

# Radiative Corrections to Production of Scalar and Vector Leptoquarks in $e^+e^-$ Annihilation

Hawaii, April 26 – 30 1993

J. Blümlein

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# 1 Classification of Leptoquarks

- B and L conserving
- family-diagonal
- $SU(3)_c \times SU(2)_L \times U(1)_Y$  invariant couplings

$$\mathcal{L} = \mathcal{L}_{|F|=2}^f + \mathcal{L}_{F=0}^f + \mathcal{L}^{\gamma, Z, g}$$

BUCHMÜLLER,  
RÜCKL, WYLER;  
Z.B., RÜCKL  
Phys. Lett. B, May 93.

$$\mathcal{L}^{\gamma, Z, g} = \sum_{scalars} [(D^\mu \Phi)^\dagger (D_\mu \Phi) - M^2 \Phi^\dagger \Phi] + \sum_{vectors} \left[ -\frac{1}{2} G_{\mu\nu}^\dagger G^{\mu\nu} + M^2 \Phi^{\mu\dagger} \Phi_\mu \right]$$

$$D_\mu = \partial_\mu - ieQ^\gamma A_\mu - ieQ^Z Z_\mu - ig_s \frac{\lambda_a}{2} \mathcal{A}_\mu^a$$

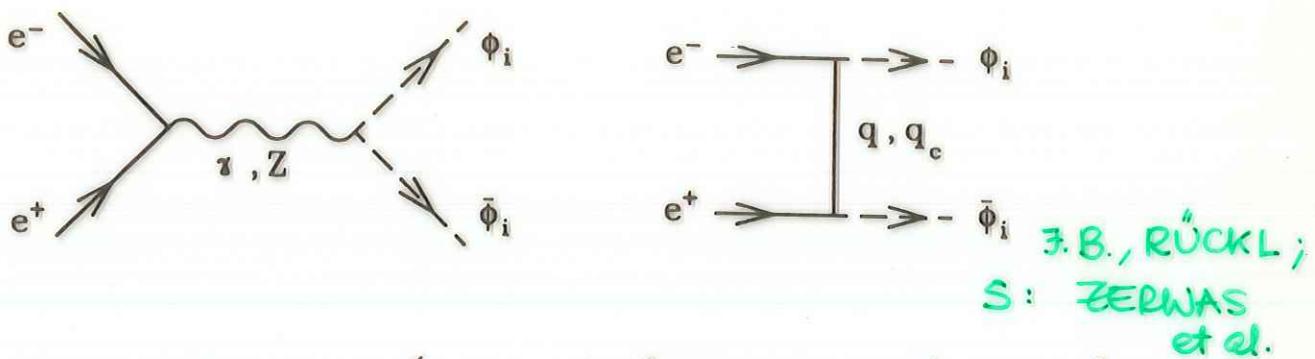
$$\begin{aligned} \mathcal{L}_{F=0}^f &= (h_{2L}\bar{u}_R l_L + h_{2R}\bar{q}_L i\tau_2 e_R)R_2 + \tilde{h}_{2L}\bar{d}_R l_L \tilde{R}_2 \\ &+ (h_{1L}\bar{q}_L \gamma^\mu l_L + h_{1R}\bar{d}_R \gamma^\mu e_R)U_{1\mu} \\ &+ \tilde{h}_{1R}\bar{u}_R \gamma^\mu e_R \tilde{U}_{1\mu} + h_{3L}\bar{q}_L \vec{\tau} \gamma^\mu l_L \tilde{U}_{3\mu} + h.c. \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{|F|=2}^f &= (g_{1L}\bar{q}_L^c i\tau_2 l_L + g_{1R}\bar{u}_R^c e_R)S_1 \\ &+ \tilde{g}_{1R}\bar{d}_R^c e_R \tilde{S}_1 + g_{3L}\bar{q}_L^c i\tau_2 \vec{\tau} l_L \tilde{S}_3 \\ &+ (g_{2L}\bar{d}_R^c \gamma^\mu l_L + g_{2R}\bar{q}_L^c \gamma^\mu e_R)V_{2\mu} \\ &+ \tilde{g}_{2L}\bar{u}_R^c \gamma^\mu l_L \tilde{V}_{2\mu} + h.c., \end{aligned}$$

leptoquark ( $\Phi$ )	spin	$F$	colour	$T_3$	$Q_{em}$	$\lambda_L(lq)$	$\lambda_R(lq)$	$\lambda_L(\nu q)$
$S_1$	0	-2	3	0	1/3	$g_{1L}$	$g_{1R}$	$-g_{1L}$
$S_1$	0	-2	3	0	4/3	0	$\tilde{g}_{1R}$	0
$\tilde{S}_3$	0	-2	3	+1	4/3	$-\sqrt{2}g_{3L}$	0	0
				0	1/3	$-g_{3L}$	0	$-g_{3L}$
				-1	-2/3	0	0	$\sqrt{2}g_{3L}$
$R_2$	0	0	3	1/2	5/3	$h_{2L}$	$h_{2R}$	0
$\tilde{R}_2$	0	0	3	-1/2	2/3	0	$-h_{2R}$	$h_{2L}$
				1/2	2/3	$\tilde{h}_{2L}$	0	0
				-1/2	-1/3	0	0	$\tilde{h}_{2L}$
$V_{2\mu}$	1	-2	3	1/2	4/3	$g_{2L}$	$g_{2R}$	0
				-1/2	1/3	0	$g_{2R}$	$g_{2L}$
$\tilde{V}_{2\mu}$	1	-2	3	1/2	1/3	$\tilde{g}_{2L}$	0	0
				-1/2	-2/3	0	0	$\tilde{g}_{2L}$
$U_{1\mu}$	1	0	3	0	2/3	$h_{1L}$	$h_{1R}$	$h_{1L}$
$\tilde{U}_{1\mu}$	1	0	3	0	5/3	0	$\tilde{h}_{1R}$	0
$\tilde{U}_{3\mu}$	1	0	3	+1	5/3	$\sqrt{2}h_{3L}$	0	0
				0	2/3	$-h_{3L}$	0	$h_{3L}$
				-1	-1/3	0	0	$\sqrt{2}h_{3L}$

18 States

## 2 The Born Cross Section



$$\frac{d\sigma_{scalar}}{d\cos\theta} = \frac{3\pi\alpha^2}{8s}\beta^3 \sin^2\theta \sum_{a=L,R} \left\{ |\kappa_a(s)|^2 + \left(\frac{\lambda_a}{e}\right)^2 \frac{4Re[\kappa_a(s)]}{t(\beta, \cos\theta)} + \left(\frac{\lambda_a}{e}\right)^4 \frac{4}{t^2(\beta, \cos\theta)} \right\}$$

$$\frac{d\sigma_{vector}}{d\cos\theta} = \frac{3\pi\alpha^2}{8M_\Phi^2}\beta \sum_{a=L,R} \left\{ |\kappa_a(s)|^2 \tilde{F}_1(\theta, \beta) + \left(\frac{\lambda_a}{e}\right)^2 Re[\kappa_a(s)] \tilde{F}_2(\theta, \beta) + \left(\frac{\lambda_a}{e}\right)^4 \tilde{F}_3(\theta, \beta) \right\}$$

$$\kappa_a(s) = \sum_{V=\gamma,Z} Q_a^V(e) \frac{s}{s - M_V^2 + iM_V\Gamma_V} Q^V(\Phi)$$

$$\tilde{F}_1(\theta, \beta) = \beta^2 \left[ 1 + \frac{1}{4}(1 - 3\beta^2) \sin^2\theta \right]$$

$$\tilde{F}_2(\theta, \beta) = 2 \left[ 1 - \frac{1 - \beta^2}{t(\beta, \cos\theta)} \right] (1 - \beta^2) + 4\beta^2 - \beta^2 \left[ 1 - 2 \frac{1 - \beta^2}{t(\beta, \cos\theta)} \right] \sin^2\theta$$

$$\tilde{F}_3(\theta, \beta) = 4 + \frac{\beta^2}{4} \left\{ (1 - \beta^2) \left[ \frac{4}{t(\beta, \cos\theta)} \right]^2 + \frac{s}{M_\Phi^2} \right\} \sin^2\theta$$

$$\begin{aligned}\sigma_{scalar}(s) &= \frac{\pi\alpha^2\beta^3}{2s} \sum_{a=L,R} \left\{ |\kappa_a(s)|^2 + \left(\frac{\lambda_a}{e}\right)^2 Re[\kappa_a(s)]F_1(\beta) + \left(\frac{\lambda_a}{e}\right)^4 F_2(\beta) \right\} \\ \sigma_{vector}(s) &= \frac{\pi\alpha^2\beta}{2M_\Phi^2} \sum_{a=L,R} \left\{ |\kappa_a(s)|^2 \tilde{F}_1(\beta) + \left(\frac{\lambda_a}{e}\right)^2 Re[\kappa_a(s)]\tilde{F}_2(\beta) + \left(\frac{\lambda_a}{e}\right)^4 \tilde{F}_3(\beta) \right\}\end{aligned}$$

