



On Bunch Compressor Optimization against Microbunching Instability and CSR

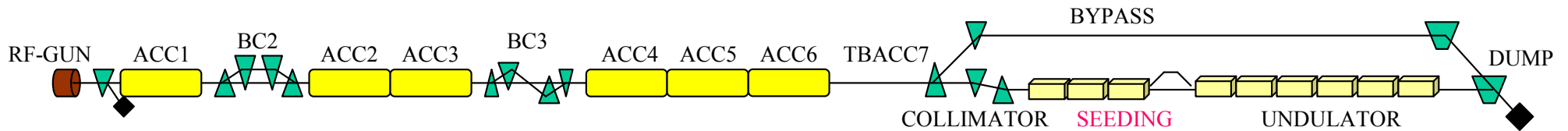
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DESY Hamburg, Germany

- Introduction to **Microbunching Instabilities** in X-ray FEL Facilities
 - **CSR** Microbunching Instability and **LSC** Microbunching Instability
- S2E Simulations on Microbunching Instability at TTF-2 Project
- Several Bunch Compressor Optimizations for XFEL Projects
 - Problems in old Linac Layout for TESLA XFEL
 - New Optimized Linac Layout for TESLA XFEL with TTF2 Injector
 - New Optimized Linac Layout for TESLA XFEL with XFEL Injector
 - Optimized Layout for PAL XFEL Project with S-band Photoinjector
- Summary
- Acknowledgments

TESLA Test Phase-2 FEL (TTF-2 FEL)



Layout of TESLA Test Facility Phase-2 (TTF-2)



courtesy of Prof. J. Rossbach

- 2003 Linac commissioning
Injector installation
- 2004 Injector commissioning
Complete vacuum installation
Beam through undulator
First Lasing @ 30 nm
- 2005 Saturation @30 nm
First users @ 30 nm
Tuning 30 – 120 nm
Full beam current
- 2006 Installation of seeding
and 3rd Harm. RF system
- 2007 6 nm
- 2008

Beam energy : 1.0 GeV

Peak current : 2.5 kA

rms bunch length : 50 μm ~ 166 fs

Bunch separation : 111 ns

Number of bunch per train : 7200

Bunch train repetition rate : 10 Hz

Slice transverse normalized rms emittance : 2 μm

rms beam size : 68 μm

rms energy spread : 1 MeV @ 1.0 GeV

Undulator period : 27.3 mm

Undulator gap : 12 mm

K-parameter : 1.17

Undulator length : ~ 30 m

Peak magnetic field : 0.495 T

SASE source wavelength : 6.4 nm ~ 120 nm

Saturation length : 27 m

Peak power : 2.8 GW

Average power : 40 W

FEL pulse length (FWHM) : ca. 100 fs

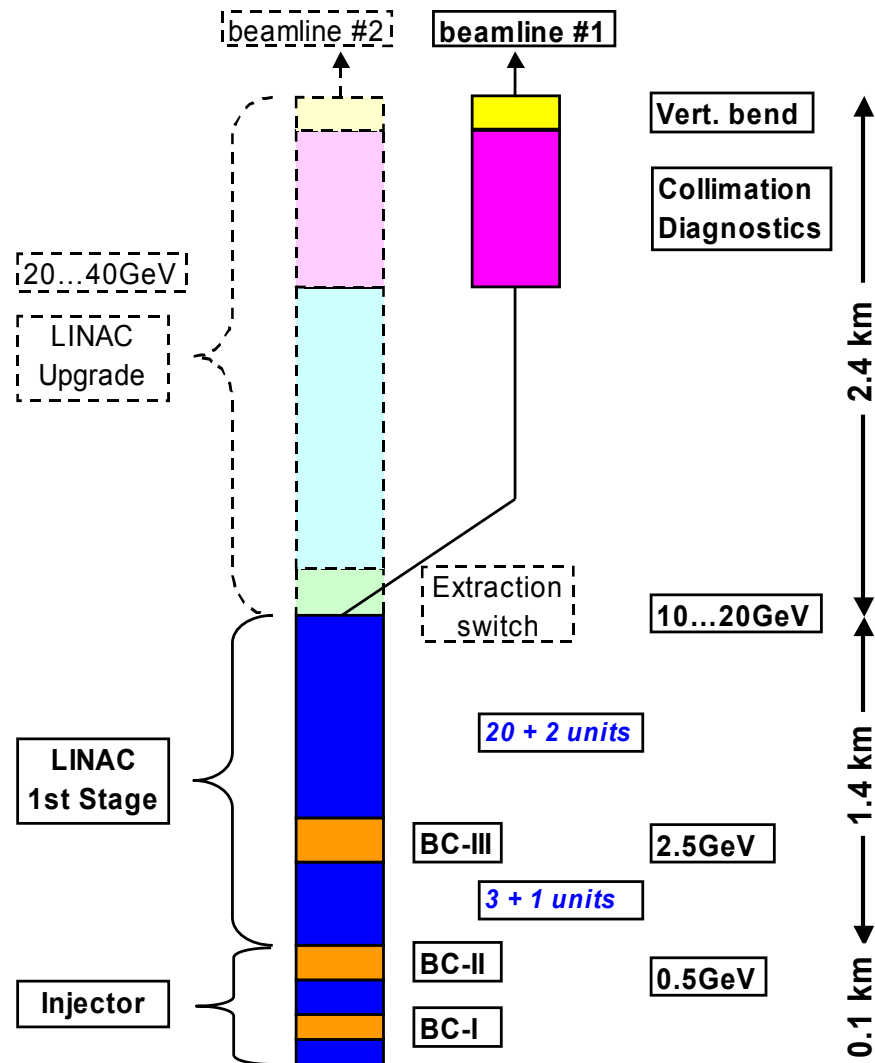
Peak (average) spectral brightness : 2.4×10^{30} (3.5×10^{22}) photons/sec/mrad²/mm²/0.1%BW

Under construction now !

Short Introduction to TESLA XFEL Project



Layout of 20 GeV TESLA XFEL Project

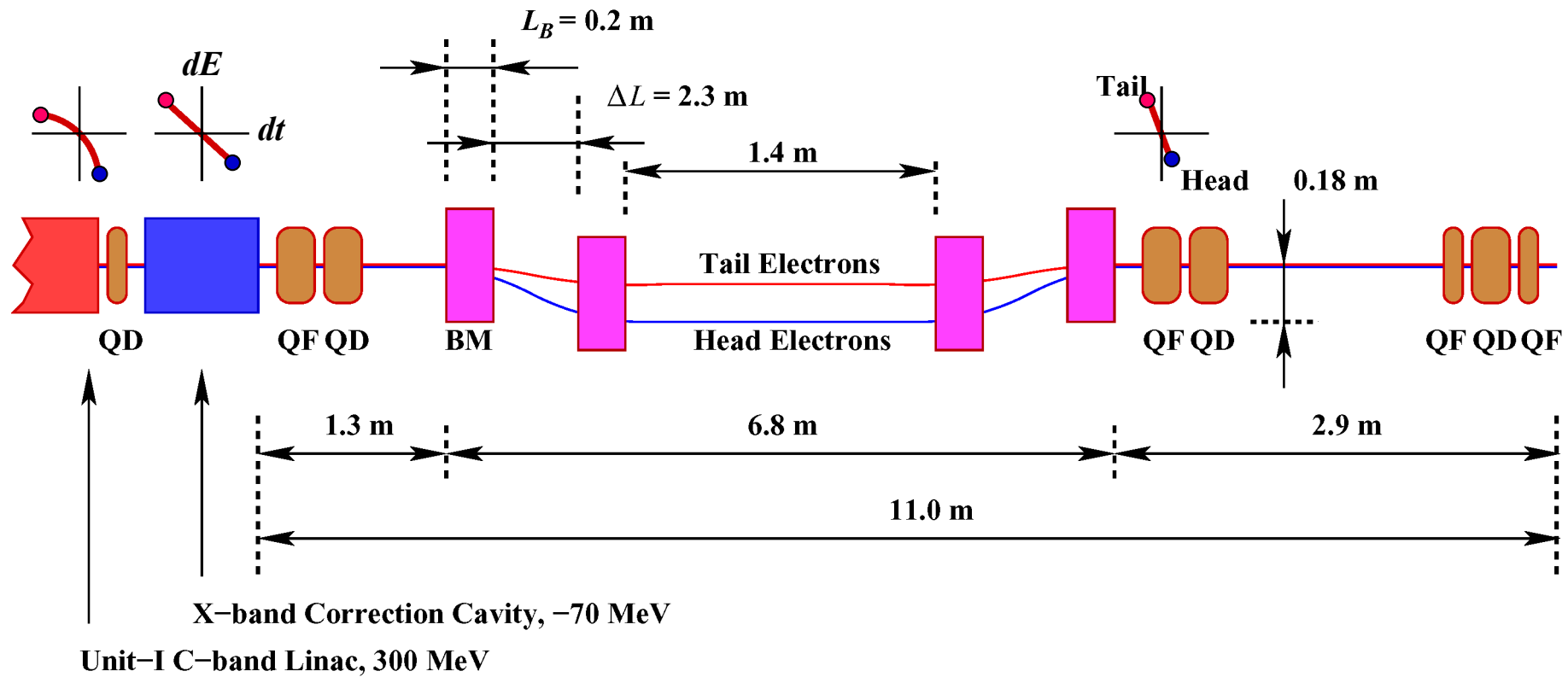


URL : <http://XFEL.DESY.de>

Already DESY got 50% of the total budget (684 Million Euro) for TESLA XFEL project from German government in February 2003.

Variable	Unit	Value
electron energy range	GeV	10-20
electron bunch length (rms)	fs	80
electron bunch charge	nC	1
normalized emittance	mmrad	1.4
uncorrelated energy spread (rms)	MeV	2.5
photon pulse length (FWHM)	fs	100
photon energy	keV	0.2 3.0 12.4
wavelength	nm	6.4 0.4 0.1
number of photons per bunch	$\times 10^{12}$	430 20 1.2
peak brilliance	$\times 10^{33}$	0.06 1.7 5.4
peak power	GW	135 100 24

Working Principle of Bunch Compressor



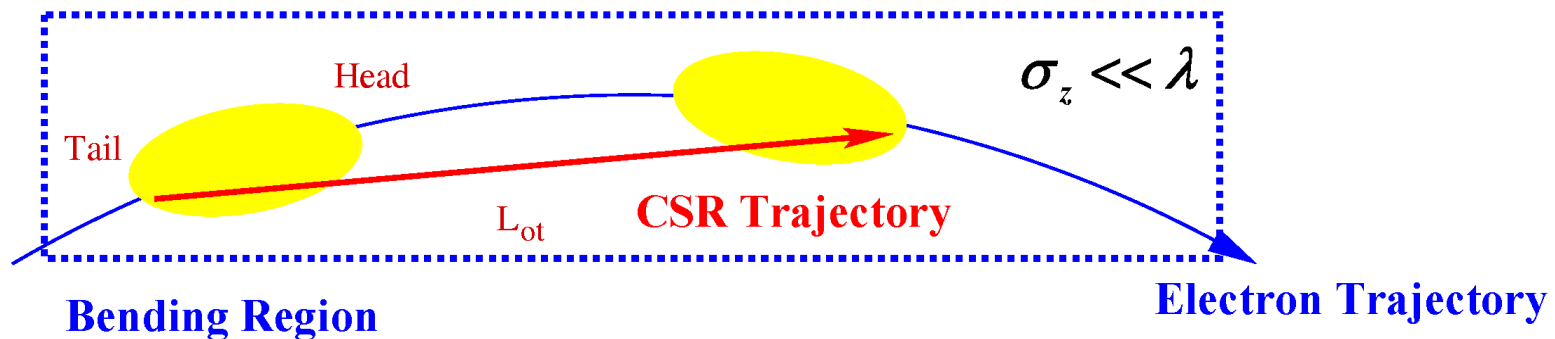
Introduction to CSR in BC



When the bunch length is smaller than the radiation wavelength, coherent synchrotron radiation (CSR) can be generated. **CSR from the tail region can overtake head electrons after the overtaking length.** Generally, head electrons are accelerated by CSR, and tail electrons are decelerated due to their own CSR loss.

Since CSR generates correlated energy spread along bunch length, electrons are transversely kicked at the nonzero dispersion region or chicane. **Hence, the projected normalized rms emittance is increased in the bunch compressor due to CSR.**

- Projected emittance growth effects brightness of the SASE source.
- Slice emittance growth effects saturation length and brightness of SASE source.



$$L_{ot} \approx \left(24\sigma_z R^2\right)^{1/3} \quad \sigma_{\delta,1\text{dipole}} \approx 0.22 \frac{r_e N L_B}{R^{2/3} \sigma_z^{4/3} \gamma}$$

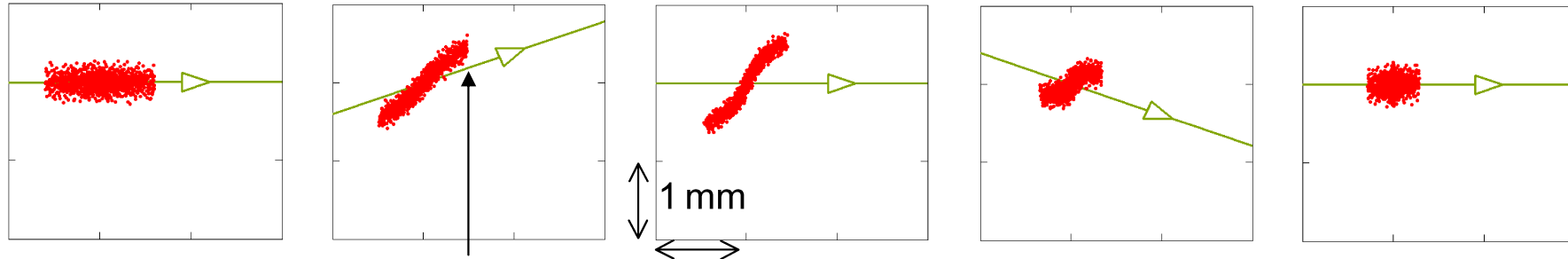
where σ_z is RMS bunch length, and R is bending radius.

Short Introduction to CSR in BC



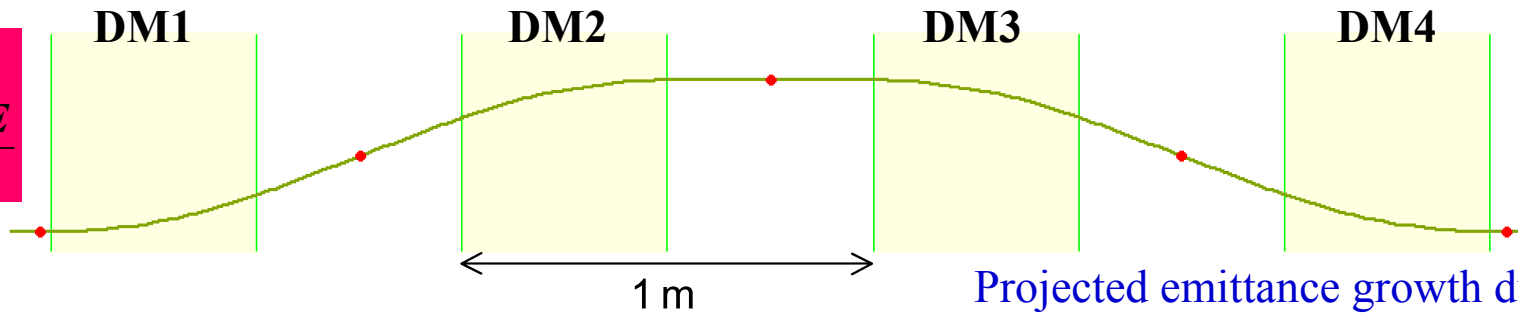
Without CSR self-interaction

Courtesy of M. Dohlus



Head with lower energy

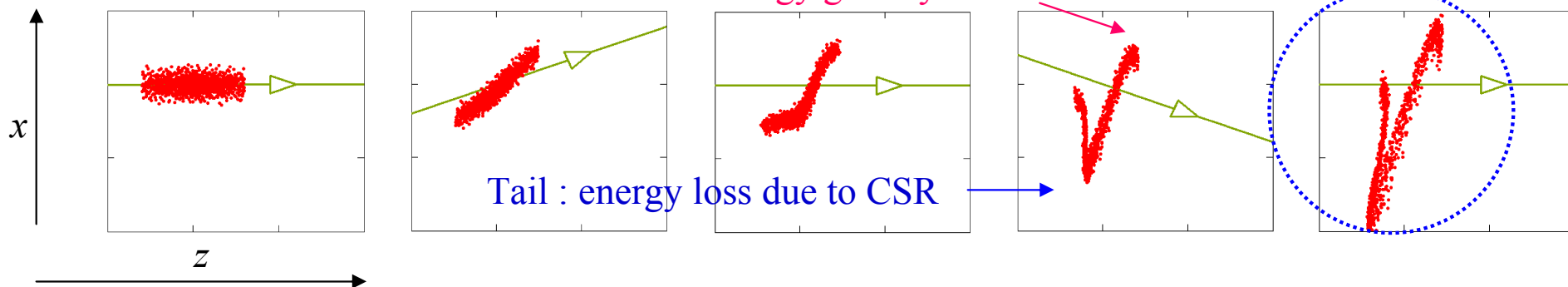
$\eta_x \neq 0$ in BC
 $x = x_\beta + \eta_x \frac{dE}{E}$



