{8}$$

Solution:
$$F(N) = \sum_{k=1}^N \frac{1}{k^3} = S_3(N)$$

Guessing Difference Equations



Solution of large problems

Assume you would like to calculate the massless 3-loop Wilson coefficients in deep-inelastic scattering using this method. How many moments would you need and how do they look like ?

About 5200 moments are needed. The largest ones are ratios of #13000/#13000 digits. They can be calculated within 15 min.

After 3 weeks you will find a difference equation of degree ~ 1000 and order 35, if you have a reasonable computer (100 Gbyte RAM).

After another week you have the solution as function of N .

Problem: It is sophisticated to obtain the input a priori. Solved: JB & C. Schneider, 2017.

Recent result: a 3-loop anomalous dimension computed from scratch.

[arXiv:1701.04614, 1705.01508, 1908.00357]: For contemporary non-trivial cases we have calculated 12000 moments.

The above numbers apply to applications in 2009. Now Sage-Implementations allow solutions within hours.

The IBPs deliver a vast amount of differential equations forming systems, which are nested hierarchically.

Provide boundary conditions [usually using other methods]

Perform uncoupling of these systems

- In the case of complete 1st order uncoupling: \exists complete solution algorithm in the case of any basis choice for 1 parameter systems

All solutions are iterative integrals over whatsoever alphabet: $\int_0^x dy f_a(y) H_b(y)$

- Irreducible n th order systems ($n \geq 2$): present target of research even in mathematics; good prospects in the case of 2nd order systems [elliptic and ${}_2F_1$ solutions]

At least one function is given by a definite integral, others iterate on.

\implies iterated integral algebras \implies bases

Risch algorithms



Risch-algorithm: [Integrability of elementary functions in terms of elementary functions.](#)

Analyse whether or not certain indefinite integrals appearing in QFT calculations are elements of **higher pre-defined function spaces**.

This can be decided by solving associated differential equation systems.

Iterative integrals over certain alphabets belong to this class and the associated alphabets are even coined in this way.

$$H_{b,\bar{a}}(x) = \int_0^x dy f_b(y) H_{\bar{a}}(y), H_{\emptyset} = 1$$

$$f_c(x) \in \{f_1(x), \dots, f_m(x)\}$$

The algorithm tries to transform the corresponding integrals into iterative integrals. This also allows to use algebraic shuffle relations.



Summation Techniques

The integrals can usually be traded for a **lower number of sums** (finite or infinite).

Solve these sums for N and/or in terms of special constants.

Principle Idea:

- 1 Sums may be represented in vector spaces, algebras, and finally fields/rings
- 2 Rephrase the sums in the setting of difference fields and rings
- 3 Apply telescoping, creative telescoping, and other principles in this setting to compute recurrences
- 4 Try to solve the recurrences; possible for most sums occurring from Feynman integrals
- 5 In addition, use nested sums algebras to speed up calculations

Telescoping: Find a function $g(k)$ such

$$f(k) = g(k+1) - g(k)$$
$$F(N) = \sum_{k=1}^N f(k) = g(N+1) - g(1)$$

\implies **nested sums algebras** \implies **bases**

Sigma solves large scale problems running over months and using several hundred Gb RAM.

Almkvist-Zeilberger Algorithm



- Given a multiple integral over hyperexponential terms:

$$F(n) = \int_0^1 dx_1 \dots dx_j \prod_{k=1}^l (P(x_i, n))^{r_k, \epsilon}, \quad r_k \in \mathbb{R} \text{ and } n \in \mathbb{N} \text{ a parameter.}$$

- Find a recurrence:

$$\sum_{k=0}^m p_k(n, \epsilon) F(n+k) = H(n, \epsilon) \text{ with some inhomogeneity } H(n, \epsilon).$$

- Correspondingly $n \rightarrow x$, a differential equation:

$$\sum_{k=0}^m p_k(x, \epsilon) \frac{d^k}{dx^k} F(x) = K(x, \epsilon) \text{ with some inhomogeneity } K(x, \epsilon).$$

Either the inhomogeneities can be forced to vanish, or a hierarchy of equations has to be solved using summation techniques and DEQ-solvers (which may also be summation techniques).

The AZ-Theorem states, that for recurrent quantities in a variable N one can always find a recurrence successively. It may have of course a very high order and/or degree. This applies also to multi-integrals depending on N . Conversely, one can consider the associated differential equation to $x \in \mathbb{C}$, e.g.

What can be achieved by all that?