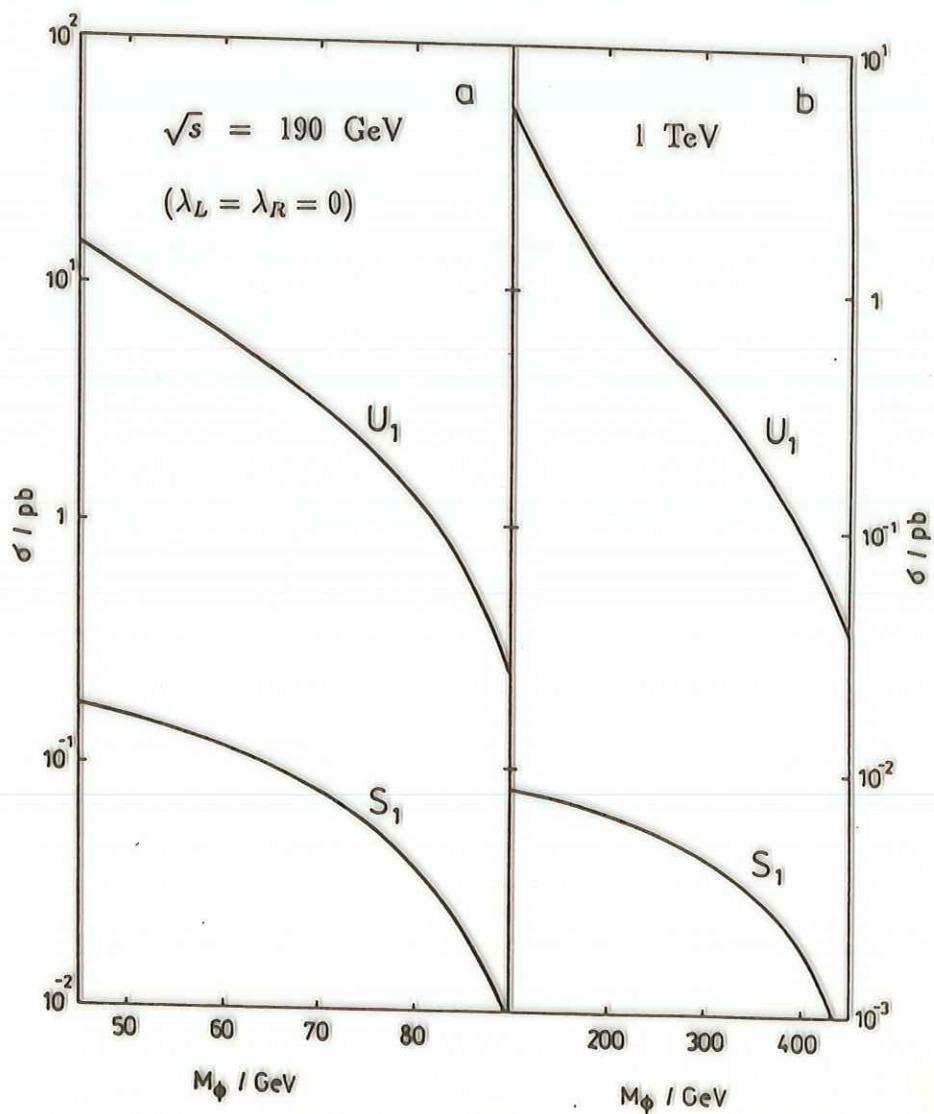
$$F_1(\beta) = \frac{3}{2} \left( \frac{1+\beta^2}{\beta^2} - \frac{(1-\beta^2)^2}{2\beta^3} \ln \frac{1+\beta}{1-\beta} \right)$$

$$F_2(\beta) = 3 \overline{\left( -\frac{1}{\beta^2} + \frac{1+\beta^2}{2\beta^3} \ln \frac{1+\beta}{1-\beta} \right)}$$

$$\tilde{F}_1(\beta) = \underline{\beta^2} \left( \frac{7-3\beta^2}{4} \right)$$

$$\tilde{F}_2(\beta) = \frac{15}{4} + 2\beta^2 - \frac{3}{4}\beta^4 - \frac{3}{8\beta}(1-\beta^2)^2(5-\beta^2) \ln \frac{1+\beta}{1-\beta}$$

$$\tilde{F}_3(\beta) = 3(1+\beta^2) + \frac{\beta^2}{4} \frac{s}{M_\Phi^2} + \frac{3}{2\beta}(1-\beta^4) \ln \frac{1+\beta}{1-\beta}.$$



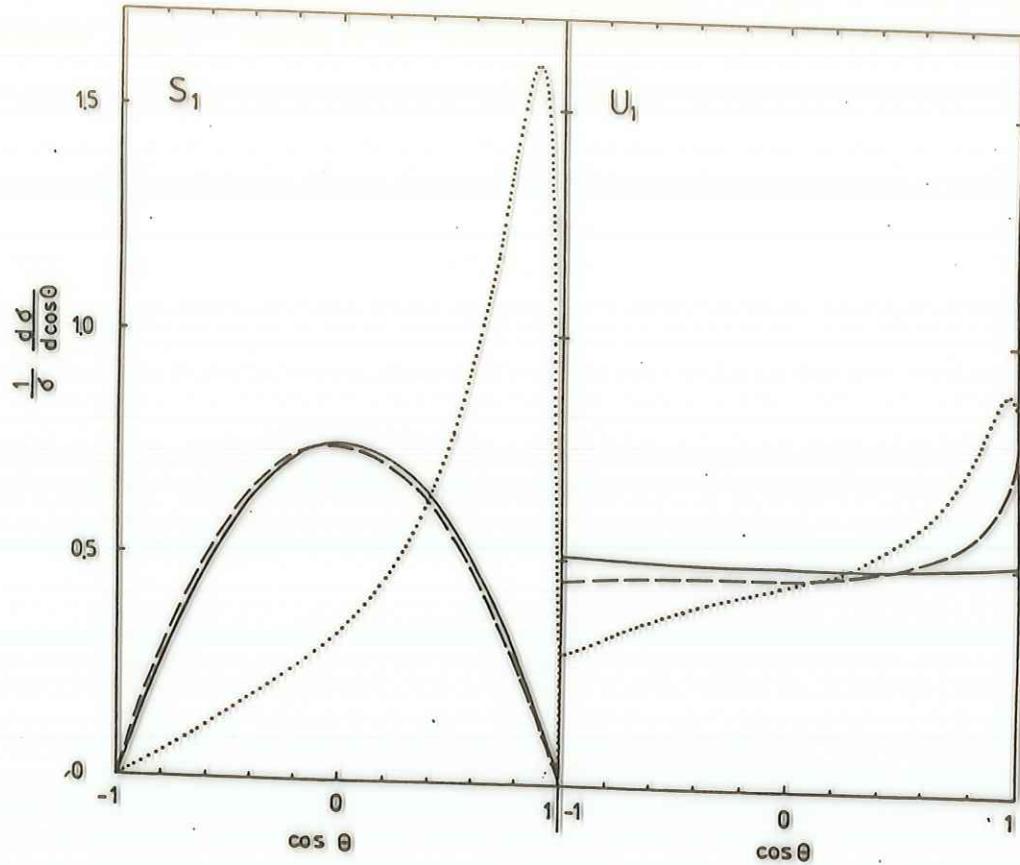


Figure 4: Angular distributions of the scalar leptoquark  $S_1$  and the vector leptoquark  $U_1$  at  $\sqrt{s} = 1$  TeV and for  $M_\Phi = 400$  GeV:  $\lambda_L = \lambda_R = 0$  (solid);  $\lambda_L/e = 0.3, \lambda_R = 0$  (dashed);  $\lambda_R/e = 1, \lambda_L = 0$  (dotted).