Projected emittance growth due to CSR

With CSR self-interaction

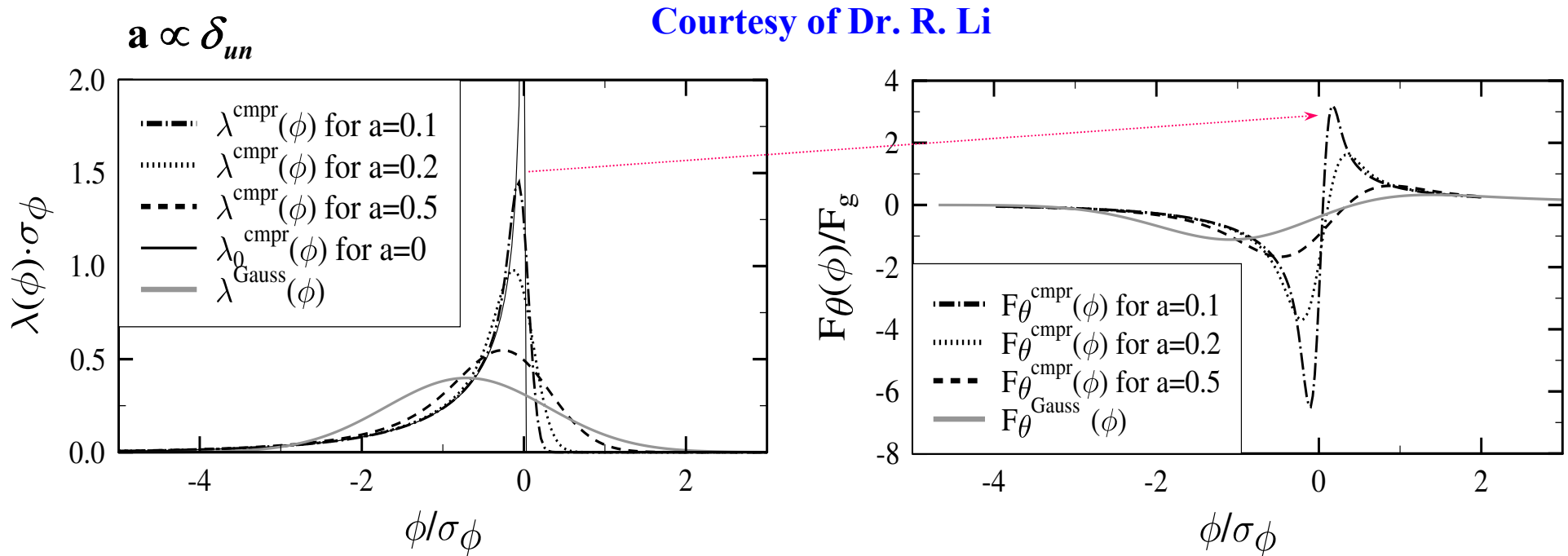
Head : energy gain by CSR



Short Introduction to CSR in BC



If there is a nonlinearity in current distribution, the local charge concentration or spike is generated during the bunch compression process. In this case, CSR is enhanced due to the local charge concentration. (R. Li, LINAC2000)



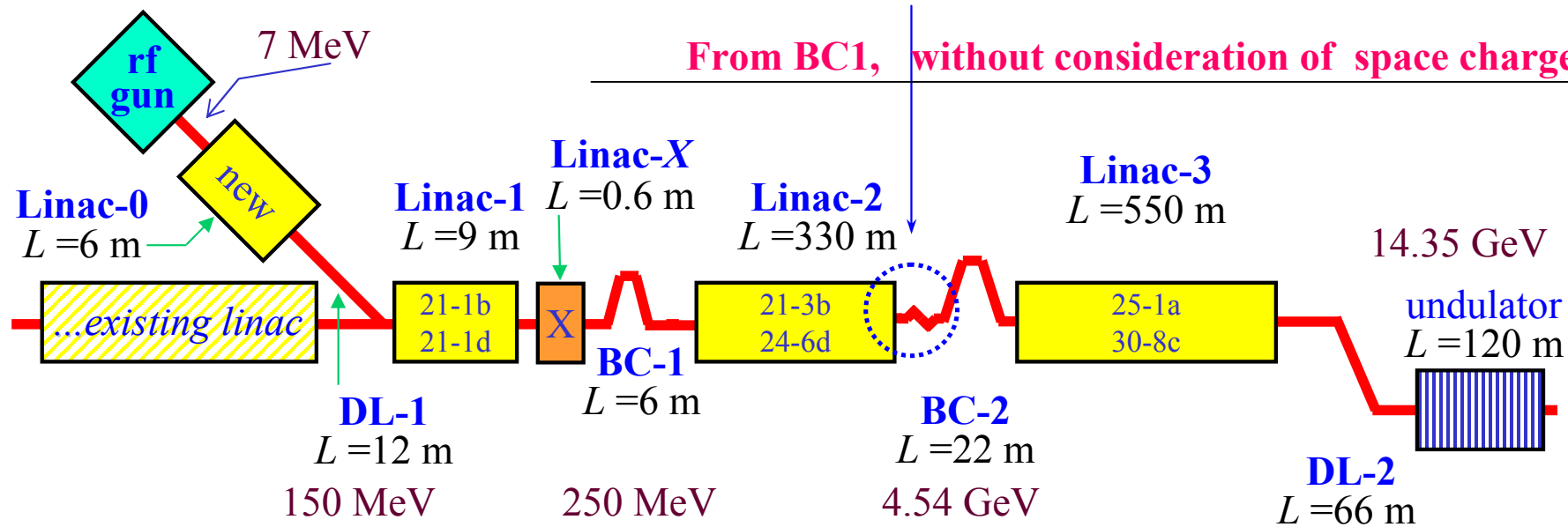
[Longitudinal Charge Distribution and CSR force for a compressed bunch]

Why does LCLS use SC-Wiggler before BC2 ?



Courtesy of Dr. P. Emma

$\sigma_\delta = 3E-5$ with SC-wiggler before LCLS BC-2
 σ_δ is increased by ISR in the wiggler, hence gain is reduced !



SLAC linac tunnel

Limitations in ELEGANT Tracking

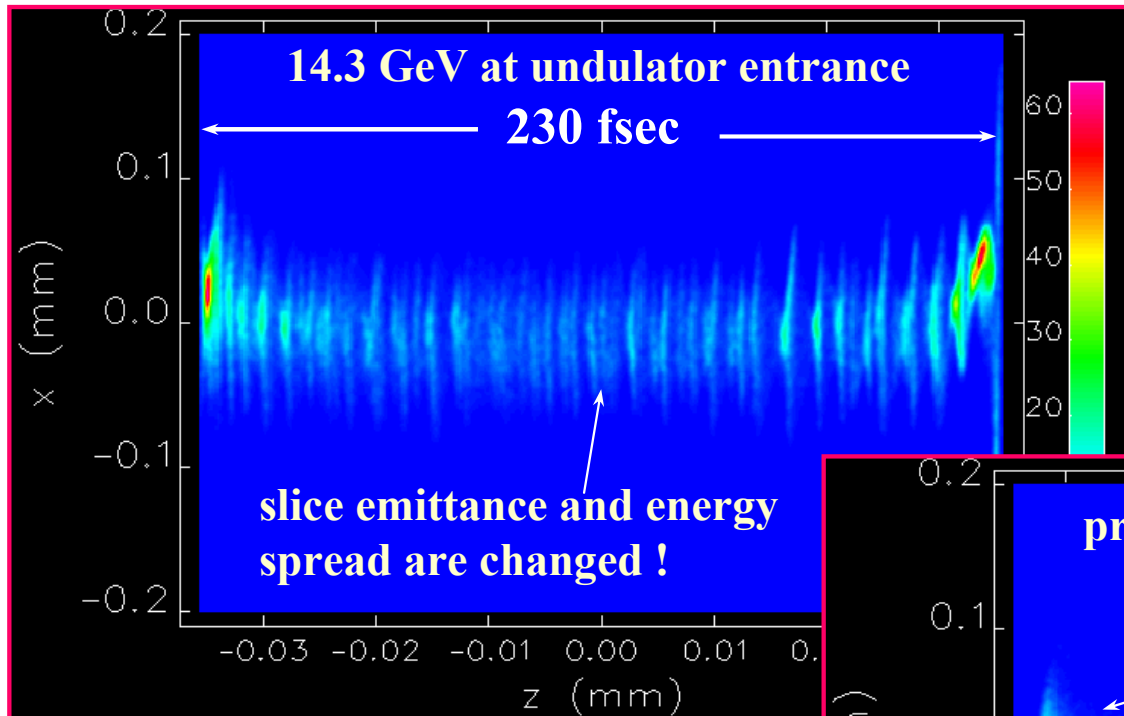
- Linear & nonlinear space charge effects
- Numerical error for small macro particles
- 1D CSR model

LCLS CSR Microbunching Instability

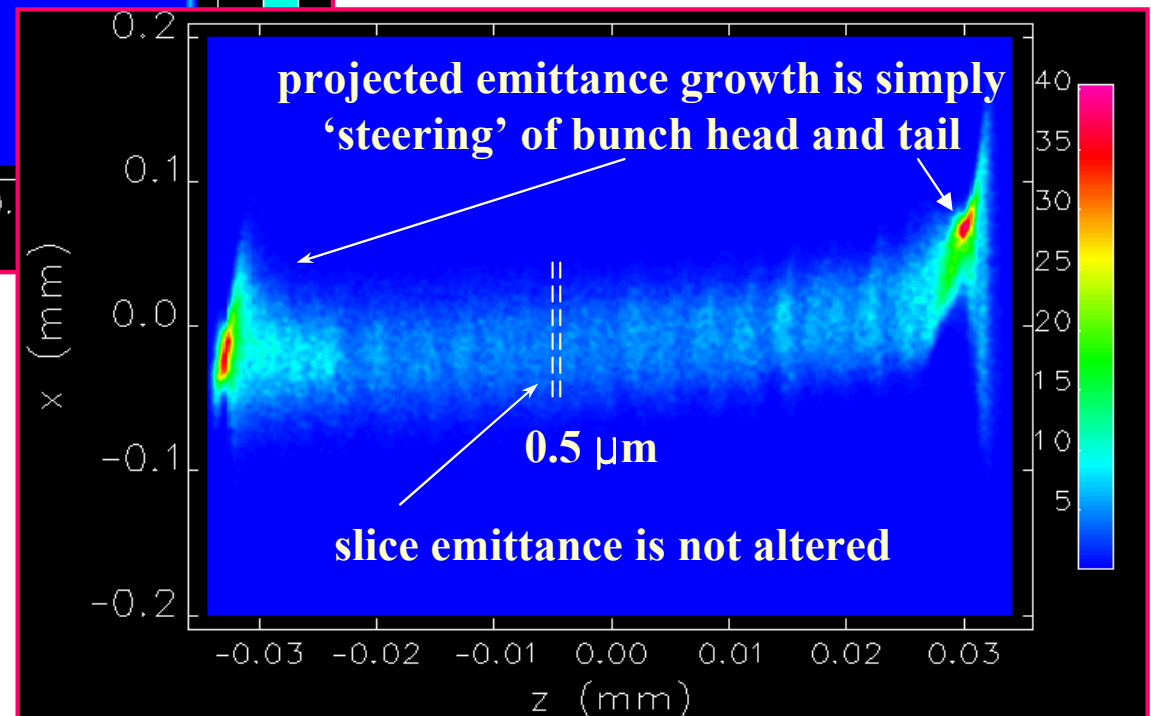


[x versus z without SC-wiggler]

Courtesy of Dr. M. Borland and P. Emma



[x versus z with SC-wiggler]



ELEGANT Tracking

without space charge force



- If initially, bunch has the density or energy modulation, CSR can be generated even for the wavelengths much shorter than the bunch length, and the initial modulation can be amplified in the bunch compressor by the CSR force (Microbunching Instability).
- Initial modulation can be converted in the bunch compressor
 - Energy modulation can be converted to density modulation via dispersion.
 - Density modulation can be converted to energy modulation via CSR.
- If the uncorrelated energy spread and transverse emittance are high enough, the amplification gain G can be reduced substantially, and the microbunching instability can be controlled by Landau damping.

Gain of CSR Microbunching Instability



Dr. Huang's PRST-AB Vol5, 074401, Gain of CSR Microbunching Instability G_f :

$$G_f = \left| b(k_f; f) / b(k_0; 0) \right|$$

$$\begin{aligned} \frac{b(k_f; f)}{b_0(k_0; 0)} &\approx \exp \left[-\frac{\bar{\sigma}_\delta^2}{2(1 + \chi)^2} \right] + A\bar{I}_f \left[\frac{\sqrt{\pi} \operatorname{erf}(\bar{\sigma}_x)}{2\bar{\sigma}_x} \right. \\ &\times \exp \left(-\frac{\bar{\sigma}_\delta^2}{2(1 + \chi)^2} \right) + F_1(\chi, \bar{\sigma}_x, \alpha_0, \phi, \bar{\sigma}_\delta) \left. \right] \\ &+ A^2 \bar{I}_f^2 F_0(\bar{\sigma}_x) F_2(\chi, \bar{\sigma}_x, \alpha_0, \phi, \bar{\sigma}_\delta). \end{aligned} \quad (1)$$

Here $\chi = hR_{56}$, $\bar{\sigma}_\delta = k_0 R_{56} \sigma_\delta$, $A = 1.63i - 0.94$, $\bar{I}_f = I_f k_0^{4/3} R_{56} L_b / (\gamma I_A \rho_0^{2/3})$, I_f is the compressed beam current, $I_A \approx 17$ kA, $\bar{\sigma}_x = k_0 L_b \sqrt{\varepsilon_x \beta_0} / \rho_0$, $\phi = 2\Delta L / \beta_0$, the error function $\operatorname{erf}(x) = 2\pi^{-1/2} \int_0^x dt \exp(-t^2)$, and

Limitations of Formula

Uncorrelated energy spread growth in BC due to ISR, CSR, and LSC

Uncorrelated energy spread growth in LINAC due to LSC, and IBS (later added !)

$$\begin{aligned} F_0 &= \frac{e^{-\bar{\sigma}_x^2} + \bar{\sigma}_x \sqrt{\pi} \operatorname{erf}(\bar{\sigma}_x) - 1}{2\bar{\sigma}_x^2}, & k_f &= k_0 / (1 + hR_{56}) \\ F_1 &= 2 \int_0^1 dt \frac{(1-t)}{(1+\chi t)^{4/3}} H(t), & k_0 &= 2\pi / \lambda \\ F_2 &= 2 \int_0^1 dt \frac{(1-t)t(1+\chi)}{(1+\chi t)^{7/3}} H(t), \\ H(t) &= \exp \left[-\bar{\sigma}_x^2 \frac{(1-2t + \alpha_0 \phi t)^2 + \phi^2 t^2}{(1+\chi t)^2} \right. \\ &\quad \left. - \frac{\bar{\sigma}_\delta^2}{2(1+\chi t)^2} \left(t^2 + \frac{(1-t)^2}{(1+\chi)^2} \right) \right]. \end{aligned} \quad (2)$$

b : bunching parameter, Fourier transform of the current modulation

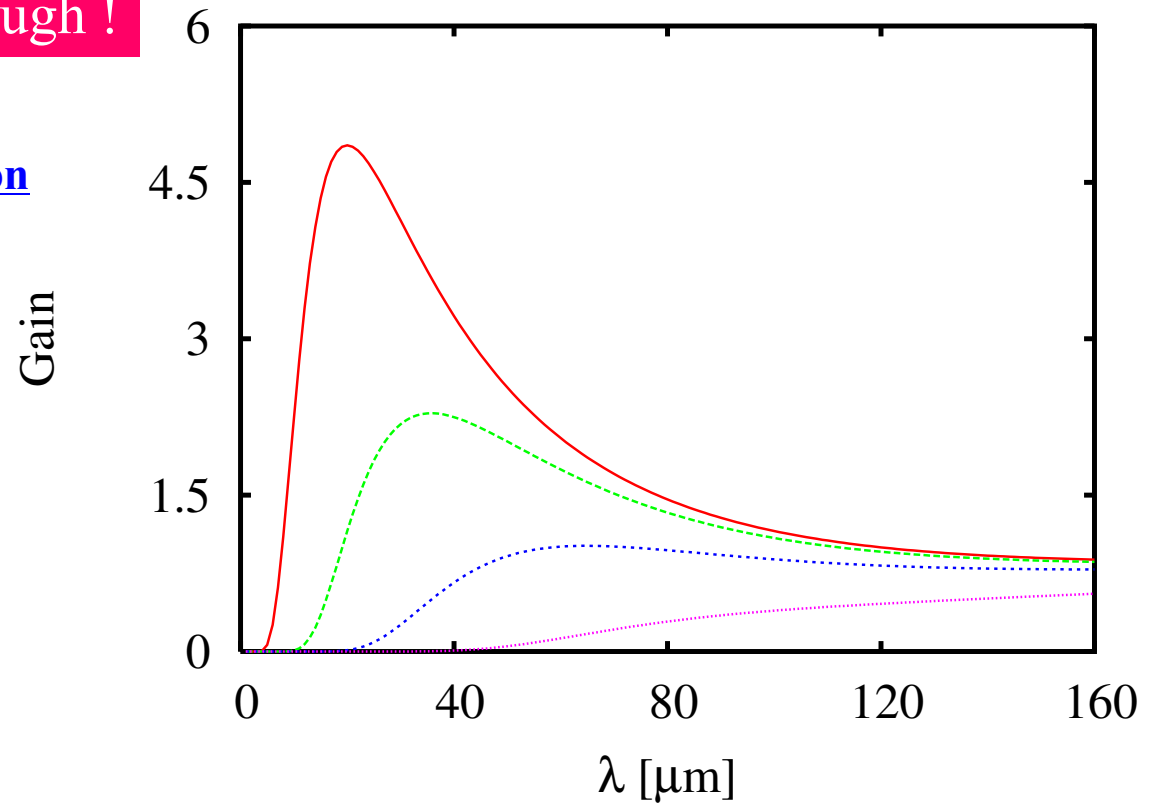
CSR Gain of Model Chicane



Gain of single chicane is small enough !