Quantum Field Theory:

- A lot of **integration technology** has been created for many analytic precision calculations for the **Large Hadron Collider** and the planned International Linear Collider.
- The results allow also for many advanced solutions in **combinatorics and number theory**.
- Within elementary particle physics the present results allow to improve the **precision of two fundamental constants** of the Standard Model:

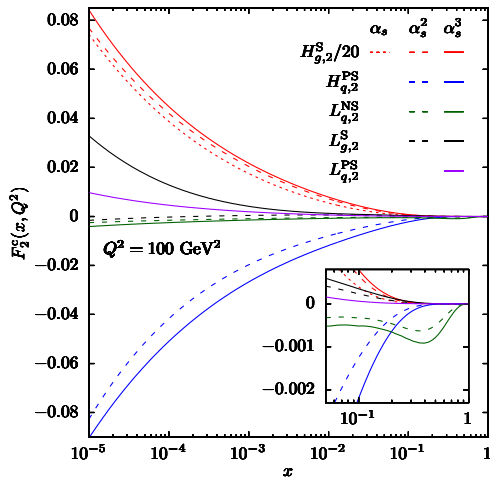
$$\frac{\delta\alpha_s(M_Z)}{\alpha_s(M_Z)} < 1\% \quad \delta m_c < 20\text{MeV}[1.6\%]$$

which may have consequences for various proposed extensions **of the SM**.

Gravity:

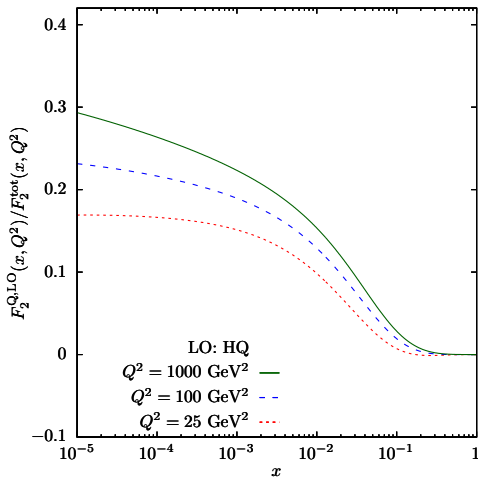
Many of the above techniques apply also in the effective field theory approach to gravity and proof to be very efficient there, in particular one wants to address even higher orders.

Single-mass contributions to $F_2^{c,b}$



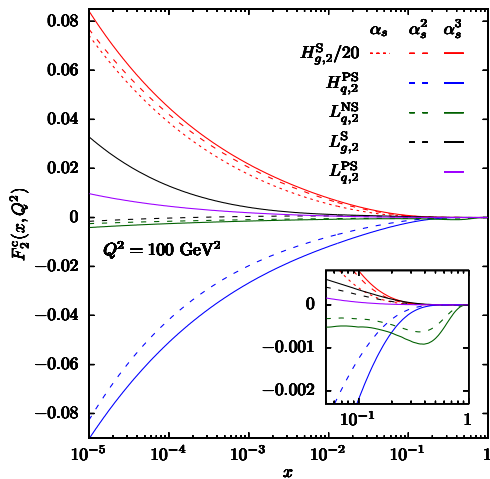
charm for $Q^2 = 100 \text{ GeV}^2$.

Allows to strongly reduce the current theory error on m_c .