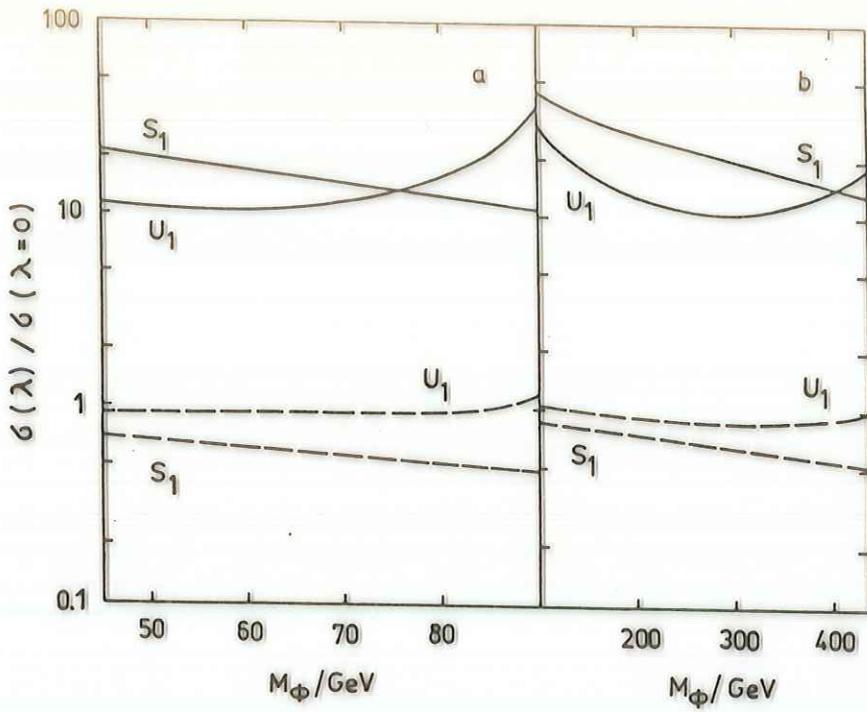


Figure 3: Effect of right-handed fermion couplings,  $\lambda_R/e = 0.3$  (dashed) and 1 (solid), on the production cross sections for scalar ( $S_1$ ) and vector ( $U_1$ ) leptoquarks at  $\sqrt{s} = 190$  GeV (a) and 1 TeV (b).

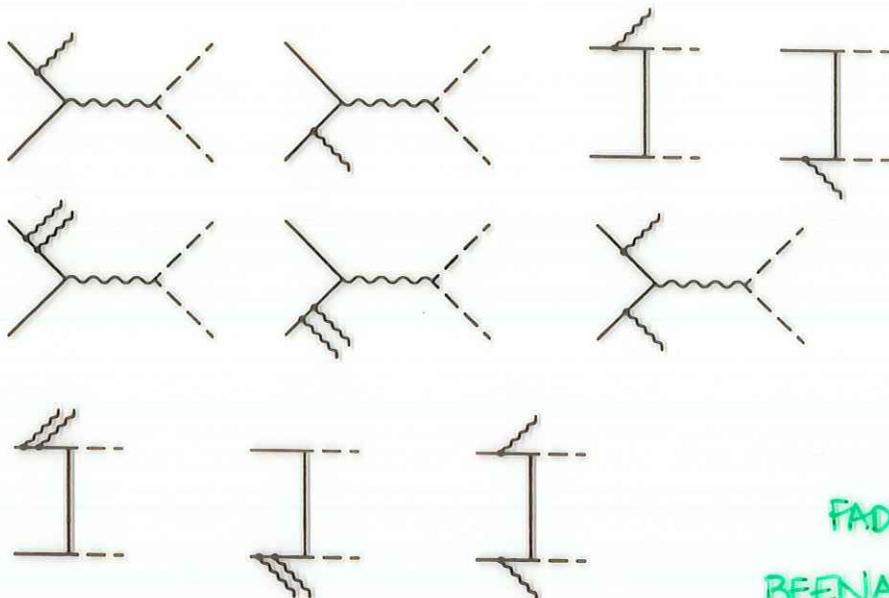
	scalars		vectors		Luminosity
	min	max	min	max	
$\sqrt{s} = 190$ GeV	$S_1, S_3^{1/3}$	$S_3^{4/3}$	$U_1, U_3^{2/3}$	$U_3^{5/3}$	$200 pb^{-1}$
$M_\Phi = 80$ GeV	0.040	1.18	1.366	12.08	
$\sqrt{s} = 1$ TeV	$S_1, S_3^{1/3}$	$S_3^{4/3}$	$U_1, U_3^{2/3}$	$U_3^{5/3}$	$10 fb^{-1}$
$M_\Phi = 430$ GeV	0.001	0.029	0.060	0.368	

Table : Maximum and minimum cross sections in  $pb$  assuming  $\lambda_{L,R} = 0$ . The cross sections for the remaining leptoquark states in table 1 lie in between these values. The masses correspond roughly to  $\beta = 0.5$ .

### 3 QED Corrections

#### 3.1 Initial State Radiation

i)



FADIN, KURAEV;

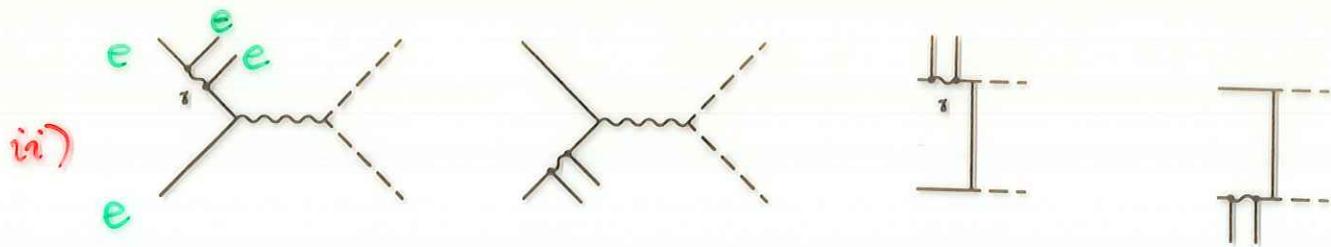
BEENAKKER,  
BERENDS, VAN  
NEERVEN;

BERENDS  
et al.

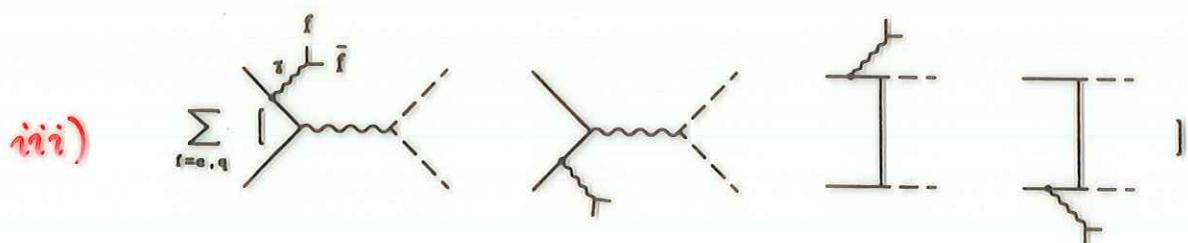
$$\begin{aligned}\Gamma_{ij}(z, L_m) = & \delta_{ij} \delta(1-z) + \frac{\alpha}{2\pi} P_{ij}^{(0)}(z) L_m \\ & + \frac{1}{2} \left( \frac{\alpha}{2\pi} \right)^2 \left[ \underline{\left( P_{ik}^{(0)} \otimes P_{kj}^{(0)} \right)(z) L_m^2} + P_{ij}^{(1)}(z) L_m \right] + \text{h.o.}\end{aligned}$$