Common Parameters for Gain Calculation

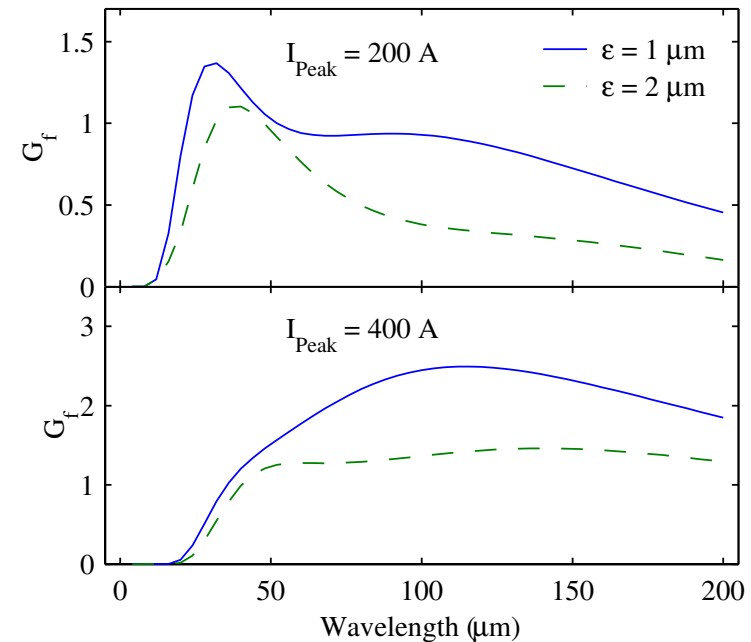
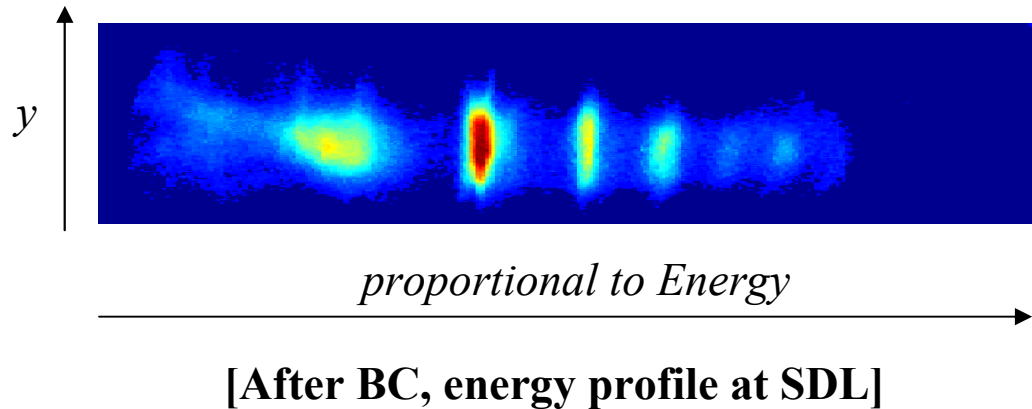
- Dipole length $L_b = 0.5$ m
- Drift length $\Delta L = 5.0$ m
- Middle drift length $L_i = 1.0$ m
- Bending radius $\rho_0 = 10.35$ m
- Bending angle = 2.77 degree
- Final peak current $I_f = 6000$ A**
- Modulation amplitude $\sim 5\%$
- Initial beta function $\beta_0 = 40$ m
- Initial alpha function $\alpha_0 = 2.6$
- Energy = 5.0 GeV
- Total chicane length = 13.0 m



Parameters used for plotting

			IBS	TTF-2		FEL
Momentum compaction	R_{56}	[mm]	+25	+25	+25	+25
Chirping constant	h	[1/m]	-36	-36	-36	-36
Normalized rms emittance	ϵ_{nx}	[μm]	1.0	1.0	1.0	1.0
Uncorrelated energy spread	$\sigma_{\delta u}$	10^{-6}	9.5	25	50	100

We observed following thing similar to the microbunching instability at SDL !
For a while, we considered their observation as a CSR microbunching instability due to laser.
But gain of CSR microbunching instability for SDL chicane is small enough !
Then what is source of their observation ?
The most strong candidate is the Longitudinal Space Charge (LSC)

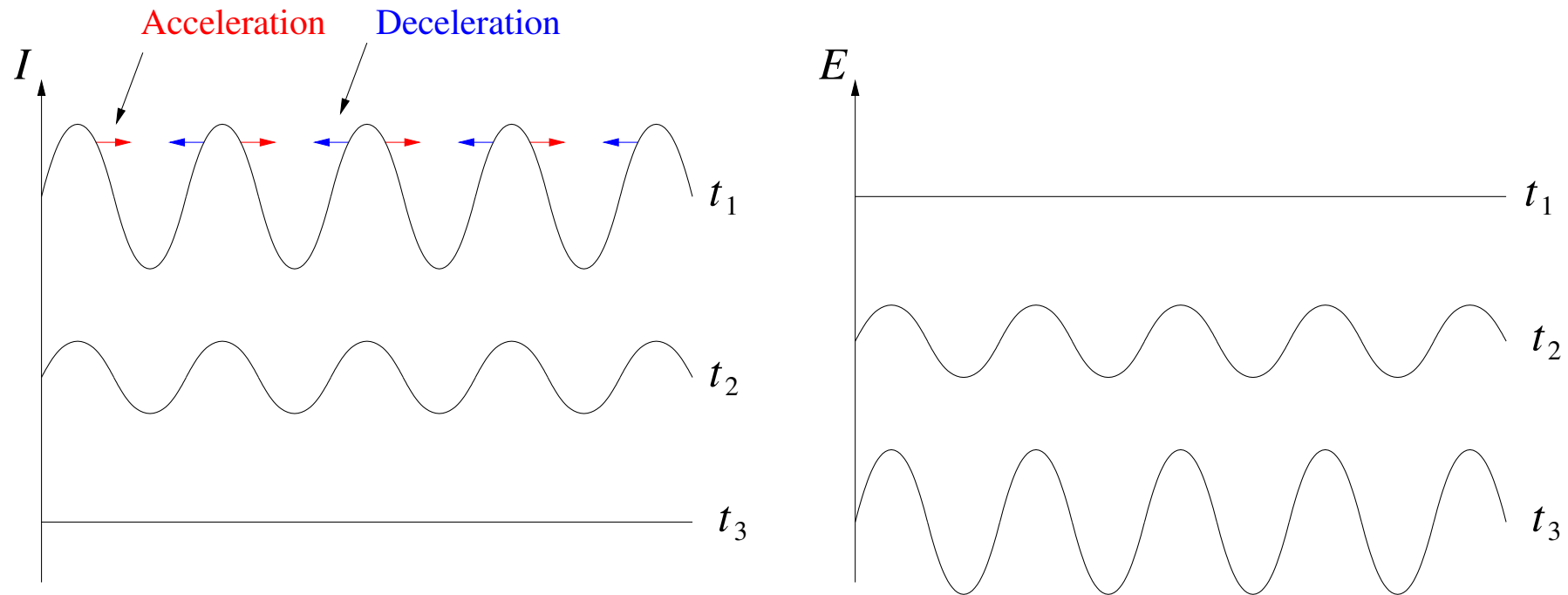


[Gain of Microbunching Instability]

Action of Longitudinal Space Charge (LSC)



If there is a density modulation, space charge pushes particles from high density to low density, creating energy modulation in the process.



Density modulation



Energy modulation

Energy modulation converts back to density modulation to complete space charge oscillation with Plasma oscillation frequency (Here Landau damping is ignored)

CSR & LSC Microbunching Instability



by **longitudinal space charge force (LSC)** in the drift & linac
by **coherent synchrotron radiation (CSR)** in the bunch compressor
by the **geometric wakefields** in Linac

**Initial current
density modulation**



**Induced
energy modulation**



$$dz_f = dz_i + R_{56}(dE/E)_i \text{ in BC}$$

Slice emittance
Projected emittance
Slice energy spread
Growth !!!

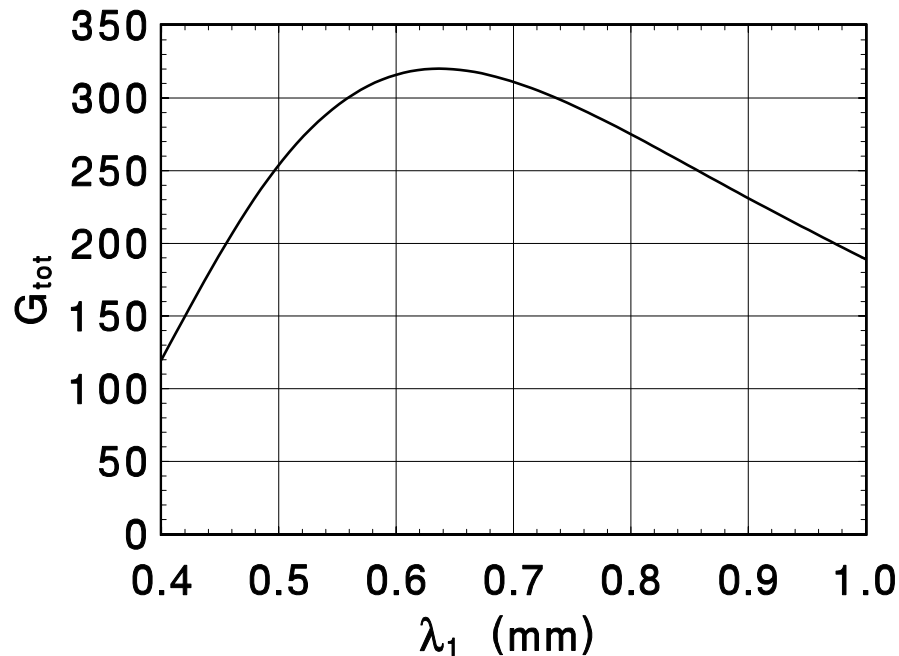
**Amplified current
density modulation**

$$\text{Gain} = \left| \frac{\text{Normalized amp. of current density}_2}{\text{Normalized amp. of current density}_1} \right|$$

□ Motivation : Longitudinal Space Charge (LSC) driven Microbunching

Instability TESLA-FEL-2003-02 by Dr. Schneidmiller *et al.*

Tu-P-49 at FEL2003



Maximum gain of the LSC driven microbunching instability in TTF-2 is about 320 for 2.0 ps initial density modulation.

Here they assumed that Plasma oscillation and Landau damping are ignorable. Therefore the initial modulation amplitude is constant and is not damped although beams go through the linac.

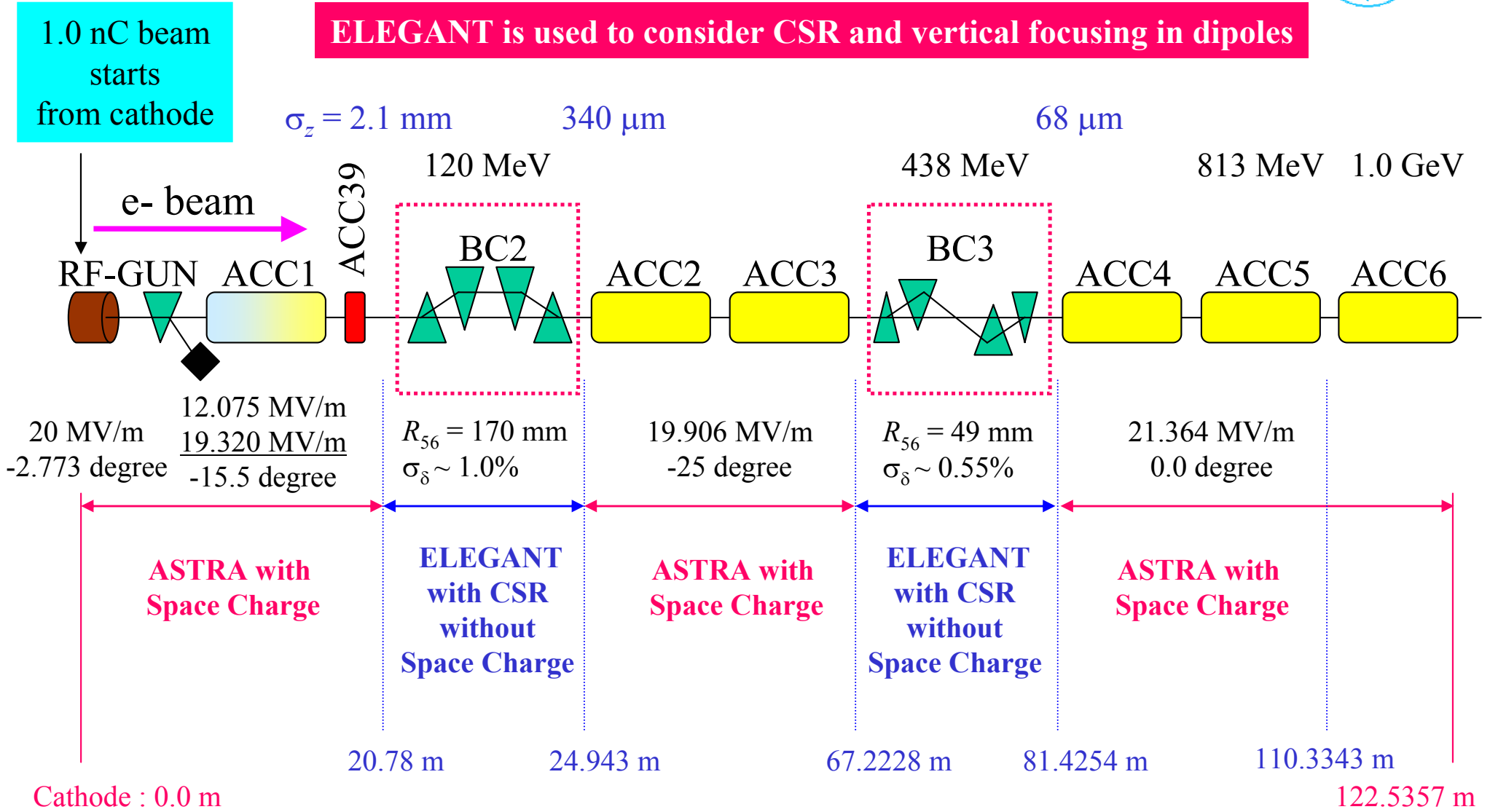
$$G = Ck |R_{56}| \frac{I_0}{\gamma_0 I_A} \frac{|Z(k)|}{Z_0} \exp\left(-\frac{1}{2} C^2 k^2 R_{56}^2 \frac{\sigma_\gamma^2}{\gamma_0^2}\right)$$

□ A S2E simulation is needed to check whether the gain is really high or not.

S2E on Microbunching Instability at TTF-2



ELEGANT is used to consider CSR and vertical focusing in dipoles



ELEGANT is called at all ends of module to apply geometric wakefields !!!

ASTRA Tracking with Laser Beam Modulation



From Cathode to Before ACC1- 1.0 m

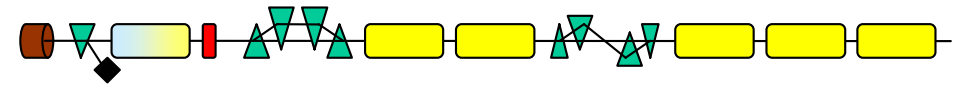
Modulation period = 2.0 ps

Modulation Amplitude = $\pm 10\%$

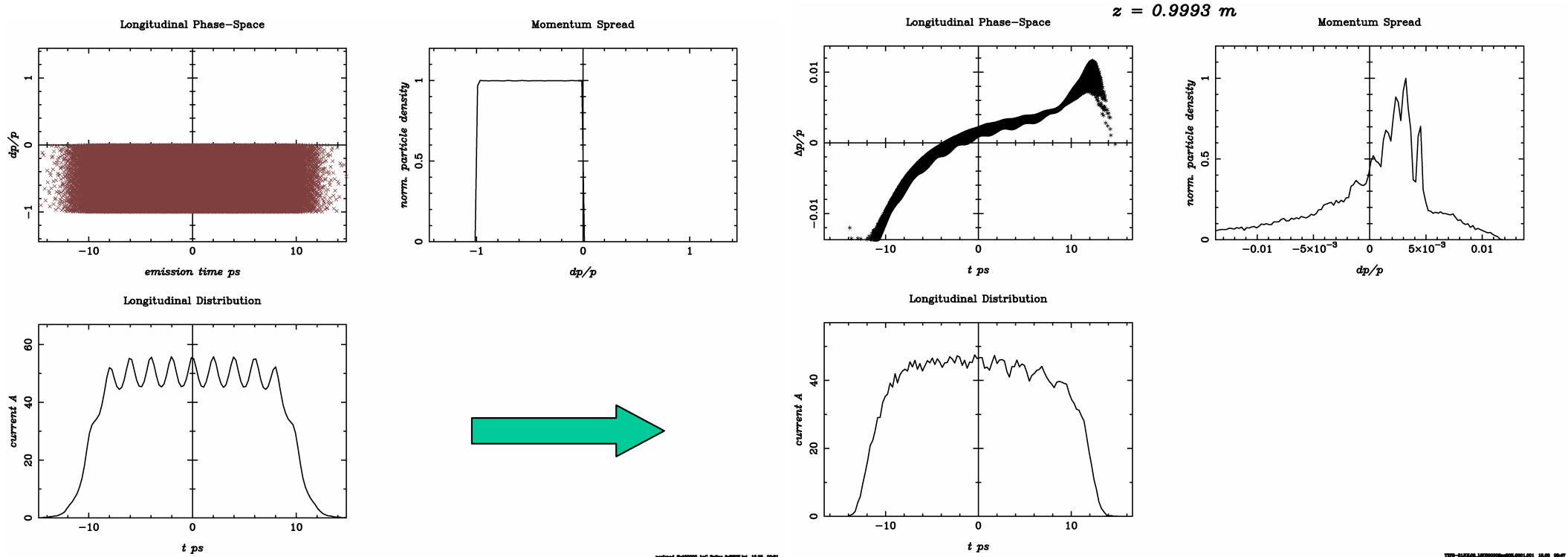
Postpro BIN number = 125

Simulation particles = 50,000

ASTRA Grids for space charge calculation = 10 rings \times 50 cells



↑ Here !