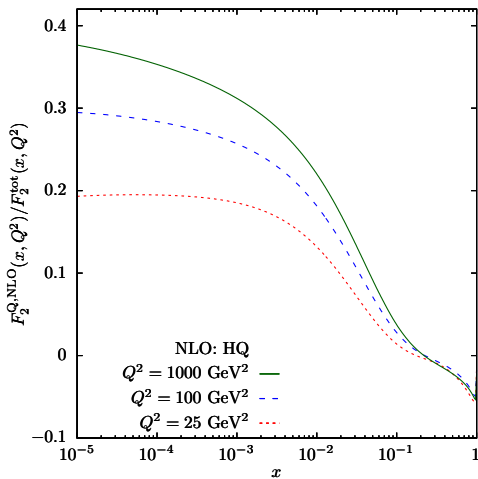
c and b single mass contributions

Single-mass contributions to $F_2^{c,b}$



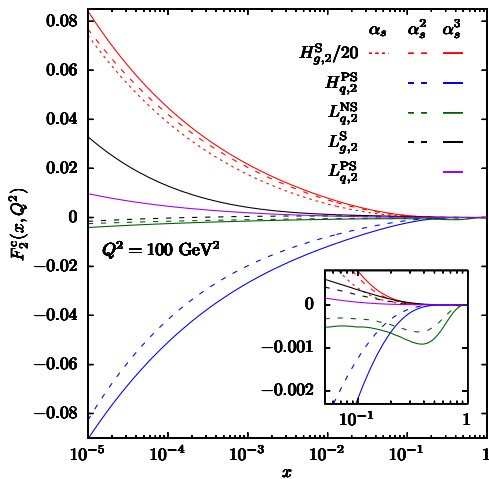
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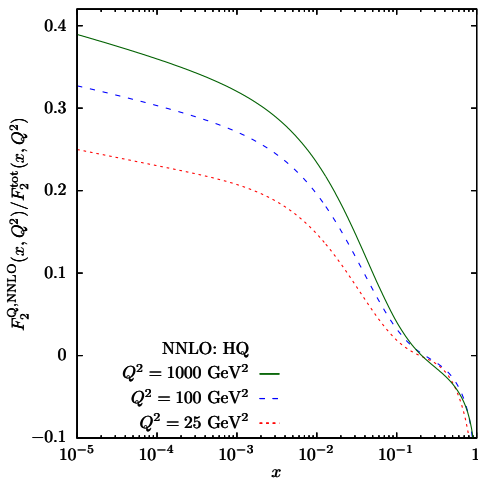
c and b single mass contributions

Single-mass contributions to $F_2^{c,b}$



charm for $Q^2 = 100 \text{ GeV}^2$.

Allows to strongly reduce the current theory error on m_c .



c and b single mass contributions

- 1 Current problems at the precision frontier of Quantum Field Theory request advanced integration (anti-differentiation) methods.
- 2 There are standardized chains down to the master integrals, with excellent computer algebraic implementations.
- 3 The calculation of the remainder integrals forms the current challenge.
- 4 The first order factorizable cases in the case of nested sums or difference equations are completely understood and there are good implementations to solve these cases.
- 5 There are techniques for 2nd order factorization cases, still in course of development.
- 6 More involved structures require a lot of further efforts.
- 7 Massive problems (from 1-scale also in gravity) map to more involved function spaces.
- 8 The present techniques form one important ingredient to proceed to higher post-Newtonian and post-Minkowskian orders using effective field theory methods.
- 9 Many of these technologies were created during our 3-loop QCD heavy flavor project for deep-inelastic scattering (unpolarized and polarized), which has now finished after 15 years.

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Backup Slides



The Individual Steps: from IBPs to Closed Form q -Series



- Generate the master integrals, determine their hierarchy, and look whether you have only 1st order factorization or also **2nd order terms**
- The latter can be trivial in case; check whether they persist in **Mellin space**
- If yes, analyze the **2nd order differential equation**
- One usually finds a ${}_2F_1$ -**solution** with rational argument $r(z)$, where $r(z)$ has additional singularities, i.e. the problem is of 2nd order, but has **more than 3 singularities**.
- Triangle group relations may be used to map the ${}_2F_1$ depending on the rational parameters a, b, c to the complete elliptic integrals **or not**.
- In the latter case return to the formalism on slide **21** and stop.
- If yes, one may walk along the **q -series avenue**.
- **Different Levels of Complexity:**
 - 1st order factorization in Mellin space:

$$\mathbf{M}[\mathbf{K}(1-z)](N) = \frac{2^{4N+1}}{(1+2N)^2 \binom{2N}{N}^2}; \quad \mathbf{M}[\mathbf{E}(1-z)](N) = \frac{2^{4N+2}}{(1+2N)^2 (3+2N) \binom{2N}{N}^2}$$

- Criteria by Herfurtnner (1991), Movasati et al. (2009) are obeyed.
 ⇒ 2-loop sunrise and kite diagrams, cf. Weinzierl et al. 2014-17.
 Only $\mathbf{K}(r(z))$ and $\mathbf{K}'(r(z))$ contribute as elliptic integrals.
- Also $\mathbf{E}(r(z))$ and $\mathbf{E}'(r(z))$, square roots of quadratic forms etc. contribute (present case)
- Transform now: $x \rightarrow q$.
- The kinematic variable x :

$$k^2 = \frac{-x^3}{(1+x)^3(1-3x)} = \frac{\vartheta_2^4(q)}{\vartheta_3^4(q)}$$

$$x = \frac{\vartheta_2^2(q)}{3\vartheta_2^2(q^3)}, \text{ i.e. } x \in [1, +\infty[$$

by a **cubic** transformation (Legendre-Jacobi).