$$P_{ee}^{(0)}(z) = \delta(1-z) \left[ \frac{3}{2} + 2 \ln \Delta \right] + \theta(1-\Delta-z) \frac{1+z^2}{1-z}$$

$$\begin{aligned}\frac{1}{2} \left[ P_{ee}^{(0)} \otimes P_{ee}^{(0)} \right](z) = & \delta(1-z) \left[ 2 \ln^2 \Delta + 3 \ln \Delta + \frac{9}{8} - 2\zeta(2) \right] \\ & + \theta(1-\Delta-z) \left\{ \frac{1+z^2}{1-z} \left[ 2 \ln(1-z) - \ln z + \frac{3}{2} \right] \right. \\ & \left. + \frac{1}{2}(1+z) \ln z - (1-z) \right\}\end{aligned}$$



$$\frac{1}{2} \left[ P_{e\gamma}^{(0)} \otimes P_{\gamma e}^{(0)} \right] (z) \equiv P_{e\gamma e}^{(1)}(z) = (1+z) \ln z + \frac{1}{2}(1-z) + \frac{2}{3} \frac{1}{z}(1-z^3)$$



$$P_{f\bar{f}}^{(1)}(z) = N_c(f) e_f^2 \frac{1}{3} P_{ee}^{(0)}(z) \theta \left( 1 - z - \frac{4m_f}{\sqrt{s}} \right)$$

$$\alpha(s) = \frac{\alpha}{1 - \frac{\alpha}{3\pi} \sum_f e_f^2 N_c(f) \ln \left( \frac{s}{m_f^2} \right)}$$

$$\begin{array}{lll} m_u & = & 62 \text{ MeV} \\ m_c & = & 1500 \text{ MeV} \end{array} \quad \begin{array}{lll} m_d & = & 83 \text{ MeV} \\ m_b & = & 4500 \text{ MeV} \end{array} \quad \begin{array}{lll} m_u & = & 215 \text{ MeV} \end{array}$$

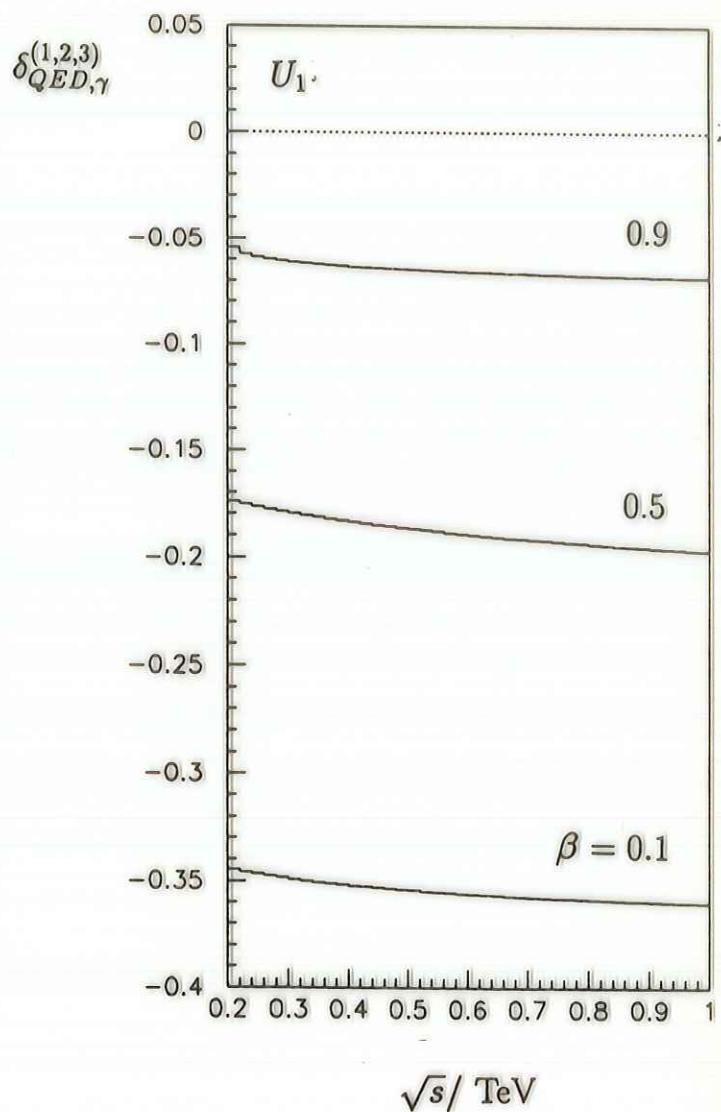
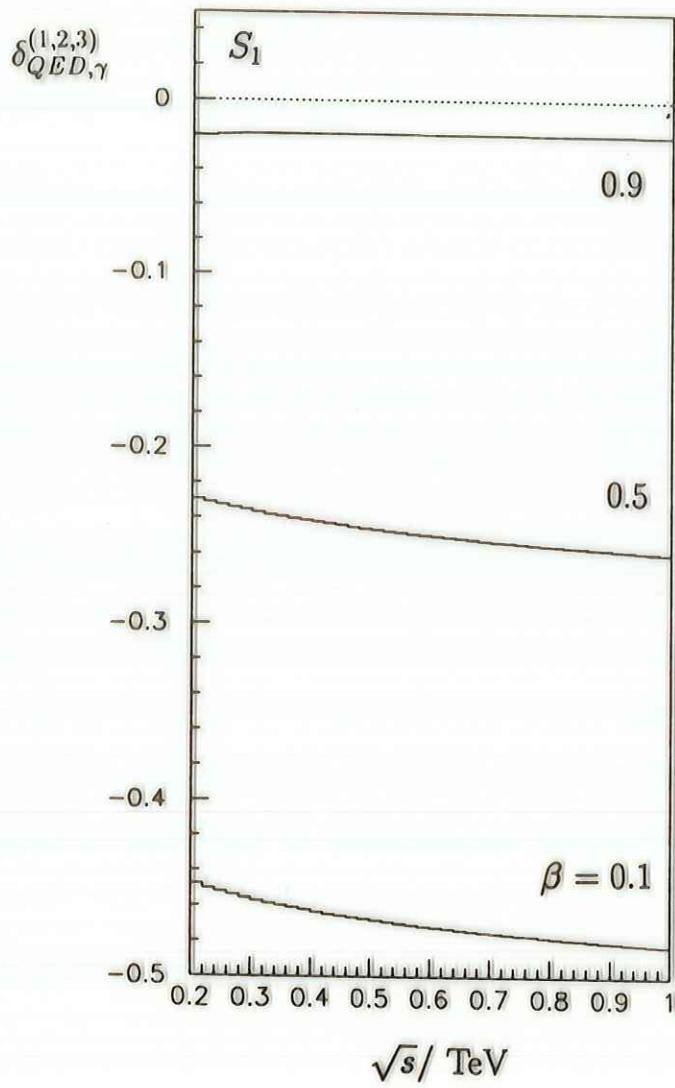
F. JEGERLEHNER

iv)

$$P_{soft}^{(3)}(z) = b(1-z)^{b-1}(1+\delta_1+\delta_2) - \frac{b(1+\delta_1)+b^2 \ln(1-z)}{1-z}$$