ASTRA Tracking with Laser Beam Modulation



From end of ACC1 to Upstream of BC2

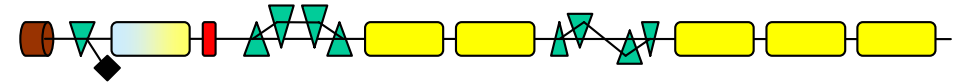
Modulation period = 2.0 ps

Modulation Amplitude = $\pm 10\%$

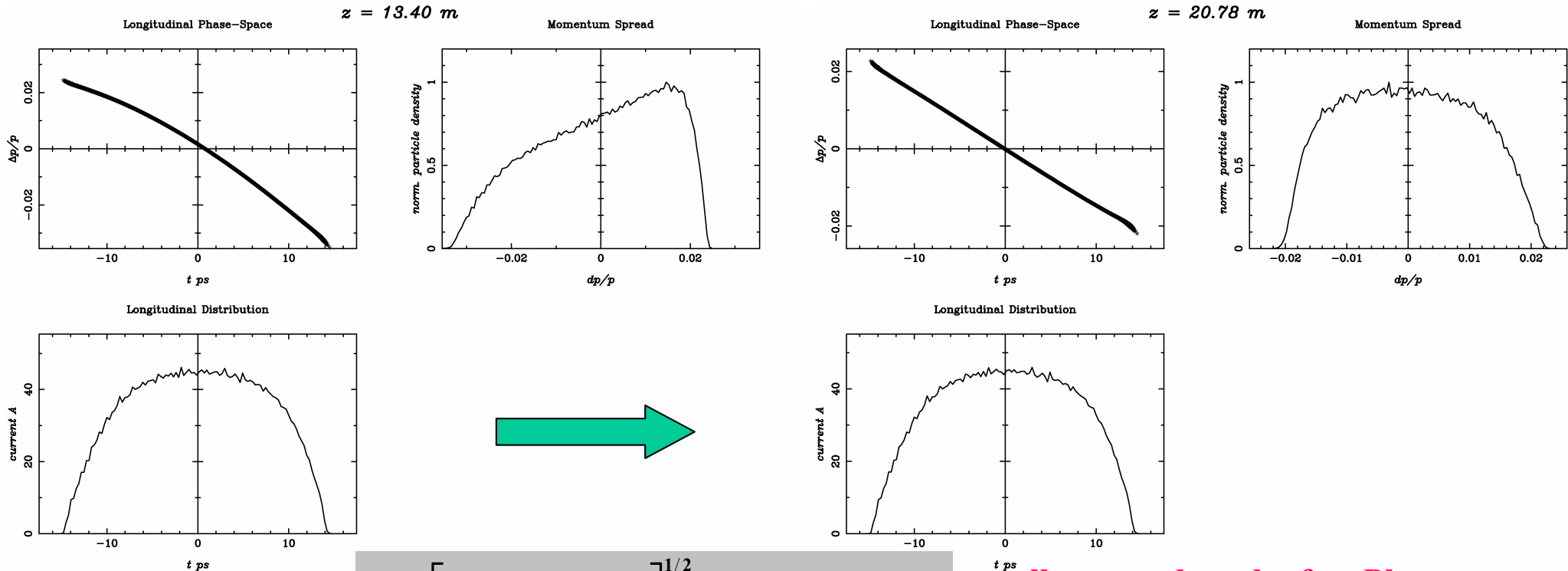
Postpro BIN number = 125

Simulation particles = 50,000

ASTRA Grids for space charge calculation = 10 rings \times 50 cells



↑ Here !



$$\omega = c \left[\frac{I_0}{\gamma^3 I_A} k |Z_{LSC}(k)| \right]^{1/2}, \text{ if } k \rightarrow \infty, \omega \rightarrow \omega_p$$

well smeared out by fast Plasma oscillation & strong Landau damping

What happened in injector ?



E-fields due to Space Charge at 1.0 m, $E=4.2$ MeV

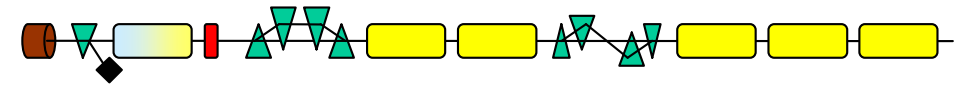
Modulation period = 2.0 ps

Modulation Amplitude = $\pm 10\%$

Postpro BIN number = 125

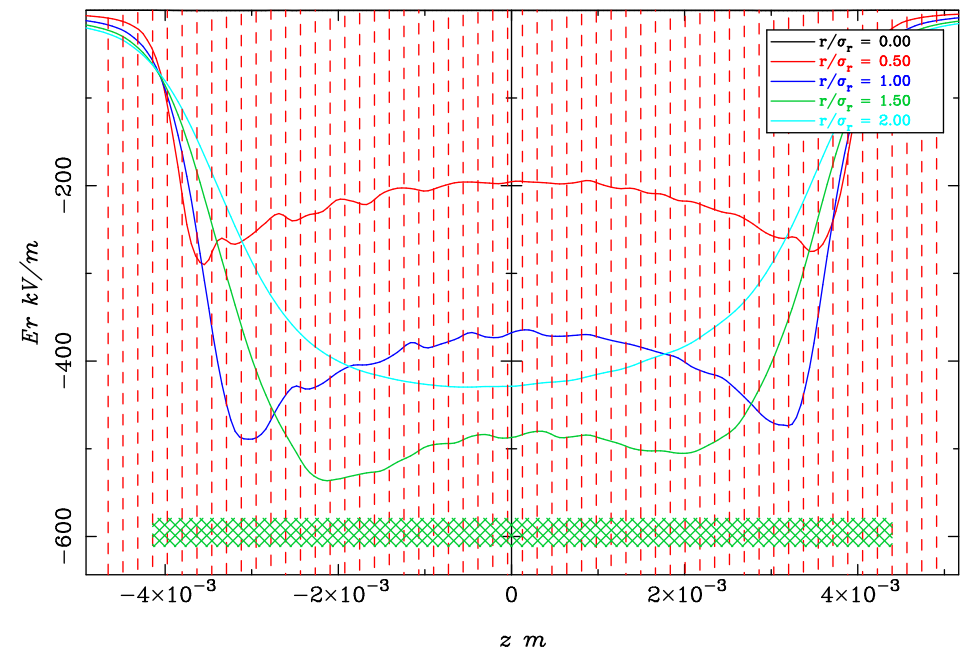
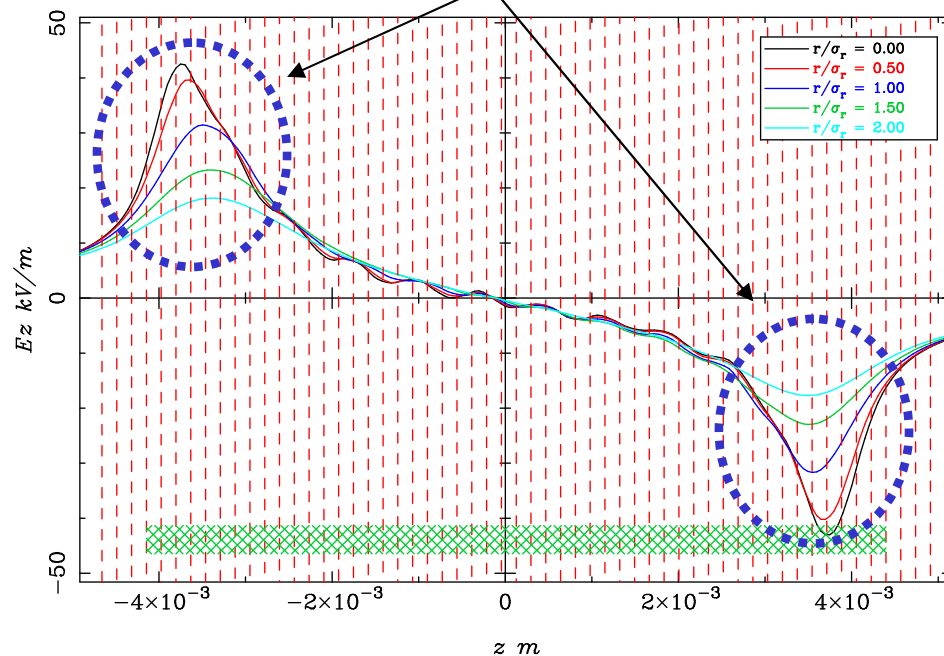
Simulation particles = 50,000

ASTRA Grids for space charge calculation = 10 rings \times 50 cells



↑ Here !

2D space charge force increases the local energy spread at the head & tail regions and helps in smearing of the initial density modulation



ASTRA Tracking with Laser Beam Modulation



At the end of BC2 with CSR

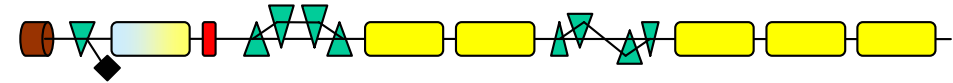
Modulation period = 2.0 ps

Modulation Amplitude = $\pm 10\%$

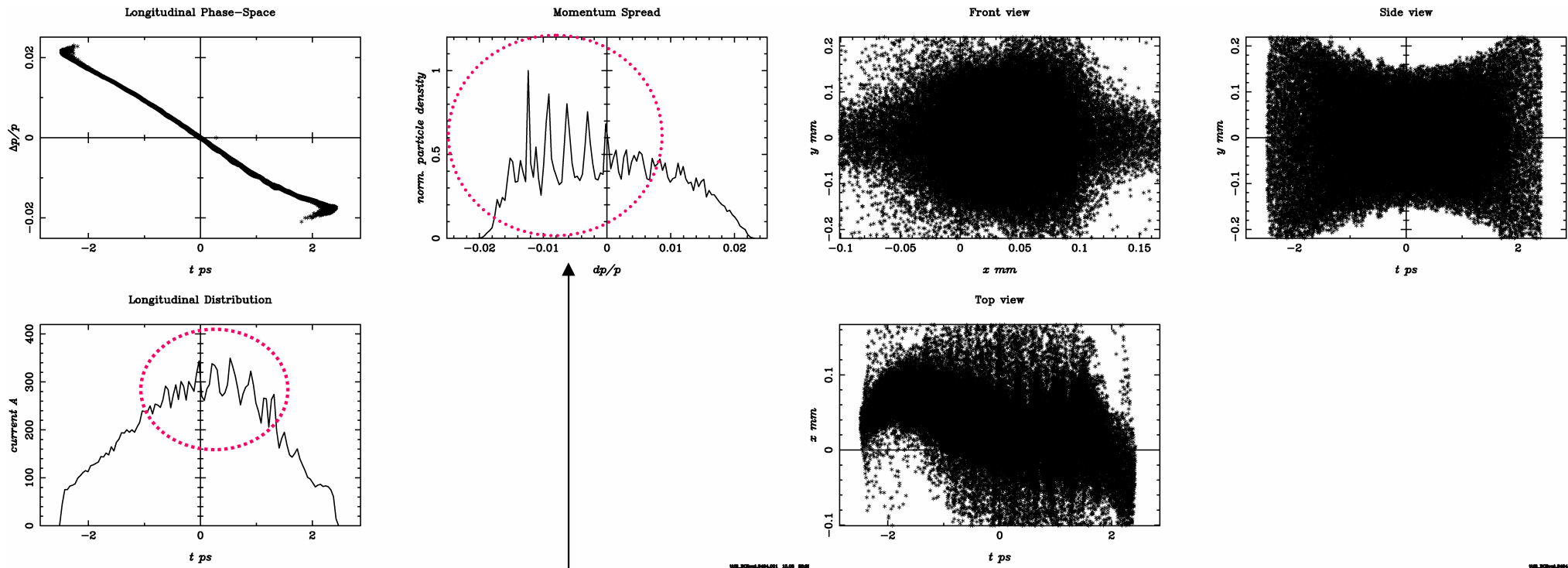
Postpro BIN number = 125

Simulation particles = 50,000

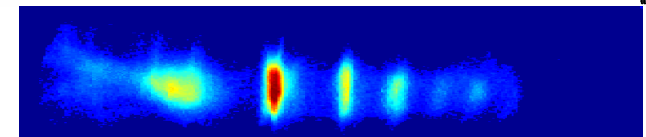
ASTRA Grids for space charge calculation = 10 rings \times 50 cells



↑ Here !



This energy modulation may be observed at SDL
We do not meet a strong density modulation yet !



ASTRA Tracking with Laser Beam Modulation



At the end of BC3

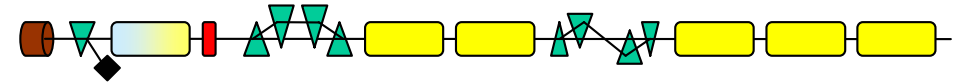
Modulation period = 2.0 ps

Modulation Amplitude = $\pm 10\%$

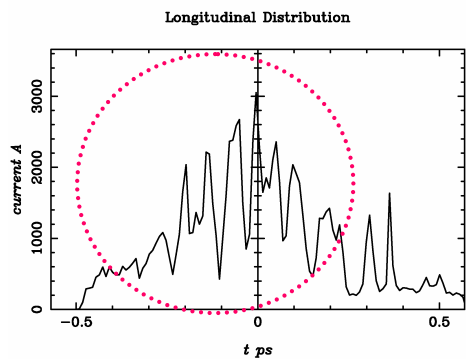
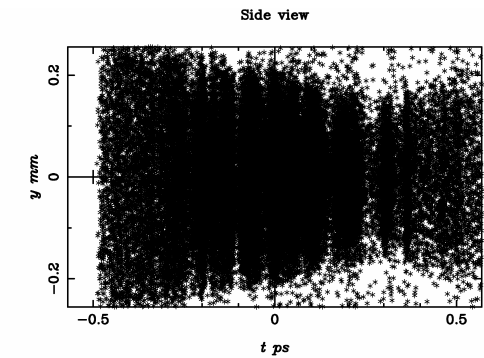
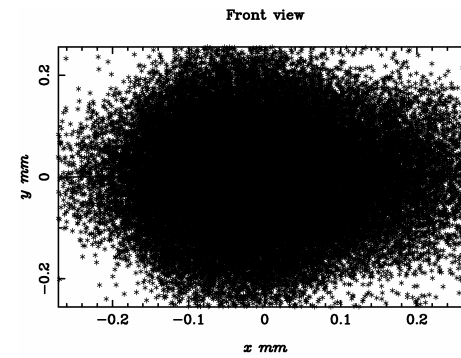
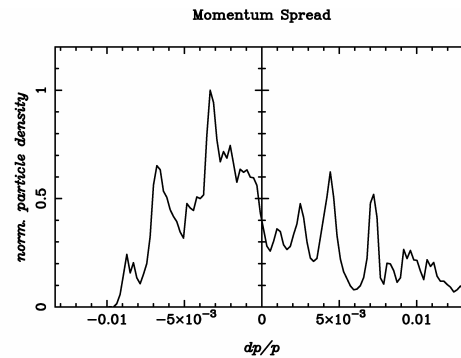
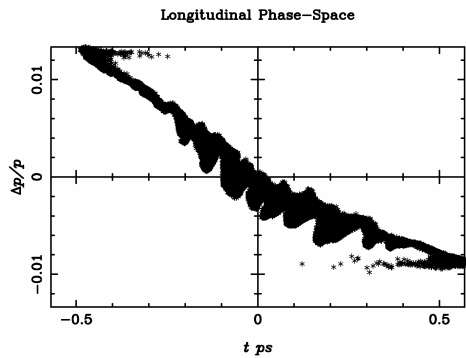
Postpro BIN number = 125

Simulation particles = 50,000

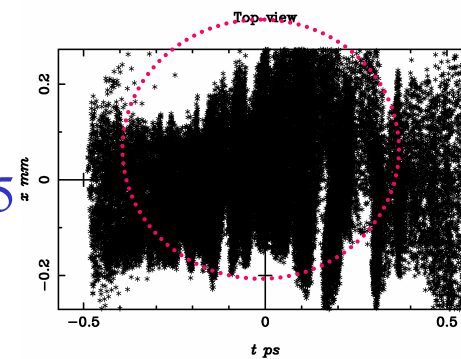
ASTRA Grids for space charge calculation = 10 rings \times 50 cells



↑ Here !



$\pm 50\%$ density modulation !
Total gain from cathode ~ 5



Unfortunately, we meet a density modulation !