[see also Borwein, Borwein: AGM; and Broadhurst (2008).]

$$x = \frac{1}{3} \frac{\eta^2(2\tau)\eta^2(3\tau)}{\eta^2(\tau)\eta^4(6\tau)}, \text{ singular, } \propto \frac{1}{q}$$

The Individual Steps: from IBPs to Closed Form q -Series



- Map to a Modular Form, which can be represented by Lambert Series

- How to find the η -ratio ? \implies Many are listed as sequences in Sloan's OEIS.
- To find a modular form, situated in a corresponding finite-dimensional vector space M_k one has to meet a series of conditions and usually split off a factor $1/\eta^k(\tau)$, $k > 0$.
- The remainder modular form is now a polynomial over \mathbb{Q} of Lambert-Eisenstein series

$$\sum_{n=0}^{\infty} \frac{m^n q^{an+b}}{1 - q^{an+b}} \cdot$$

Example:

$$\mathbf{K}(z(x)) = \frac{\pi}{2} \sum_{k=1}^{\infty} \frac{q^k}{1 + q^{2k}}$$

- In this case, two q series are equal, if both are modular forms, and agree in a series of k first terms, where k is predicted for each congruence sub-group of $\Gamma(N)$.

The Individual Steps: from IBPs to Closed Form q -Series



- Map Lambert-Eisenstein Series into the frame of Elliptic Polylogarithms
- Examples:

$$\begin{aligned} \mathbf{K}(z) &= \frac{\pi}{2} \sum_{k=1}^{\infty} \frac{q^k}{1+q^{2k}} = \frac{\pi}{i} \sum_{k=1}^{\infty} [\text{Li}_0(iq^k) - \text{Li}_0(-iq^k)] \\ &= \frac{\pi}{4} \bar{E}_{0,0}(i, 1, q), \\ q \frac{\vartheta_4'(q)}{\vartheta_4(q)} &= -\frac{1}{2} [\text{ELi}_{-1;0}(1; 1; q) + \text{ELi}_{-1;0}(-1; 1; q)] \\ &\quad + [\text{ELi}_{0;0}(1; q^{-1}; q) + \text{ELi}_{0;0}(-1; q^{-1}; q)] \\ &\quad - [\text{ELi}_{-1;0}(1; q^{-1}; q) + \text{ELi}_{-1;0}(-1; q^{-1}; q)]. \end{aligned}$$

- New type of elliptic polylogarithm, e.g.: $\text{ELi}_{-1;0}(-1; q^{-1}; q)$, $y = y(q)$!
- Argument synchronization necessary: $-q \rightarrow q$, $q^k \rightarrow q$ (cyclotomic).

The Individual Steps: from IBPs to Closed Form q -Series



- Map Lambert-Eisenstein Series into the frame of Elliptic Polylogarithms
- Examples:

$$\begin{aligned} \mathbf{K}(z) &= \frac{\pi}{2} \sum_{k=1}^{\infty} \frac{q^k}{1+q^{2k}} = \frac{\pi}{i} \sum_{k=1}^{\infty} [\text{Li}_0(iq^k) - \text{Li}_0(-iq^k)] \\ &= \frac{\pi}{4} \bar{E}_{0,0}(i, 1, q), \\ q \frac{\vartheta_4'(q)}{\vartheta_4(q)} &= -\frac{1}{2} [\text{ELi}_{-1;0}(1; 1; q) + \text{ELi}_{-1;0}(-1; 1; q)] \\ &\quad + [\text{ELi}_{0;0}(1; q^{-1}; q) + \text{ELi}_{0;0}(-1; q^{-1}; q)] \\ &\quad - [\text{ELi}_{-1;0}(1; q^{-1}; q) + \text{ELi}_{-1;0}(-1; q^{-1}; q)]. \end{aligned}$$

- New type of elliptic polylogarithm, e.g.: $\text{ELi}_{-1;0}(-1; q^{-1}; q)$, $y = y(q)$!
- Argument synchronization necessary: $-q \rightarrow q$, $q^k \rightarrow q$ (cyclotomic).

Elliptic Solutions and Analytic q -Series



■ Terms to be translated:

- rational functions in x
- \mathbf{K}, \mathbf{E}
- $\sqrt{(1-3x)(1+x)}$
- $H_{\bar{a}}(x)$