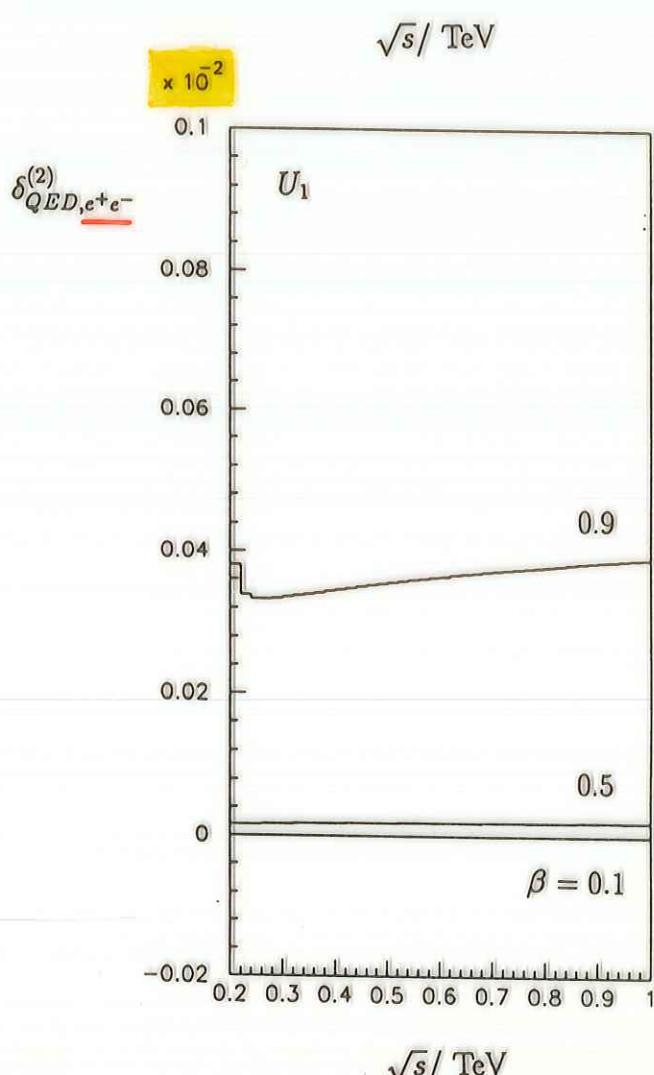
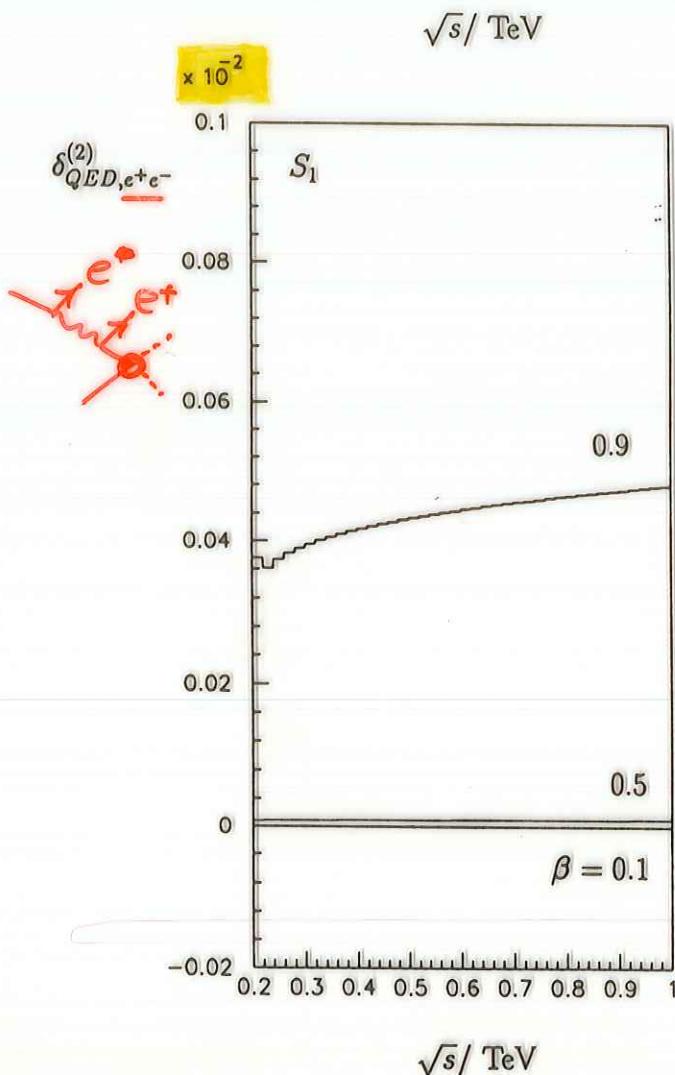
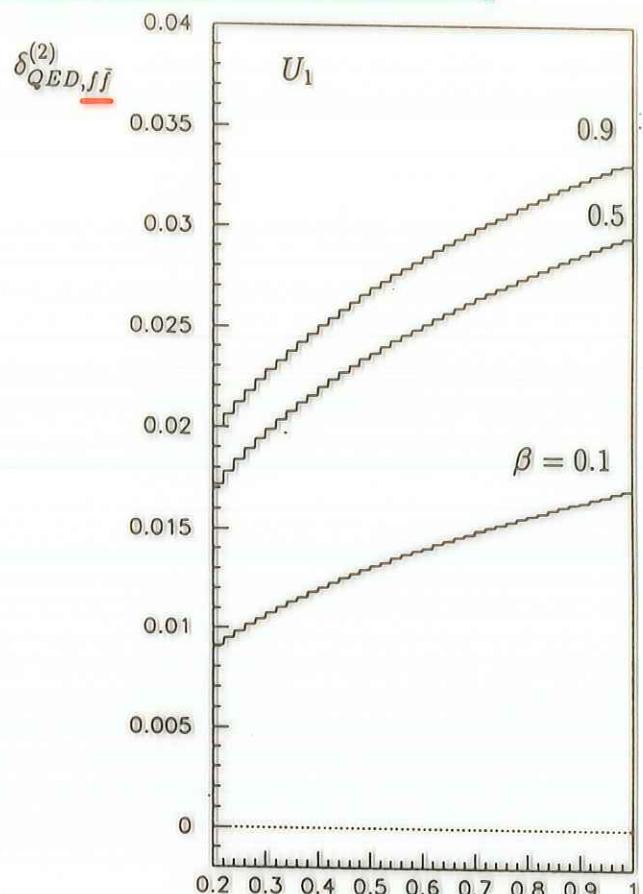
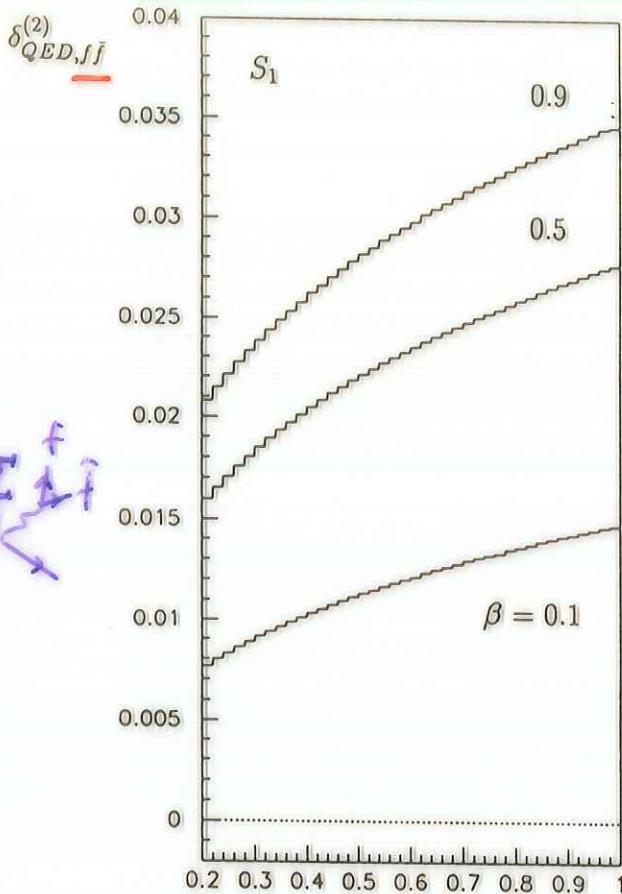
$$b = (2\alpha/\pi)(L_m - 1), \delta_1 = (3\alpha/2\pi)L_m \text{ and } \delta_2 = (\alpha/\pi)^2[9/8 - 2\zeta(2)]L_m^2$$

i) Bremsstrahlung:  
 $\mathcal{O}(\alpha)$ ,  $\mathcal{O}(\alpha^2)$ , soft exponentiation



$$\lambda_L = 0, \quad \lambda_R = e$$

### iii) $\mathcal{O}(\alpha^2)$ $e^+e^-$ and $f\bar{f}$ pair creation



### 3.2 Beamstrahlung

$$\frac{1}{\mathcal{L}} \frac{d\mathcal{L}}{dx} = \frac{x^a(1-x)^b}{B(a+1, b+1)}$$