Gain after BC3



At the end of BC3

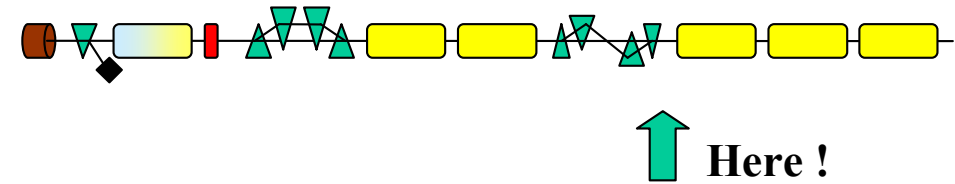
Modulation period = 2.0 ps

Modulation Amplitude = 5%, 10%, 15%, 20%

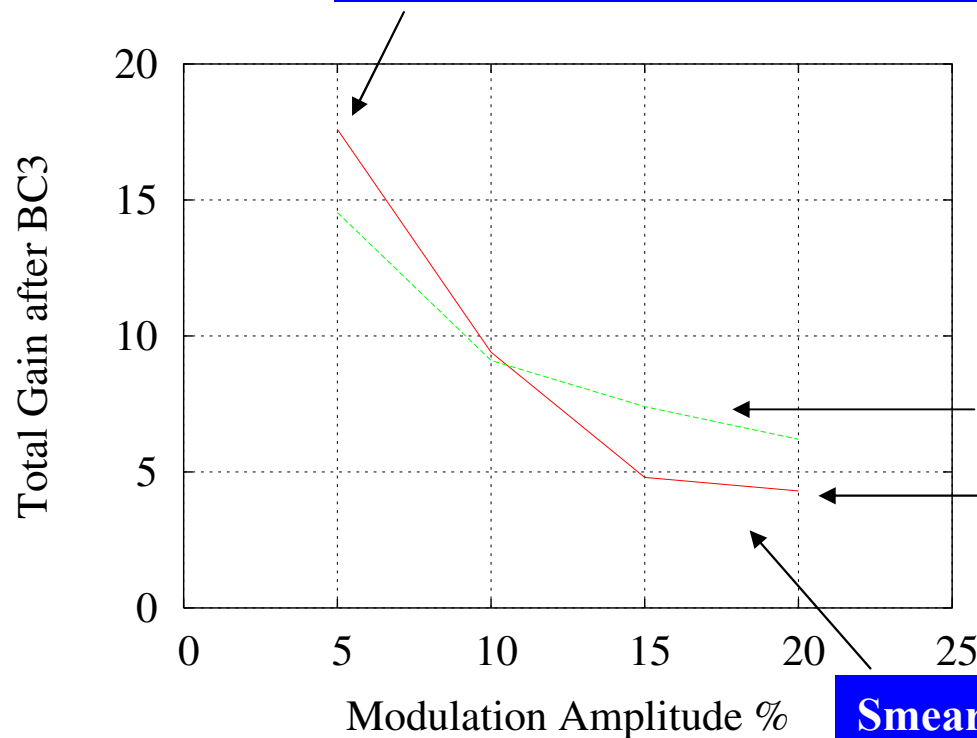
Postpro BIN number = 125

Simulation particles = 50,000

ASTRA Grids for space charge calculation = 10 rings \times 50 cells



Gain is overestimated at small amplitude due to numerical noise



CSR effects becomes strong due to the increased non-uniformity in the current density profile when the modulation amplitude is increased.

Gain ~ 7 without CSR at BC3

Gain ~ 5 with CSR at BC3

Smeared by the increased local energy spread due to CSR wakefield

Slice & Projected Emittances after BC3



At the end of BC3

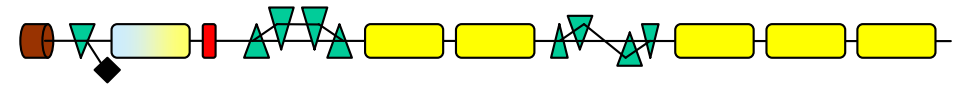
Modulation period = 2.0 ps

Modulation Amplitude = 5%, 10%, 15%, 20%

Postpro BIN number = 125

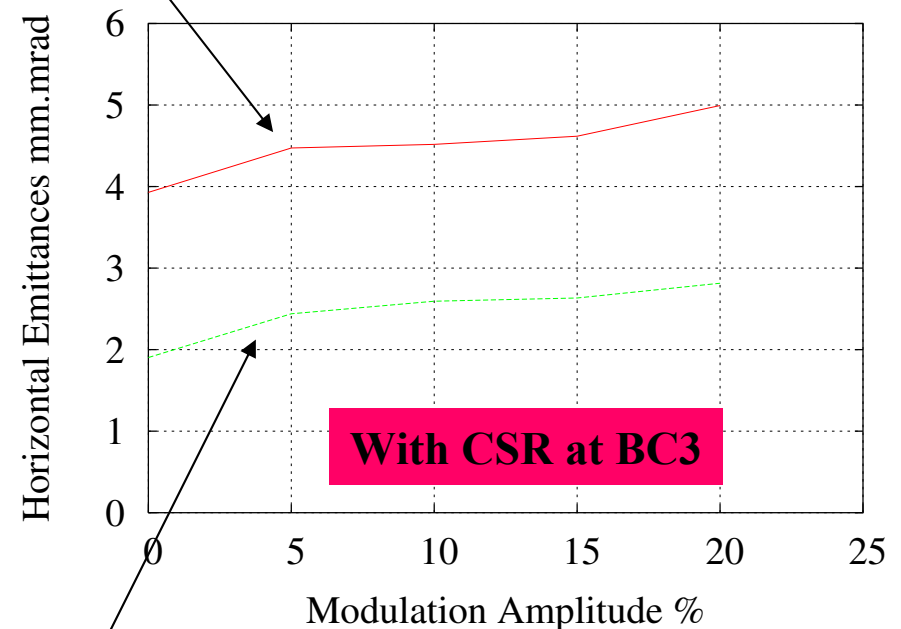
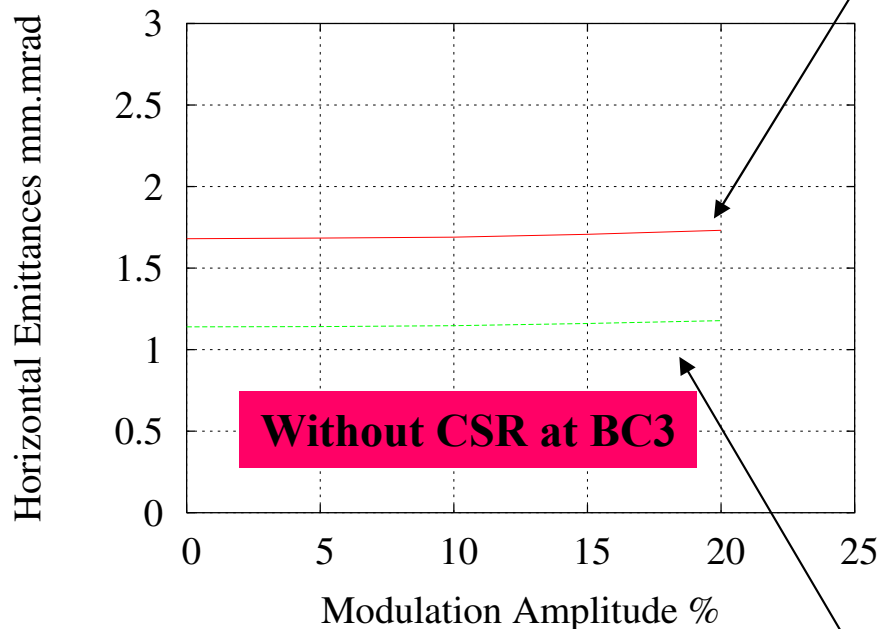
Simulation particles = 50,000

ASTRA Grids for space charge calculation = 10 rings \times 50 cells



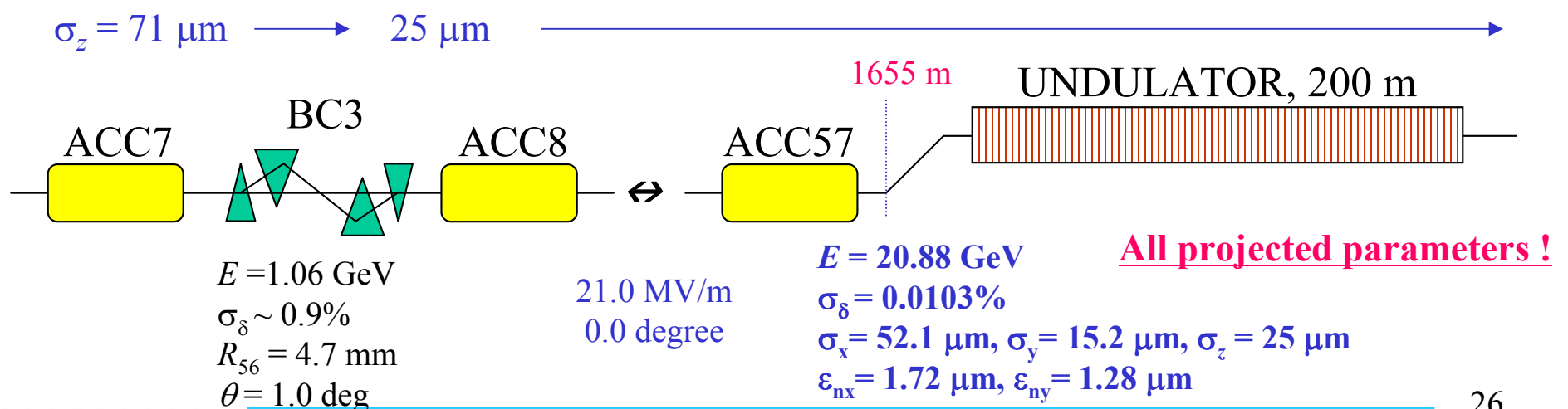
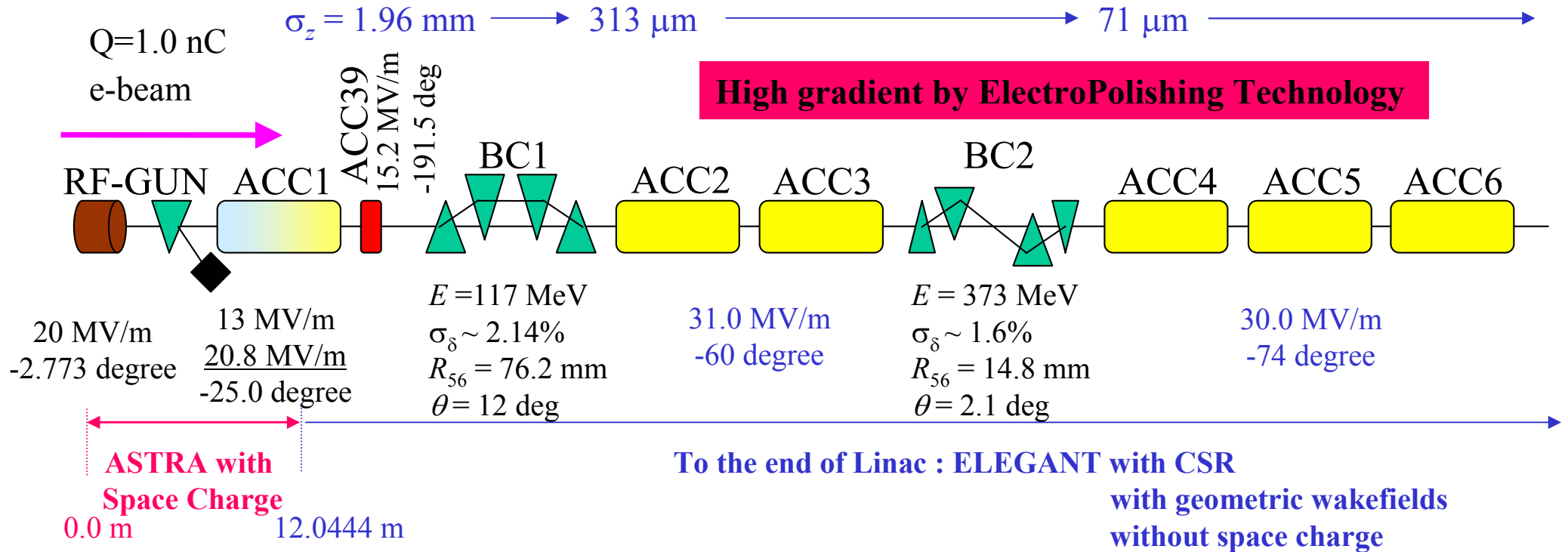
Here !

Normalized projected emittance

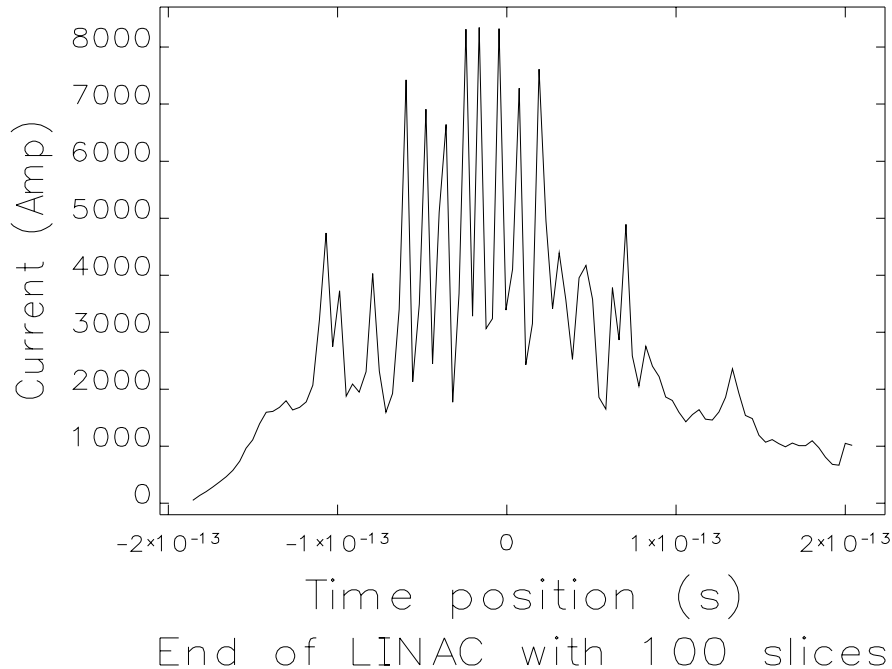


Normalized slice emittance

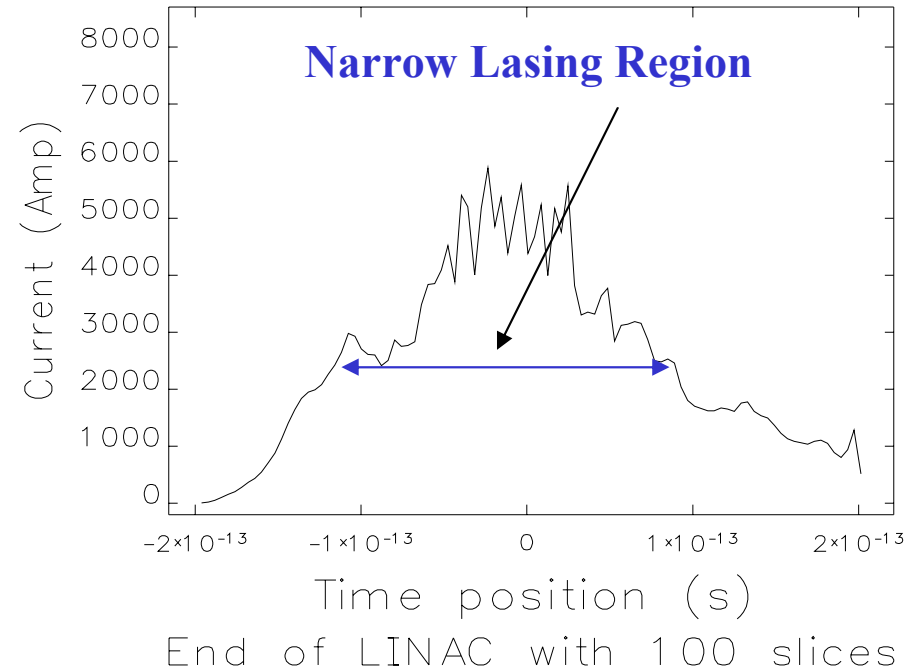
Old Lattice for TESLA XFEL - 2nd Version