Examples:

$$\mathbf{E}(k^2) = \mathbf{K}(k^2) + \frac{\pi^2 q}{\mathbf{K}(k^2)} \frac{d}{dq} \ln [\vartheta_4(q)]$$

$$\mathbf{E}'(k^2) = \frac{\pi}{2\mathbf{K}(k^2)} \left[1 + 2 \ln(q) q \frac{d}{dq} \ln [\vartheta_4(q)] \right].$$

$$\begin{aligned} \frac{1}{\mathbf{K}(k^2)} &= \frac{2}{\pi \eta^{12}(\tau)} \left\{ \frac{5}{48} \left\{ 1 - 24 \text{ELi}_{0,-1}(1; 1; q) - 4 \left[1 - \frac{3}{2} \left[\text{ELi}_{0,-1}(1; 1; q) + \text{ELi}_{0,-1}(1; i; q) \right. \right. \right. \right. \\ &\quad \left. \left. \left. + \text{ELi}_{0,-1}(1; -1; q) + \text{ELi}_{0,-1}(1; -i; q) \right] \right\} \left\{ -1 + 4 \left[-\frac{1}{2} \left[\text{ELi}_{-2,0}(i; 1/q; q) \right. \right. \right. \right. \\ &\quad \left. \left. \left. + \text{ELi}_{-2,0}(-i; 1/q; q) \right] + \left[\text{ELi}_{-1,0}(i; 1/q; q) + \text{ELi}_{-1,0}(-i; 1/q; q) \right] - \frac{1}{2} \left[\text{ELi}_{0,0}(i; 1/q; q) \right. \right. \right. \end{aligned}$$

$$\begin{aligned}
& + \text{ELi}_{0,0}(-i; 1/q; q) \Big] \Big\} - \frac{1}{16} \left\{ 5 + 4 \left[-\frac{1}{2} \left[\text{ELi}_{-4,0}(i; 1/q; q) + \text{ELi}_{-4,0}(-i; 1/q; q) \right] \right. \right. \\
& + 2 \left[\text{ELi}_{-3,0}(i; 1/q; q) + \text{ELi}_{-3,0}(-i; 1/q; q) \right] - 3 \left[\text{ELi}_{-2,0}(i; 1/q; q) \right. \\
& + \text{ELi}_{-2,0}(-i; 1/q; q) \Big] + 2 \left[\text{ELi}_{-1,0}(i; 1/q; q) + \text{ELi}_{-1,0}(-i; 1/q; q) \right] \\
& \left. \left. - \frac{1}{2} \left[\text{ELi}_{0,0}(i; 1/q; q) + \text{ELi}_{0,0}(-i; 1/q; q) \right] \right\} \right\} \\
H_{-1}(x) = \ln(1+x) &= -\ln(3q) - \bar{E}_{0,-1;2}(-1; -1; q) + \bar{E}_{0,-1;2}(\rho_6; -1; q) \\
& - \bar{E}_{0,-1;2}(\rho_3; -i; q) - \bar{E}_{0,-1;2}(\rho_3; i; q) \\
I(q) = \frac{H_1(x)}{\eta^k(\tau)} \cdot \mathbf{P} & \left[\ln(q), \text{Li}_0(q^m), \text{ELi}_{k,l}(x, y, q), \text{ELi}_{k',l'}(x, q^{-1}, q) \right] \\
& \int \frac{dq}{q} I(q)
\end{aligned}$$

is usually not an elliptic polylogarithm, due to the η -factor, but a higher transcendental function in q .
We are still in the unphysical region and have to map back to $x \in [0, 1]$.

$$\Delta\rho = \frac{3G_F m_t^2}{8\pi^2 \sqrt{2}} \left(\delta^{(0)}(x) + \frac{\alpha_s}{\pi} \delta^{(1)}(x) + \left(\frac{\alpha_s}{\pi} \right)^2 \delta^{(2)}(x) + \mathcal{O}(\alpha_s^3) \right)$$

$$\begin{aligned} \delta^{(2)}(x) = & \dots + C_F \left(C_F - \frac{C_A}{2} \right) \left[\frac{11-x^2}{12(1-x^2)^2} f_{8a}(x) + \frac{9-x^2}{3(1-x^2)^2} f_{9a}(x) + \frac{1}{12} f_{10a}(x) \right. \\ & \left. + \frac{5-39x^2}{36(1-x^2)^2} f_{8b}(x) + \frac{1-9x^2}{9(1-x^2)^2} f_{9b}(x) + \frac{x^2}{12} f_{10b}(x) \right] \\ & + \frac{C_F T_F}{9(1-x^2)^3} \left[(5x^4 - 28x^2 - 9) f_{8a}(x) + \frac{1-3x^2}{3x^2} (9x^4 + 9x^2 - 2) f_{8b}(x) \right. \\ & \left. + (9-x^2)(x^4 - 6x^2 - 3) f_{9a}(x) + \frac{1-9x^2}{3x^2} (3x^4 + 6x^2 - 1) f_{9b}(x) \right] \end{aligned}$$

Insert the respective functions, which are the 2nd order solutions and perform deep enough Taylor expansions in the respective regions in x .