à la PYTHIA

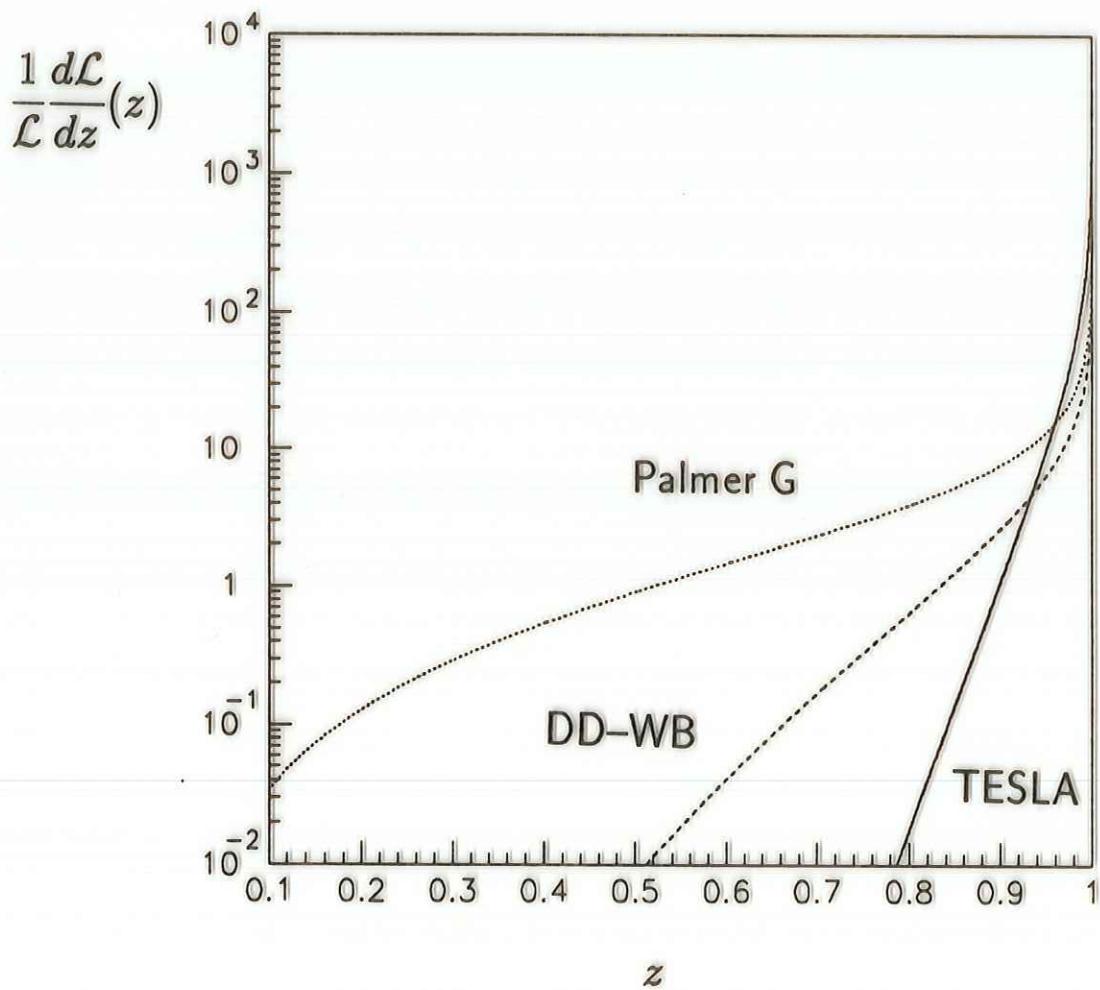
Accelerator	a	b
Palmer G	1.8	-0.67
DESY/Darmstadt WB	8.0	-0.67
TESLA	30.0	-0.85

- BARKLOW, CHEN  
KOZANECKI

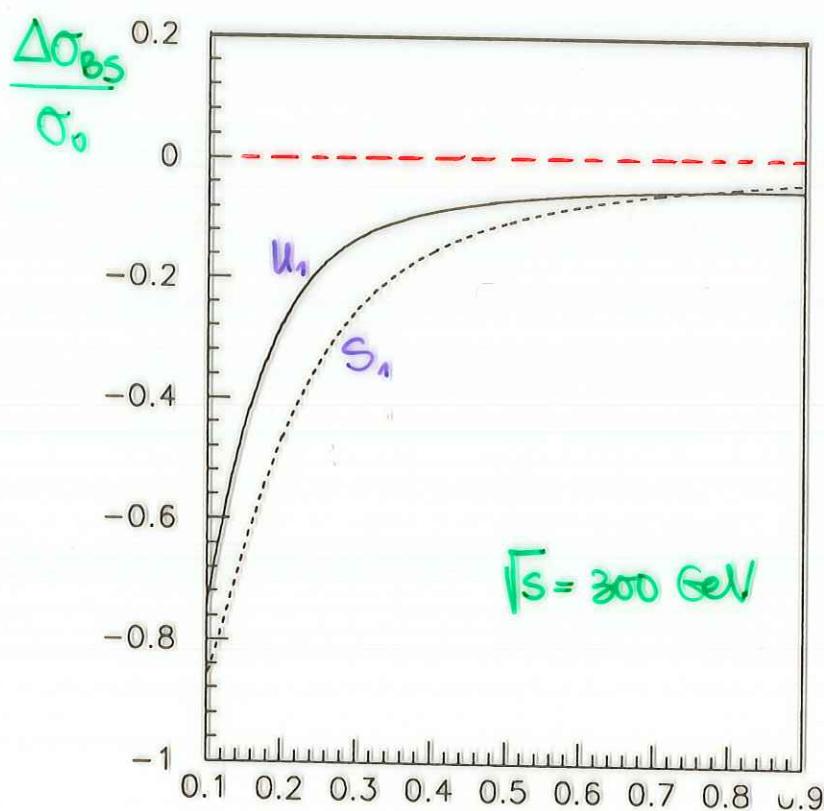
PARAM.: SCHREIBER

Table 1: Fitted parameters  $a$  and  $b$  in (35) [14]

$$\Delta\sigma_{BS}(s) = \int_0^1 dz \frac{1}{\mathcal{L}} \frac{d\mathcal{L}}{dz}(z) \sigma^{(0)}(zs) \theta\left(z - \frac{4M^2}{s}\right) - \sigma^{(0)}(s)$$



THRESH. RANGE:  $\beta \approx 0.1$        $\Delta\sigma_{BS} \approx -0.77$  ( $u_1$ )      DD, NB  
 $\Delta\sigma_{BS} \approx -0.87$  ( $s_1$ )  
(FURTHER CONFIG. TO BE INVESTIGATED.)

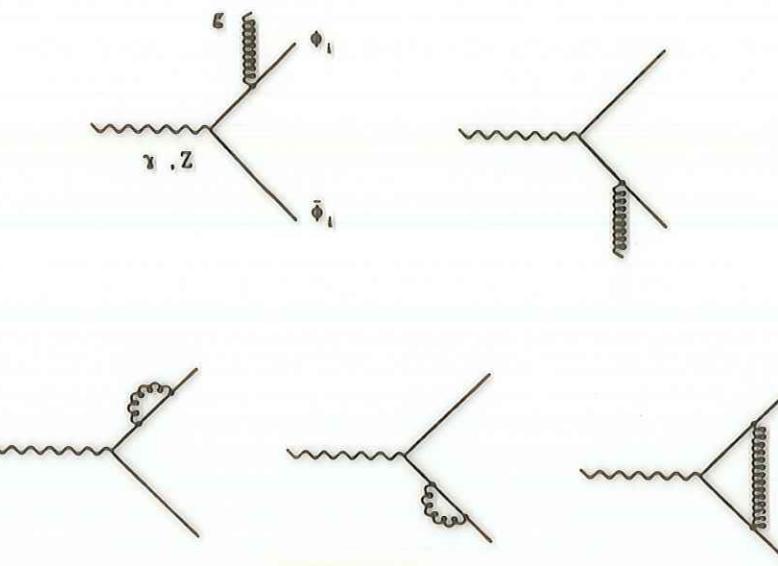


$\frac{1}{z} \frac{dL}{dz}$   
 DD-NBB  
 (ORTEU)