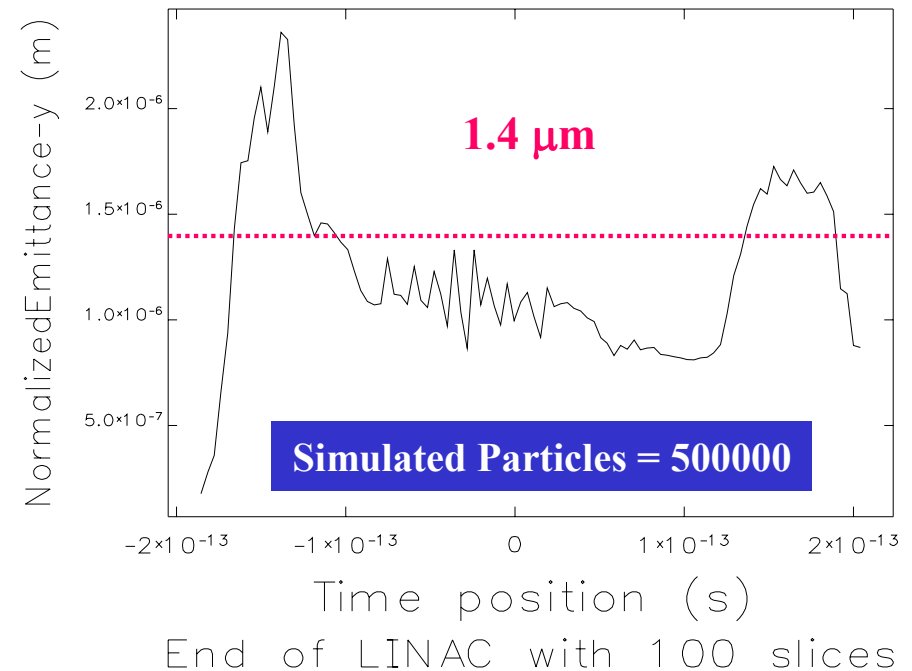
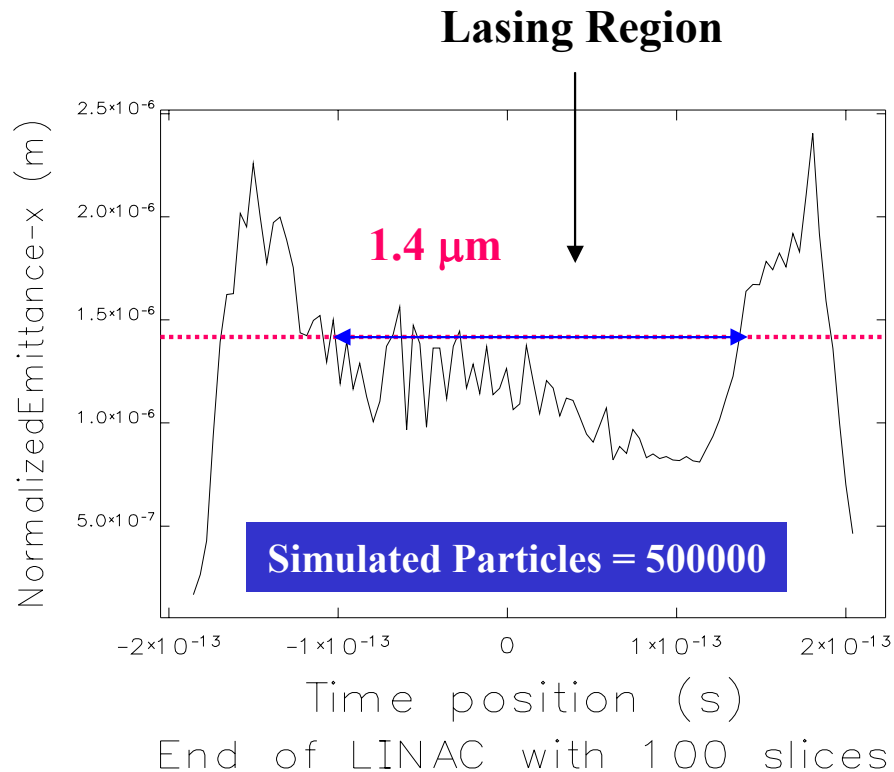
Simulated Particles = 500000



Simulated Particles = 2000000

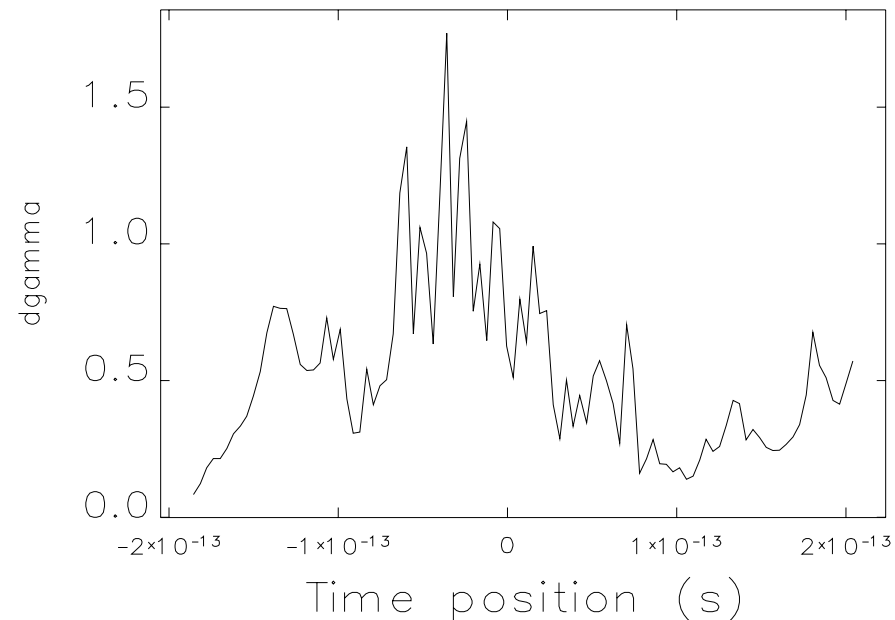


When the number of particles is 500000, there are many high spikes in the current due to the strong microbunching instability at S-type chicanes. But those spikes are somewhat damped when we use 2000000 particles !



When the number of particles is 500000, there are many high spikes in the slice emittance due to the strong microbunching instability at S-type chicanes.

Simulated Particles = 500000



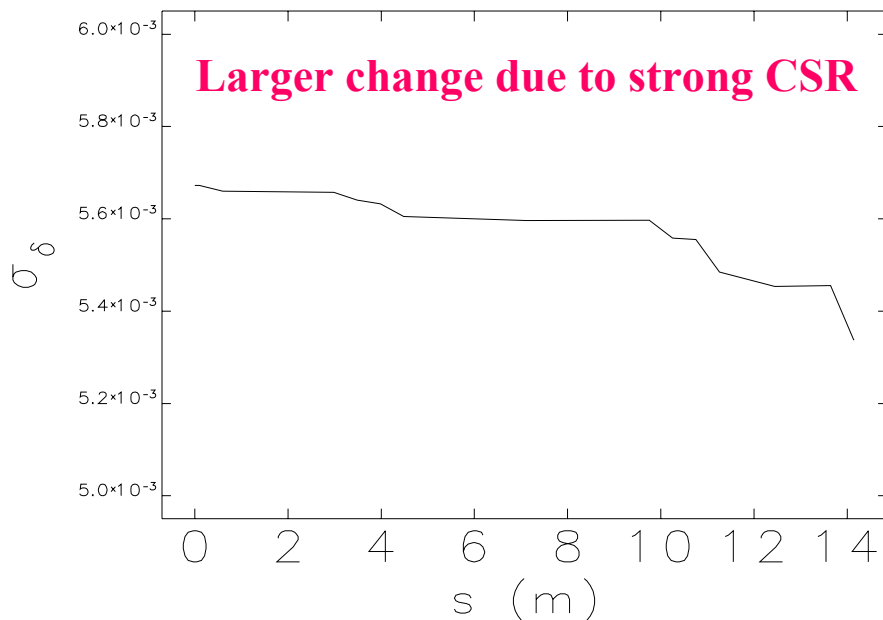
End of LINAC with 100 slices

Maximum uncorrelated energy spread : 0.767 MeV @ 20.0 GeV ~ 0.0038% < 0.0125%
Please note that although the uncorrelated energy spread is not a hot issue at 20 GeV linac, the spread is highest at the center region, i.e., the main lasing region !

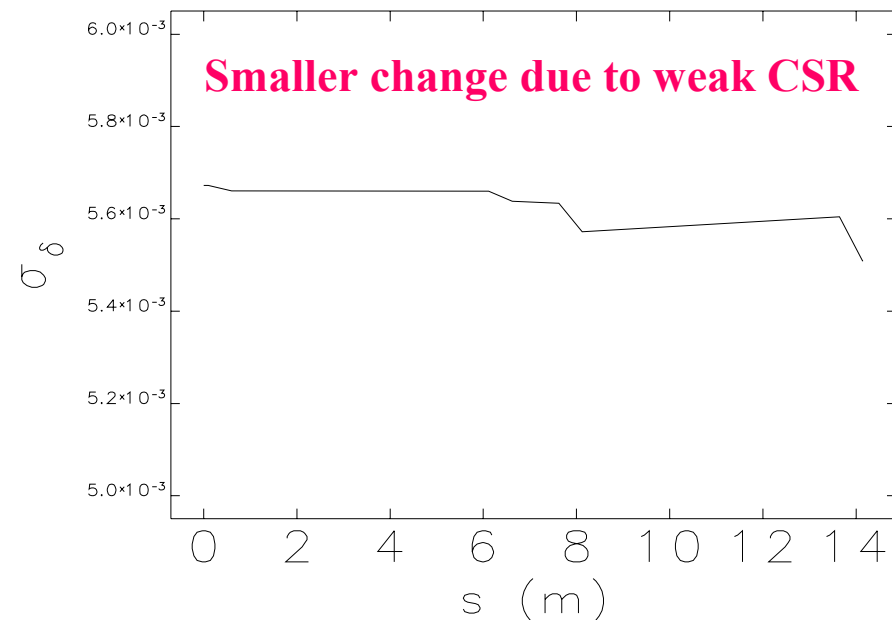
Chicane Comparison - Energy Spread Change



From the TTF-2 S2E Simulations on the microbunching instability, we found that **S-type chicane is more dangerous against the microbunching instability due to two additional dipoles.** When a 2.0 ps 20% current density modulation is applied at the cathode, we got following results at TTF-2 BC3. (Refer to Yujong's TESLA-S2E-2003-08 & TESLA-S2E-2003-10)



Along BC3 with 2.0 ps 20% modulation and with CSR



Along BC3 with 2.0 ps 20% modulation and with CSR

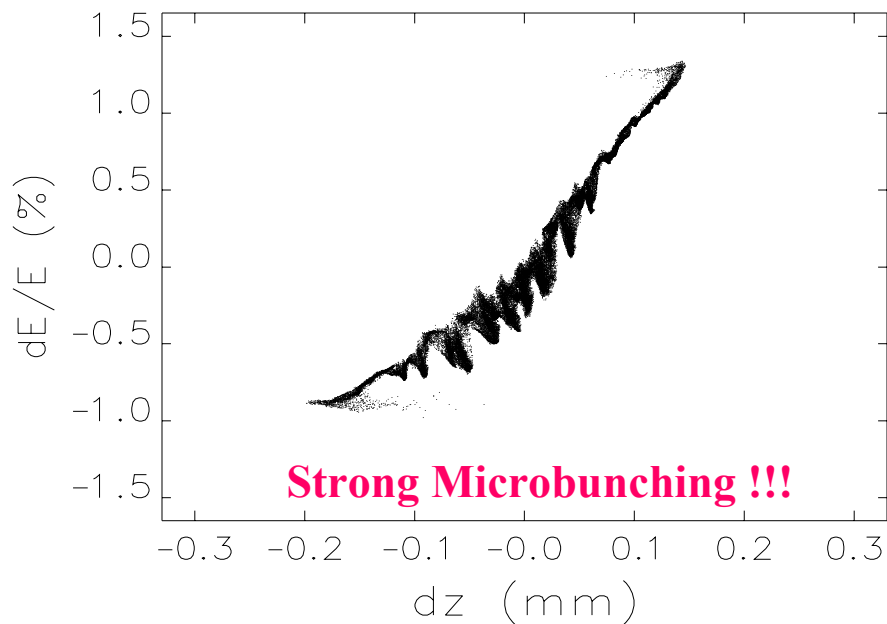
$$\sigma_{\delta\Delta} = 0.567 \sim 0.53\% \text{ for S-type chicane}$$

$$\sigma_{\delta\Delta} = 0.567 \sim 0.55\% \text{ for 4 bend chicane}$$

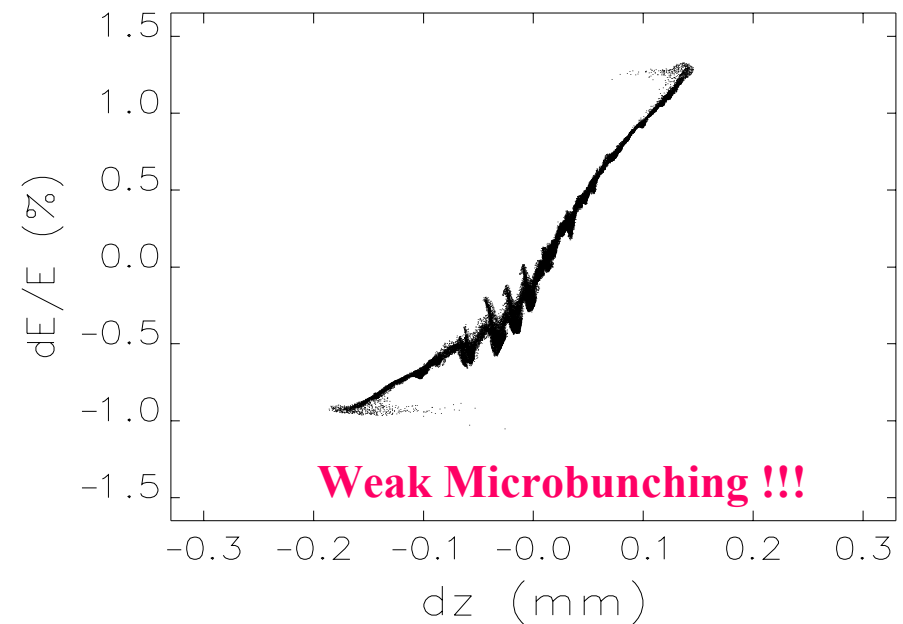
Chicane Comparison - Longitudinal Space



From the TTF-2 S2E Simulations on the microbunching instability, we found that **S-type chicane is more dangerous against the microbunching instability due to two additional dipoles.** When a 2.0 ps 20% current density modulation is applied at the cathode, we got following results at TTF-2 BC3. (Refer to Yujong's TESLA-S2E-2003-08 & TESLA-S2E-2003-10)



6-bends S-type chicane

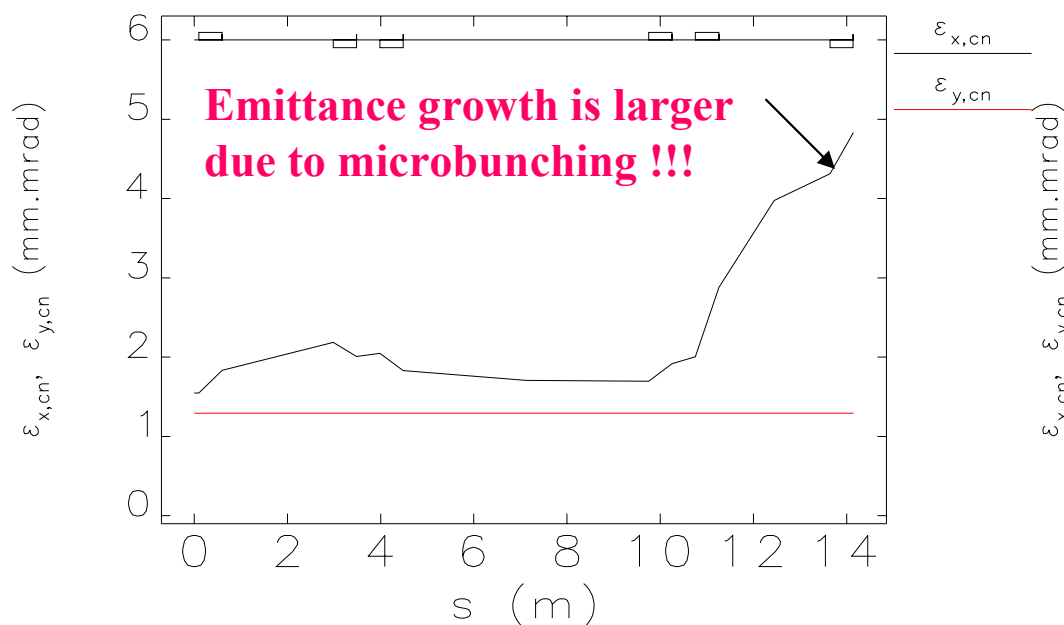


4-bends normal chicane

Chicane Comparison - Project emittance

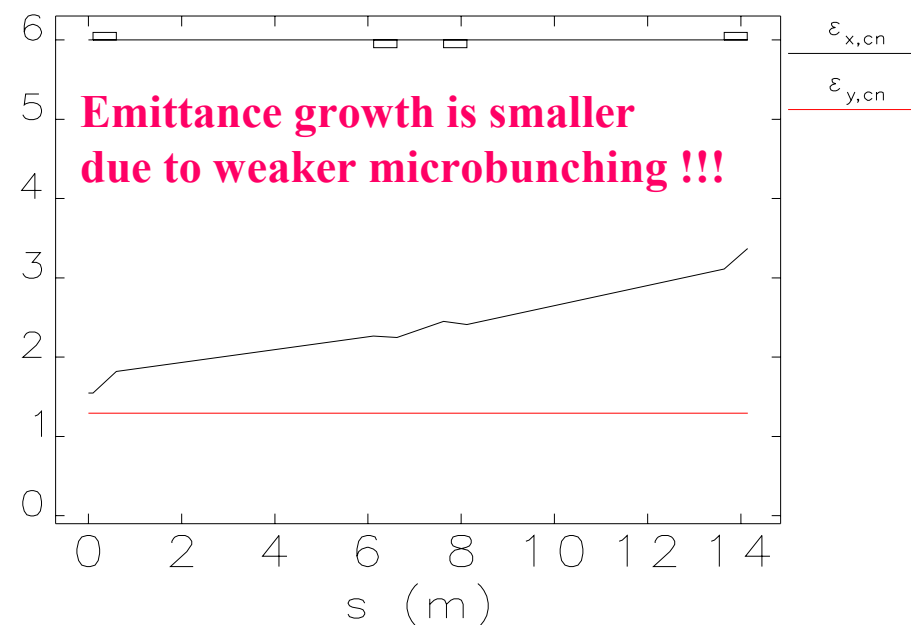


From the TTF-2 S2E Simulations on the microbunching instability, we found that **S-type chicane is more dangerous against the microbunching instability due to two additional dipoles.** When a 2.0 ps 20% current density modulation is applied at the cathode, we got following results at TTF-2 BC3. (Refer to Yujong's TESLA-S2E-2003-08 & TESLA-S2E-2003-10)



After BC3 with 2.0 ps 20% modulation and with CSR

$$\epsilon_{nx} = 4.82 \mu\text{m for S-type chicane}$$



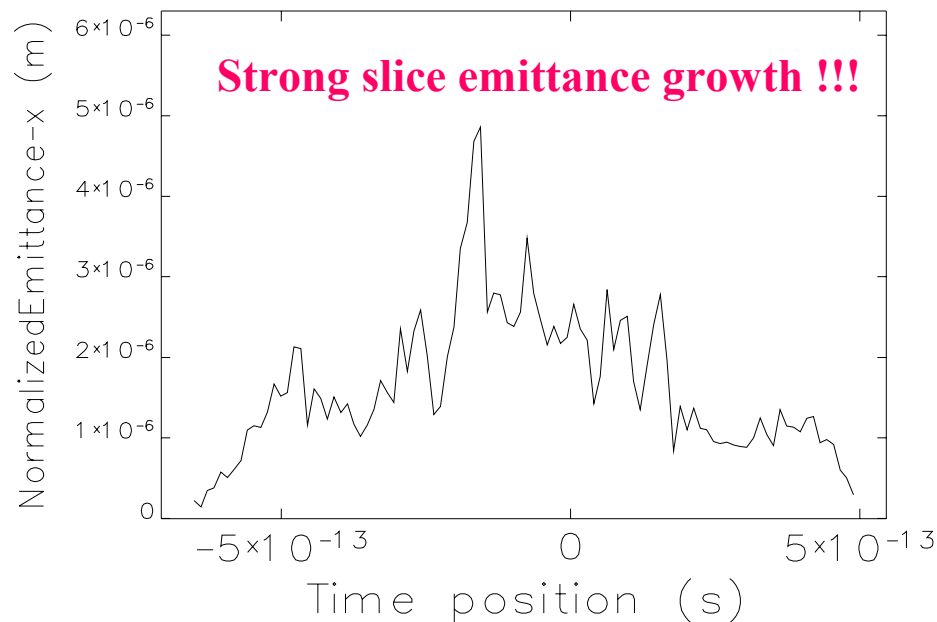
After BC3 with 2.0 ps 20% modulation and with CSR

$$\epsilon_{nx} = 3.36 \mu\text{m for 4 bend chicane}$$

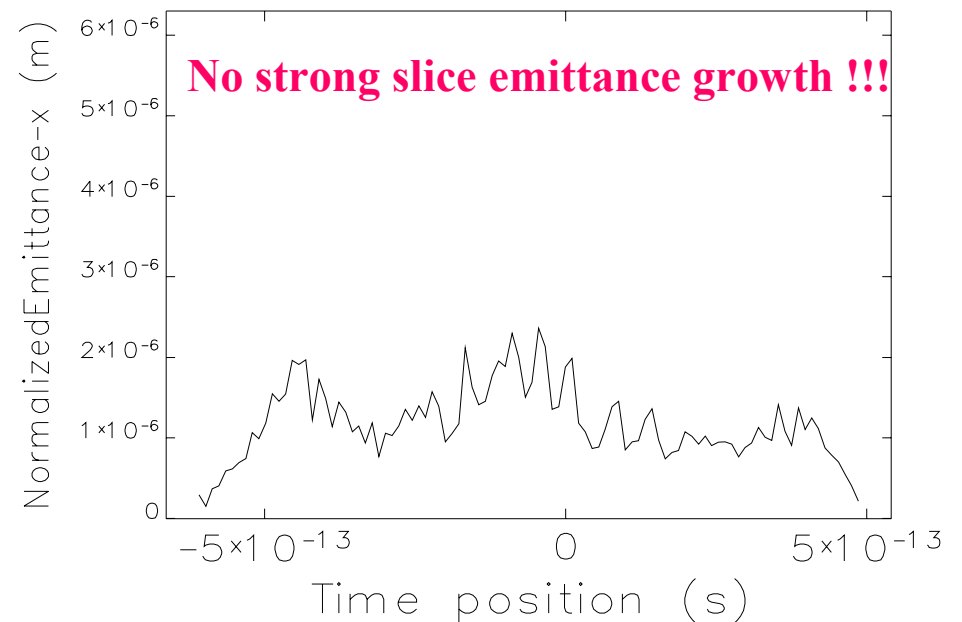
Chicane Comparison - Slice emittance



From the TTF-2 S2E Simulations on the microbunching instability, we found that **S-type chicane is more dangerous against the microbunching instability due to two additional dipoles.** When a 2.0 ps 20% current density modulation is applied at the cathode, we got following results at TTF-2 BC3. (Refer to Yujong's TESLA-S2E-2003-08 & TESLA-S2E-2003-10)



6-bends S-type chicane



4-bends normal chicane



Generally the compensation of S-type chicane helps in reducing the projected emittance growth due to CSR. However the projected emittance growth in the 4-bend chicane can be small enough if we well-optimize the normal chicane.

❑ Reducing projected emittance growth due to CSR & Chromatic Effects

- Keeping somewhat large energy spread at BCs by putting BCs at a lower energy
- Reducing QM length around BCs to reduce chromatic effect
- Removing emittance measuring FODO after BC1 to reduce chromatic effect
- Using weaker strength chicanes with a longer chicane length of ~ 20 m
- Using a longer drift space between 1st and 2nd dipoles in BCs ~ 8.9 m
- Optimizing Twiss parameters around BCs
- Selecting much smaller compression factor at BC2 against CSR there
- Reducing dipole length from 0.5 m to 0.3 m against CSR

❑ Reducing Slice Parameter Growths due to the Microbunching Instability

- Replace two 6-bends S-type chicanes into one 4-bends normal chicane
- Reducing three BC stages to one BC stage with a double chicane
- Smearing the microbunching instability at BC2 by keeping large uncorrelated energy spread at BC2 without any acceleration between BC1 and BC2

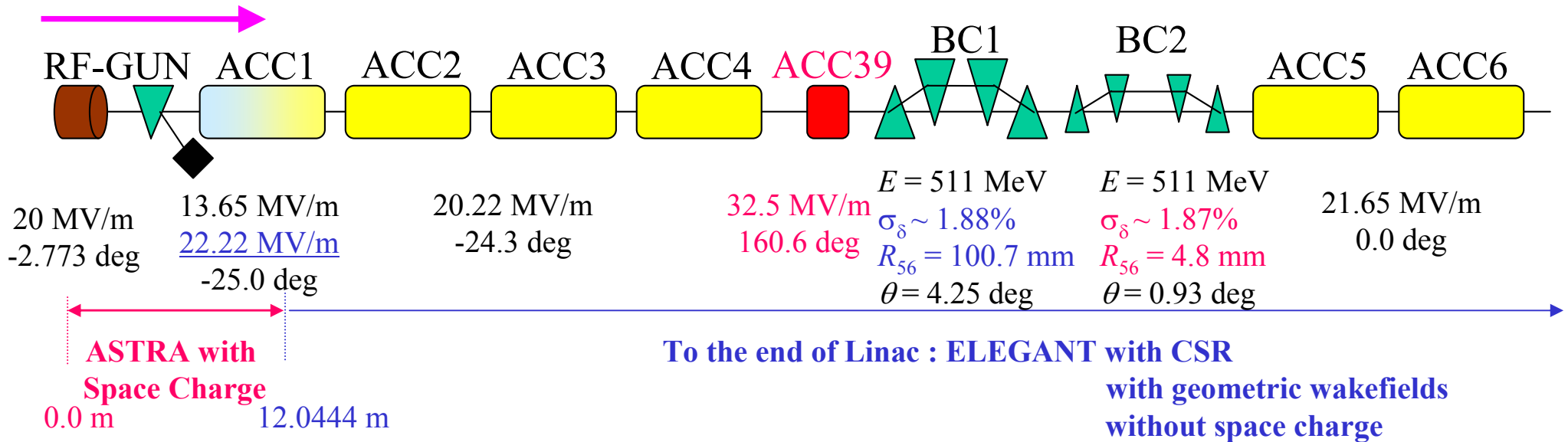
New Lattice for TESLA XFEL - 3rd Version



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

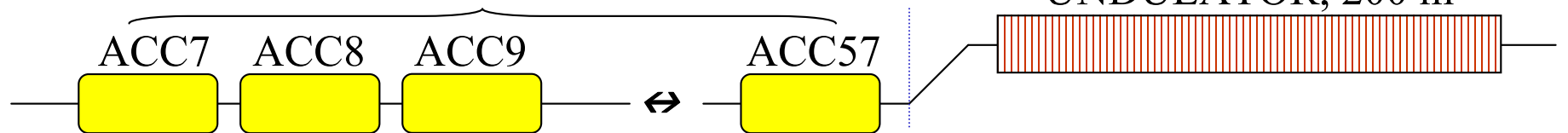
$Q = 1.0 \text{ nC}$
e-beam

$\sigma_z = 2.0 \text{ mm} \longrightarrow 112 \mu\text{m} \longrightarrow 22 \mu\text{m}$



$\sigma_z = 20 \mu\text{m}$

FODO MODULES



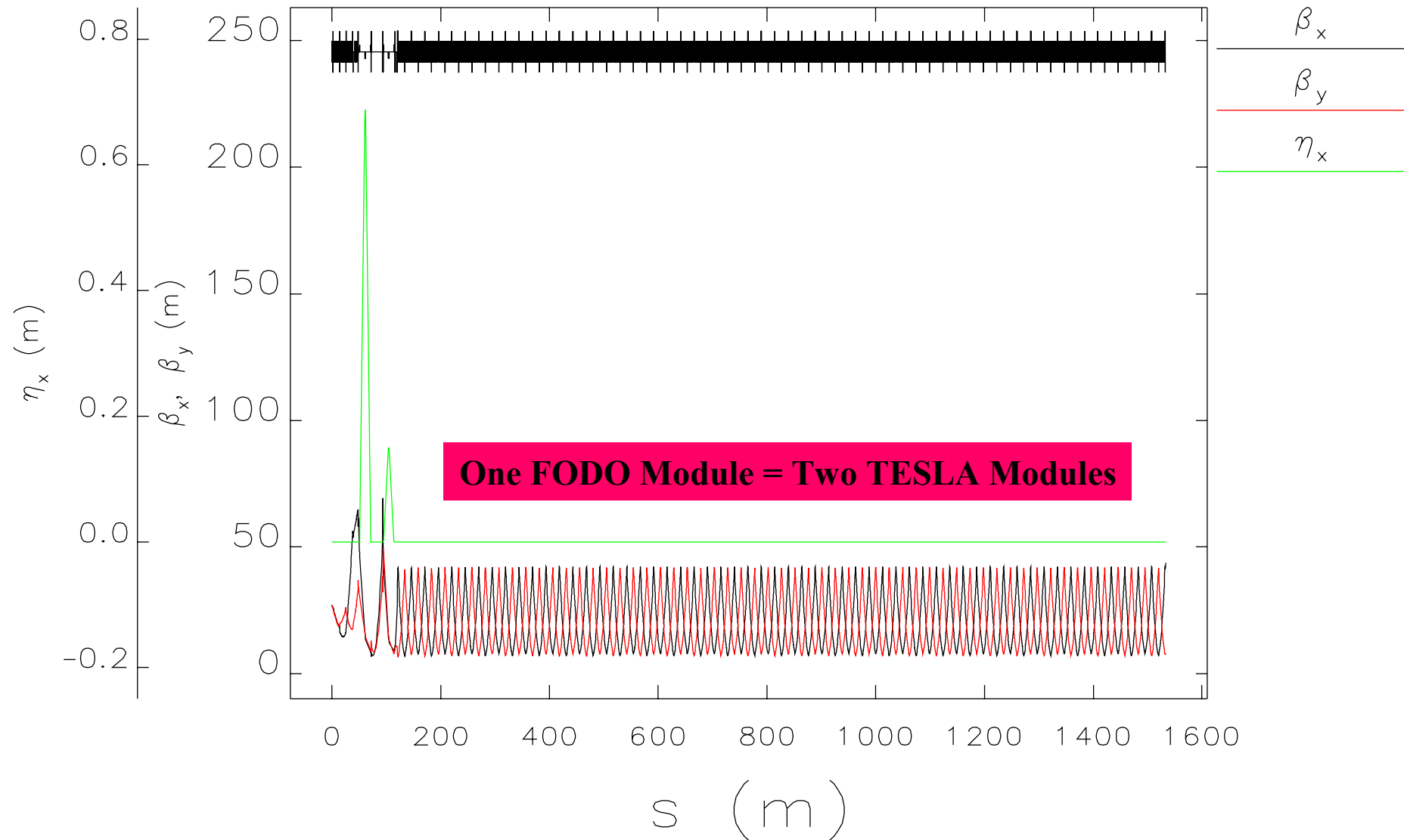
21.65 MV/m
0.0 deg

$E = 20.0 \text{ GeV}$ **All projected parameters !**
 $\sigma_\delta = 0.0059\%$
 $\sigma_x = 41.7 \mu\text{m}, \sigma_y = 15.8 \mu\text{m}, \sigma_z = 20 \mu\text{m}$
 $\epsilon_{nx} = 1.525 \mu\text{m}, \epsilon_{ny} = 1.295 \mu\text{m}$

Twiss Parameters of New Lattice



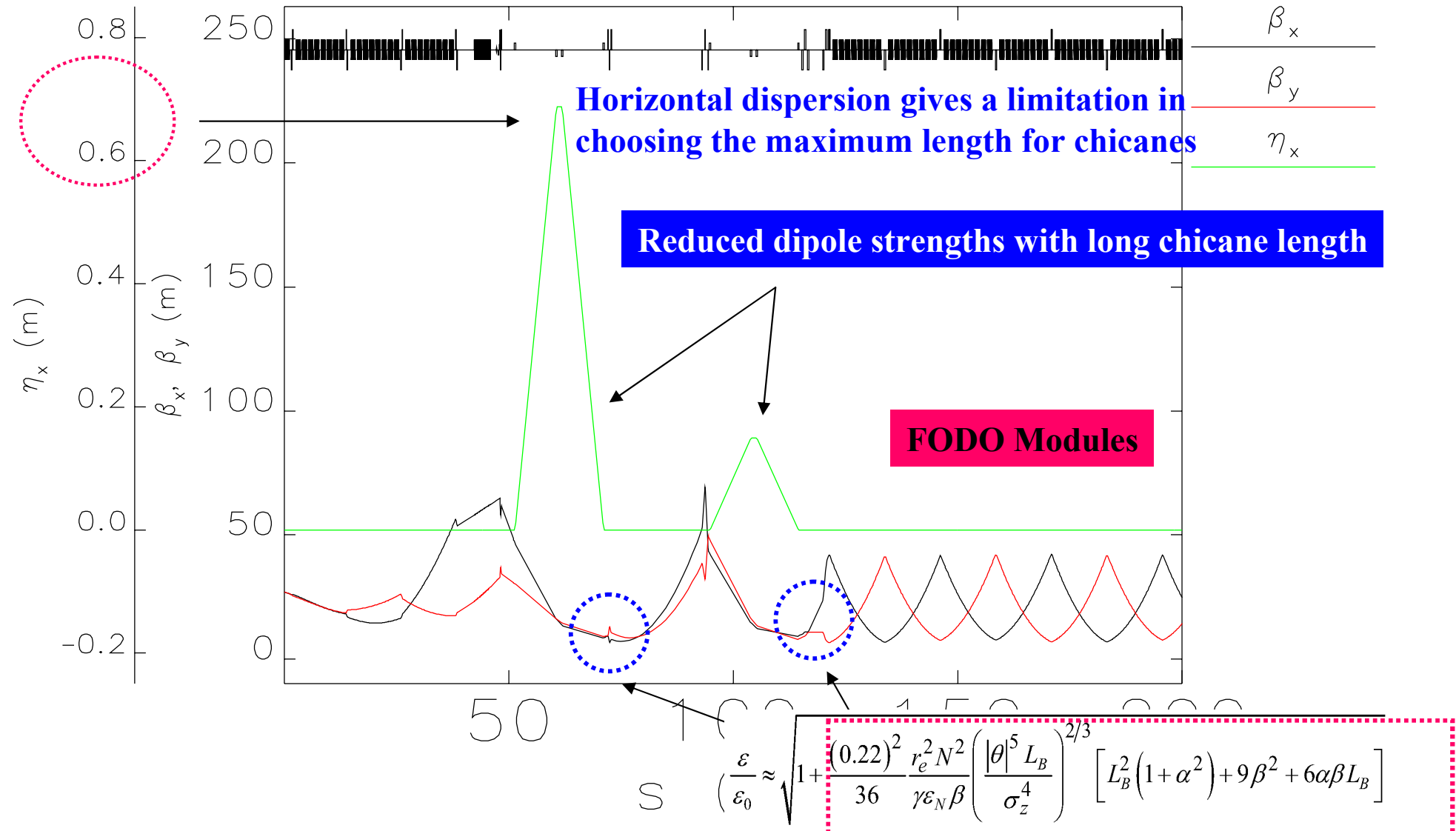
With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$