$\beta$

## 4 QCD Corrections

$$\lambda_L, \lambda_R \ll 1$$



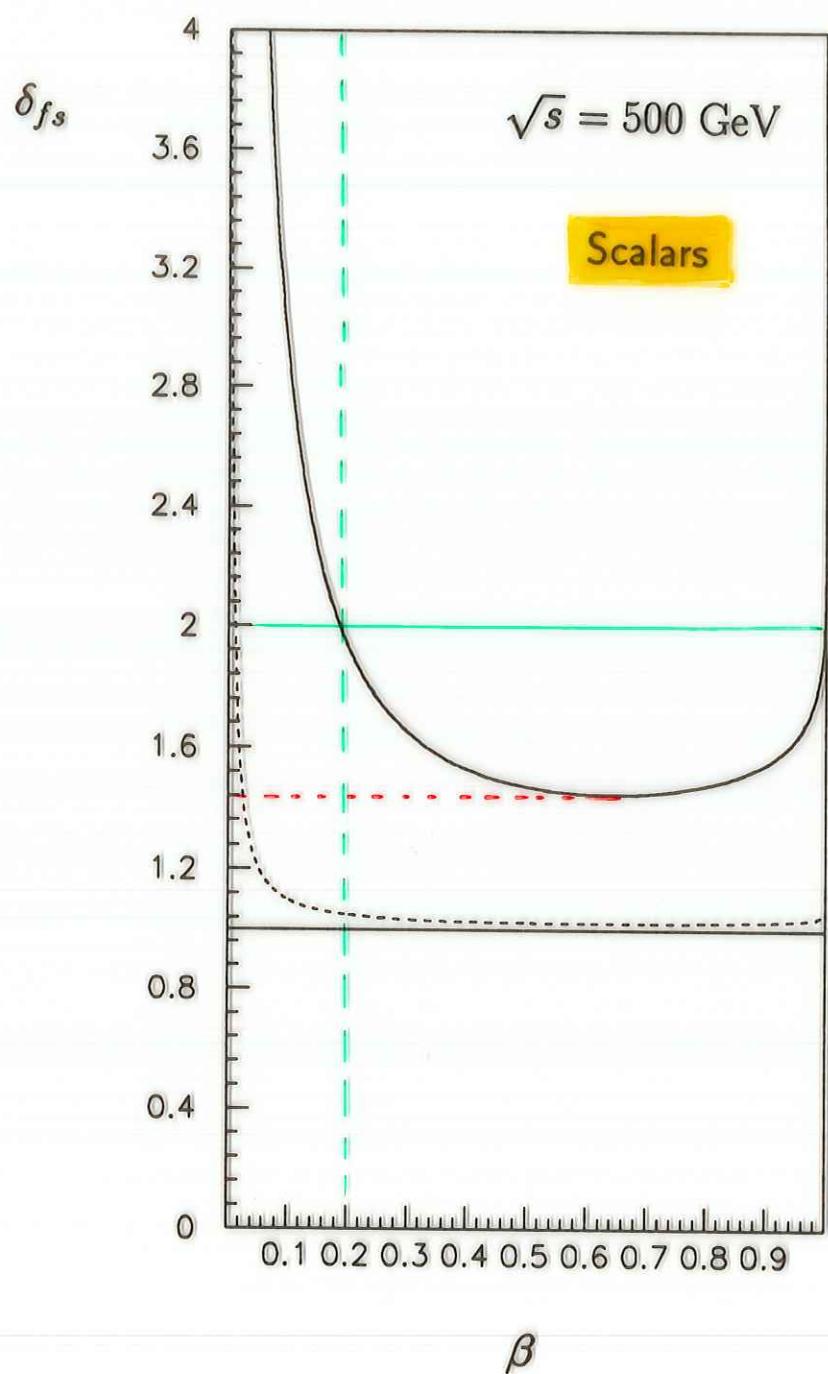
Scalars

$$\begin{aligned} \mathcal{F}_s(\beta) &= \frac{1+\beta^2}{\beta} \left[ 4Li_2\left(\frac{1-\beta}{1+\beta}\right) + 2Li_2\left(-\frac{1-\beta}{1+\beta}\right) - 3\ln\frac{2}{1+\beta}\ln\frac{1+\beta}{1-\beta} - 2\ln\beta\ln\frac{1+\beta}{1-\beta} \right] \\ &- 3\ln\left(\frac{4}{1-\beta^2}\right) - 4\ln\beta + \frac{1}{\beta^3} \left[ \frac{5(1+\beta^2)^2}{4} - 2 \right] \ln\frac{1+\beta}{1-\beta} + \frac{3}{2}\frac{1+\beta^2}{\beta^2} \end{aligned}$$

(J. SCHWINGER)

$$\sigma_{scalar}^{(1, QCD)}(s) = \sigma_{scalar}^{(0)}(s) \left\{ 1 + \frac{4\alpha_s}{3\pi} \mathcal{F}_s(\beta) \right\}$$

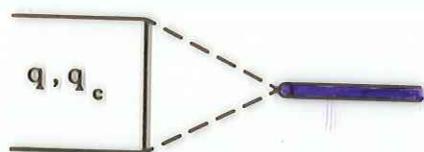
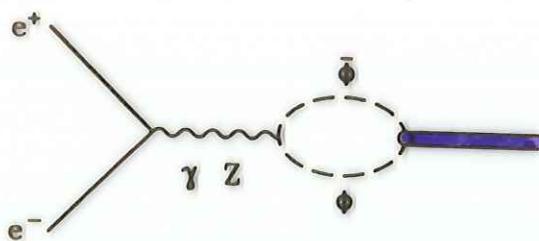
$$\sigma_{scalar}^{(1, QED, fs)}(s) = \sigma_{scalar}^{(0)}(s) \left\{ 1 + \frac{\alpha}{\pi} \mathcal{F}_s(\beta) \right\}$$



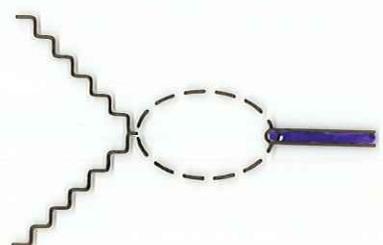
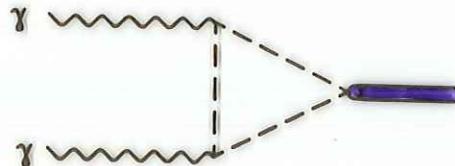
## Formation of Bound States

Leptoquarkonia

$$\beta \ll 1$$



However,  $\sigma \sim \beta^3$  !



SCHARS: → SQUARKONIA  
KÜHN, ZERWAS

$$\Gamma_0 = 0.4 \text{ GeV } f_{SV} \left(\frac{\Lambda}{e}\right)^2 \frac{M}{200 \text{ GeV}} ;$$

MASS  $\leftrightarrow$  LIFETIME

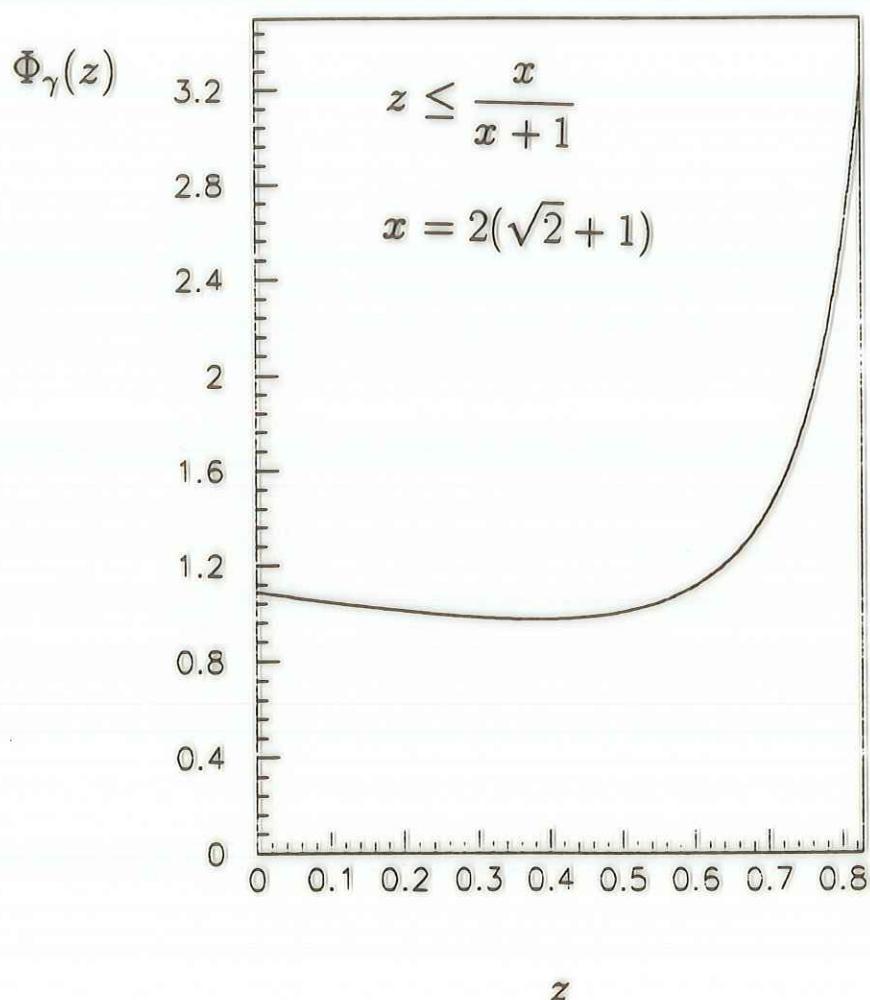
BOUND STATE FORMATION ?

DIFFICULT TO OBSERVE THE RESONANCE ,  $\sigma_{SV} \sim \beta^3$   
(ENERGY SPREAD !) for  $\beta \ll 1$ .

## 5 $\gamma\gamma$ Fusion

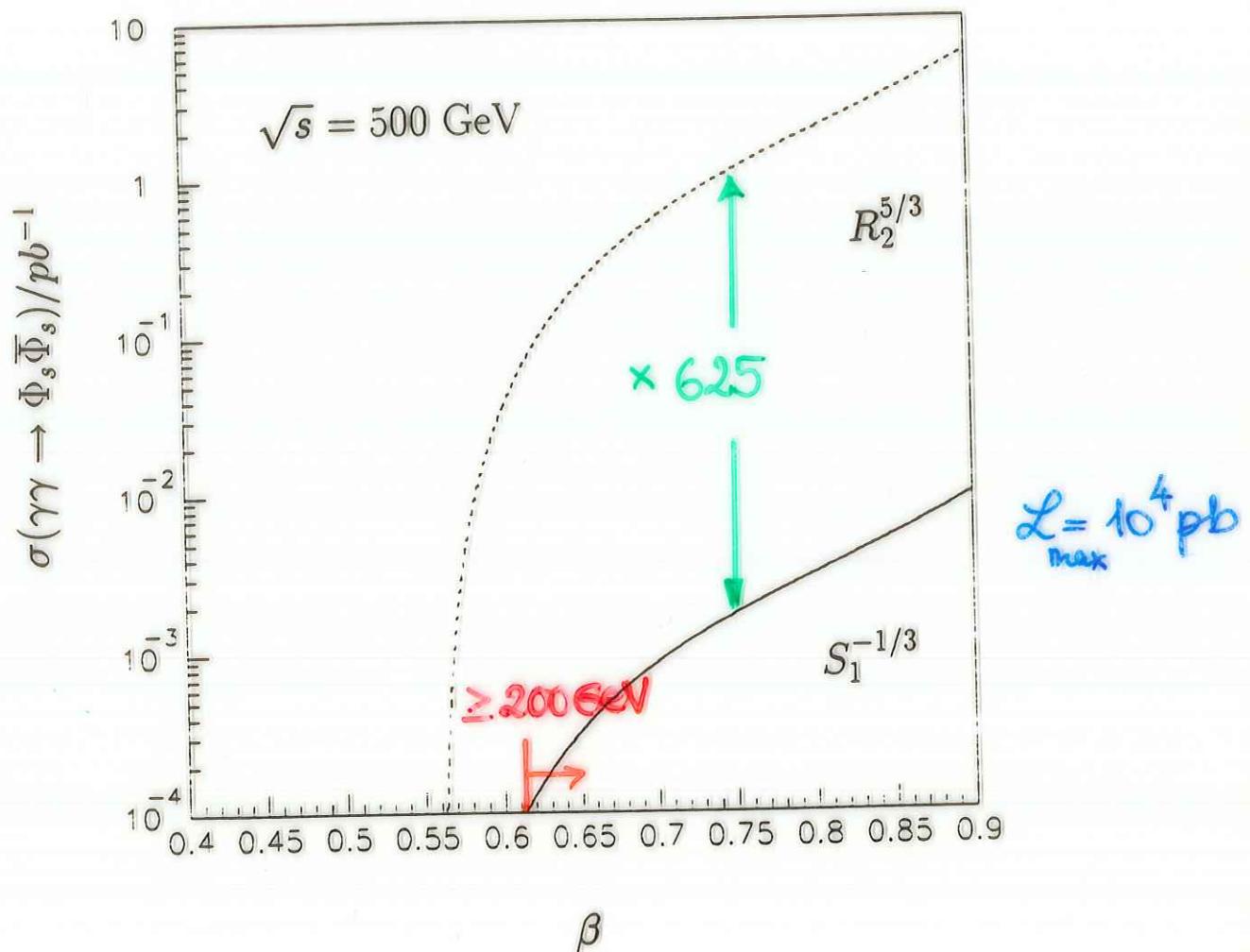
$$\Phi_\gamma(z) = \frac{1}{N(x)} \left[ 1 - z + \frac{1}{1-z} - \frac{4z}{x(1-z)} + \frac{4z^2}{x^2(1-z)^2} \right]$$

$$N(x) = \frac{16 + 32x + 18x^2 + x^3}{2x(1+x)^2} + \frac{x^2 - 4x - 8}{x^2} \ln(1+x)$$



$$\sigma_{scalar}(s) = \frac{\pi \alpha^2}{s} Q_\Phi^4 \left\{ 2(2 - \beta^2)\beta - (1 - \beta^4) \ln \left| \frac{1 + \beta}{1 - \beta} \right| \right\}$$

$$\sigma = \int_0^{z_{max}} dz_1 \int_0^{z_{max}} dz_2 \Phi_\gamma(z_1) \Phi_\gamma(z_2) \hat{\sigma}(\hat{s}) \theta(\hat{s} - 4M_\Phi^2)$$



## 6 Summary

- QED & QCD CORRECTIONS TO LEPTOQUARK PAIR PRODUCTION HAVE BEEN CALCULATED NUMERICALLY.

**QED ISR:**  $\beta = 0.1$

$\sqrt{s} = 500 \text{ GeV}$      $-47\% \text{ (incl.: } +8\% \text{ } \alpha^2 \text{ already)}$      $S_1$      $\gamma, \gamma\gamma \dots$   
 $-35\% \text{ } u_1$   
 $+3\% \text{ } e^+e^- \text{- pairs}$

**QCD:**  $\lambda_{c,e} \ll e$

$\delta \geq 42\%$     ;  $\beta \leq 0.2 \rightarrow \delta \geq 100\%$   
 ↳ BOUND STATE FORMATION.

**QED:** ISR     $\delta \approx 3\%$      $\beta = 0.1, \delta \approx 8\%$ .

- LARGE CORRECTIONS DUE TO BEAMSTRAHUNG  
 $\longrightarrow$  NBB REQUIRED;

- $\gamma\gamma$  FUSION : FAVOURS  $\frac{5}{3}$  LQ's     $\sigma \sim Q_{LQ}^4$   
 $\phi_s = R_2^{5/2}$  e.g.

$$(\sigma_i / \sigma_j) = (Q_i / Q_j)^4 \leq 625$$

RANGE:  $\beta \ll 1$  SUPPRESSED DUE TO  $\phi_\gamma(z)$ .