Twiss Parameters of New Lattice around BCs



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$



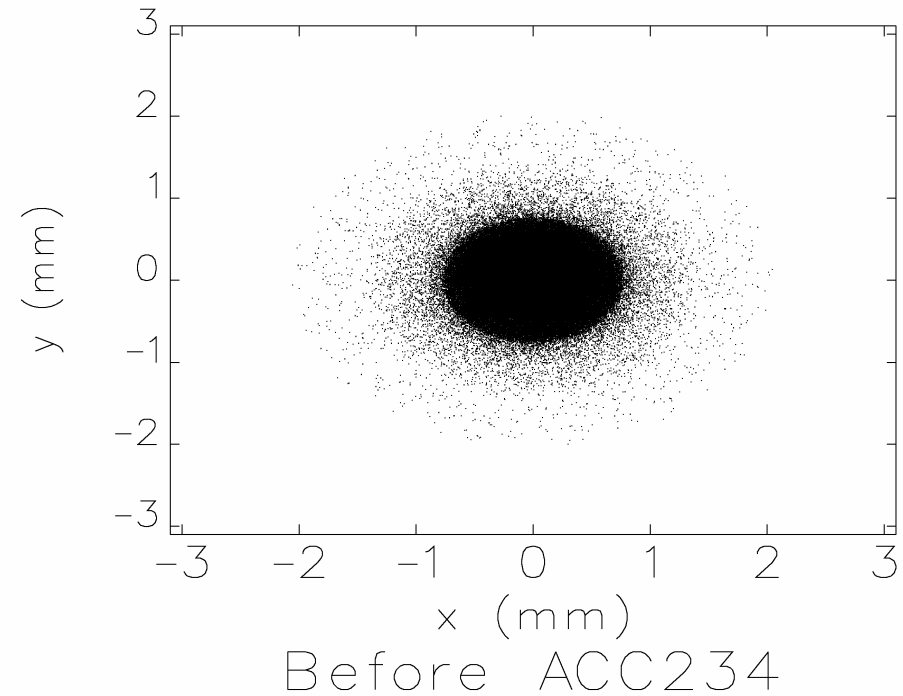
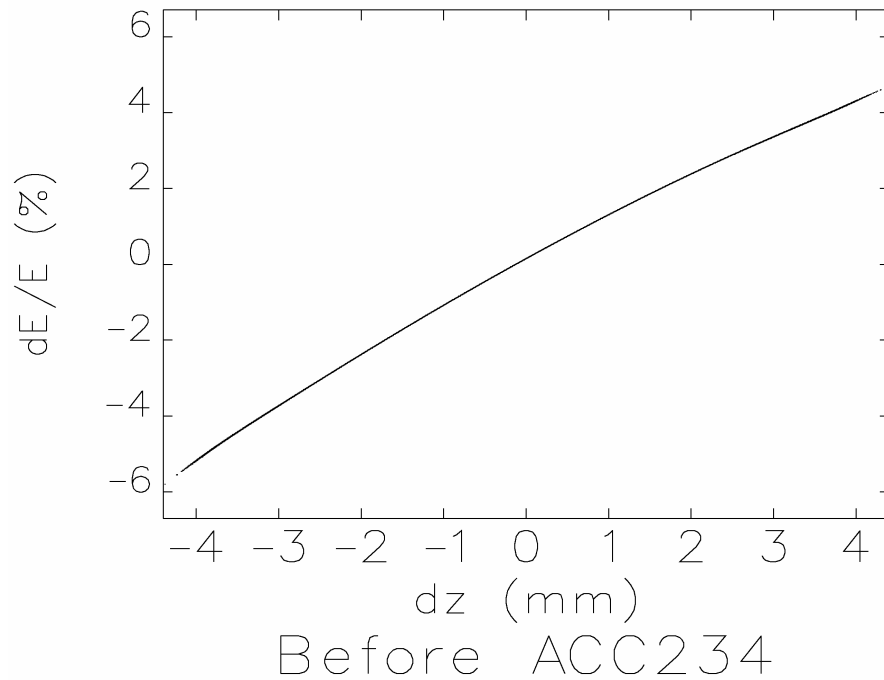
Lattice optimization to minimize this term

New Lattice - Before ACC2



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

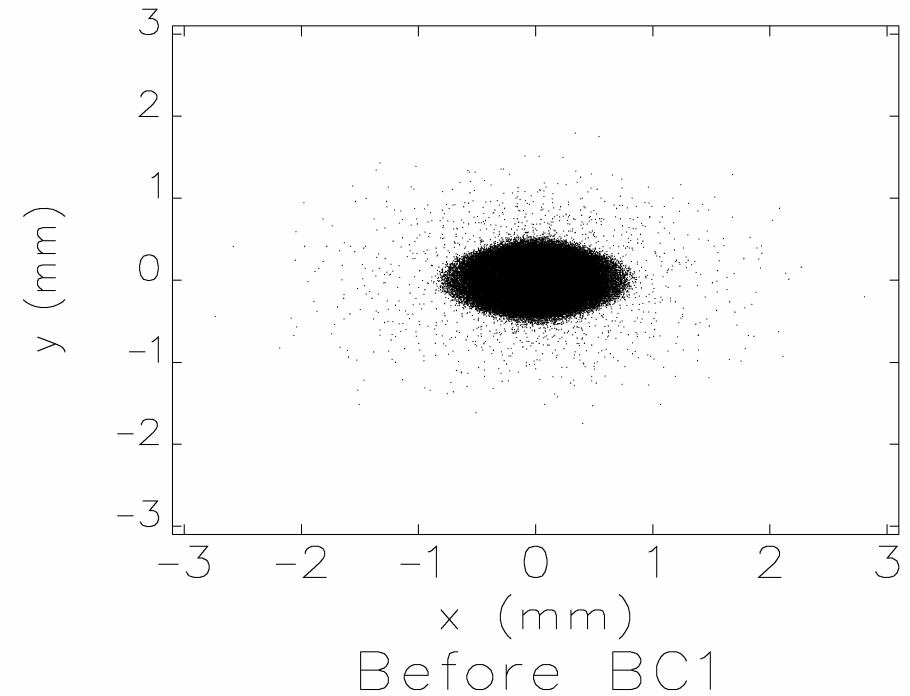
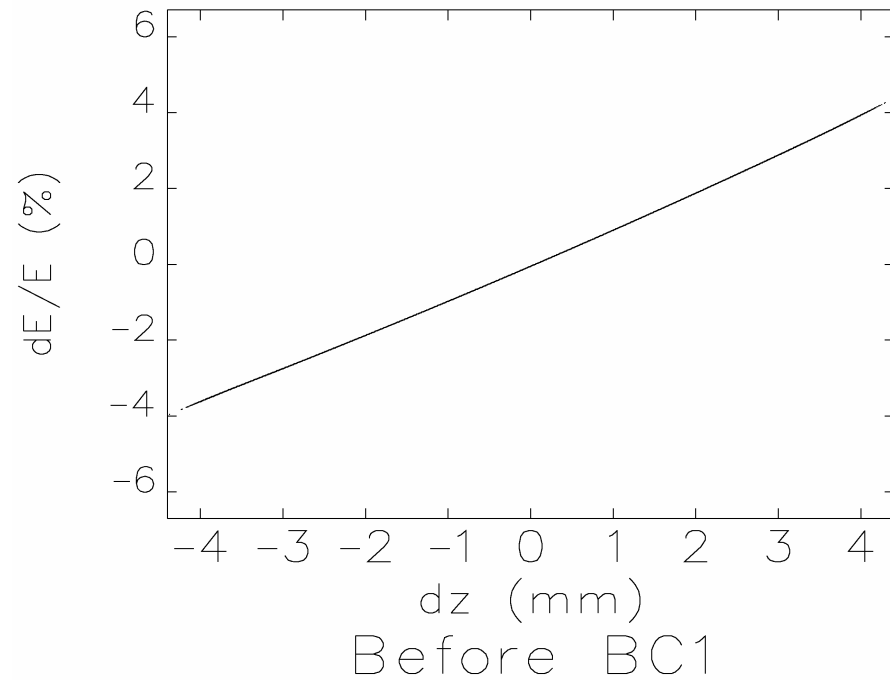


New Lattice - Before BC1



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

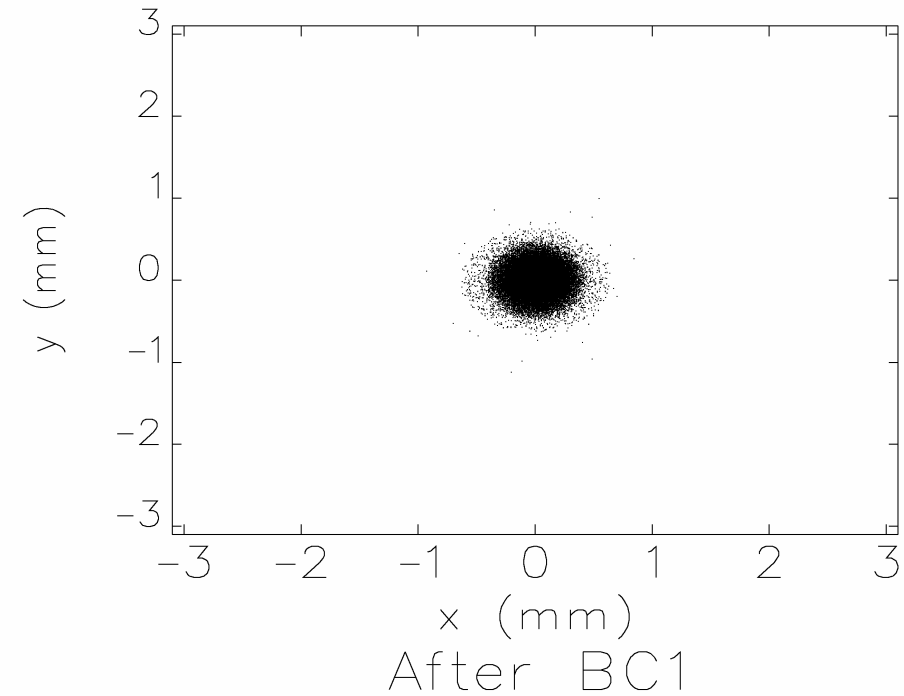
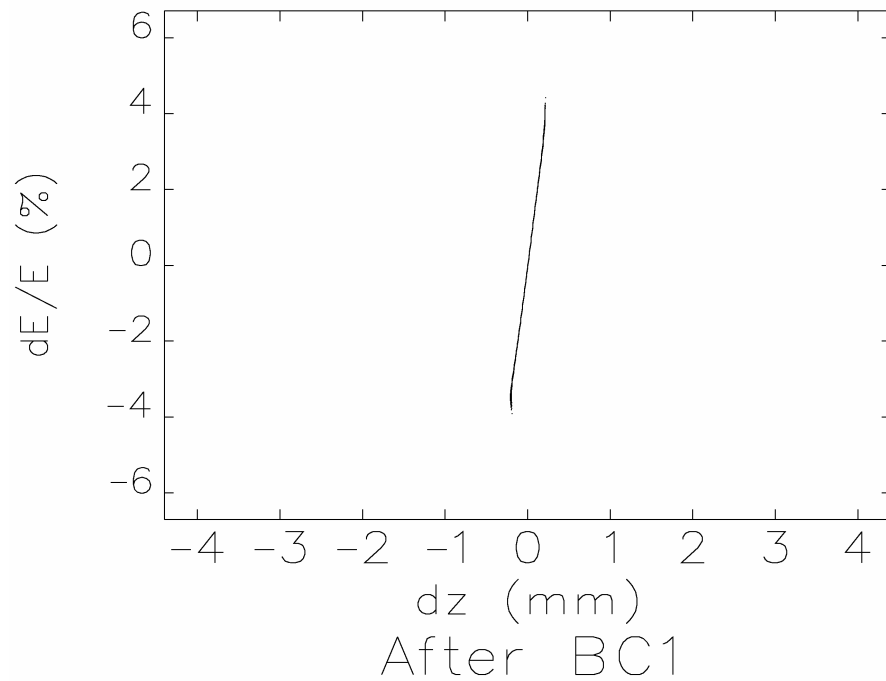


New Lattice - After BC1



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

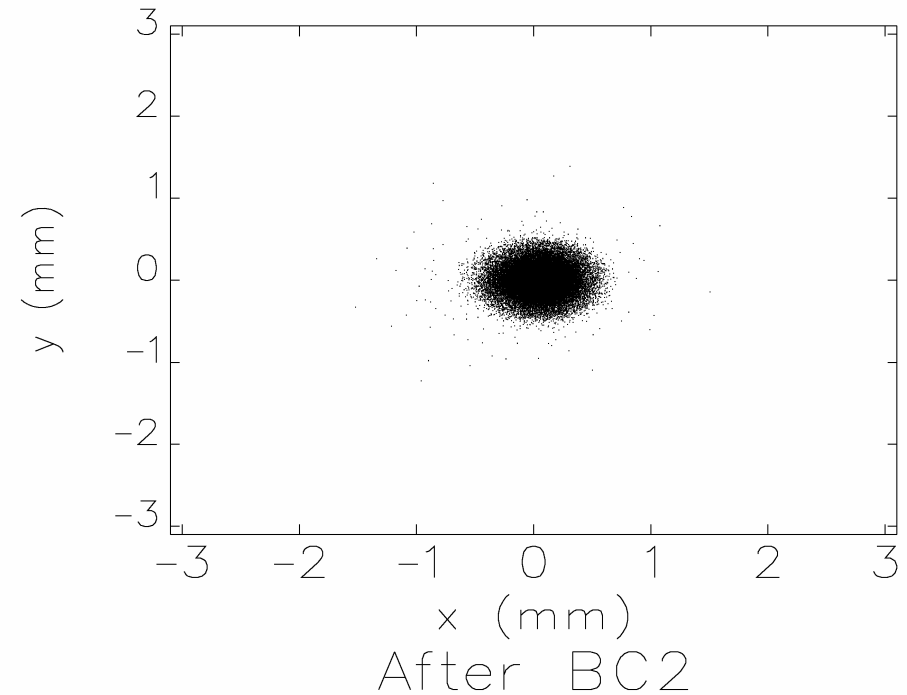
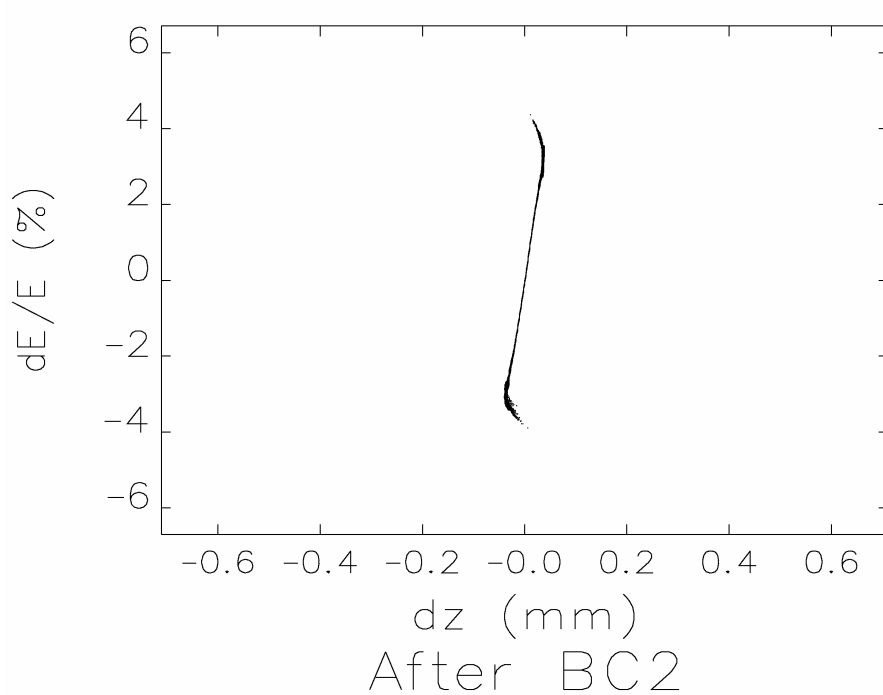


New Lattice - After BC2



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



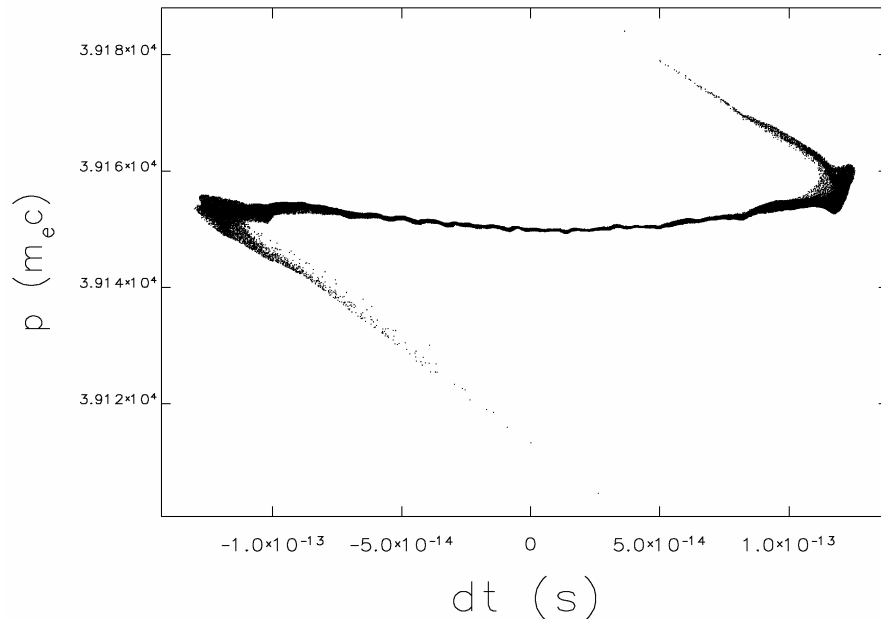
Although we use only 200000 particles, we did not meet any strong artificial modulation after BC2. This means that the amplification action of our new lattice is much weaker than that of our old lattice.

New Lattice - At the end of LINAC

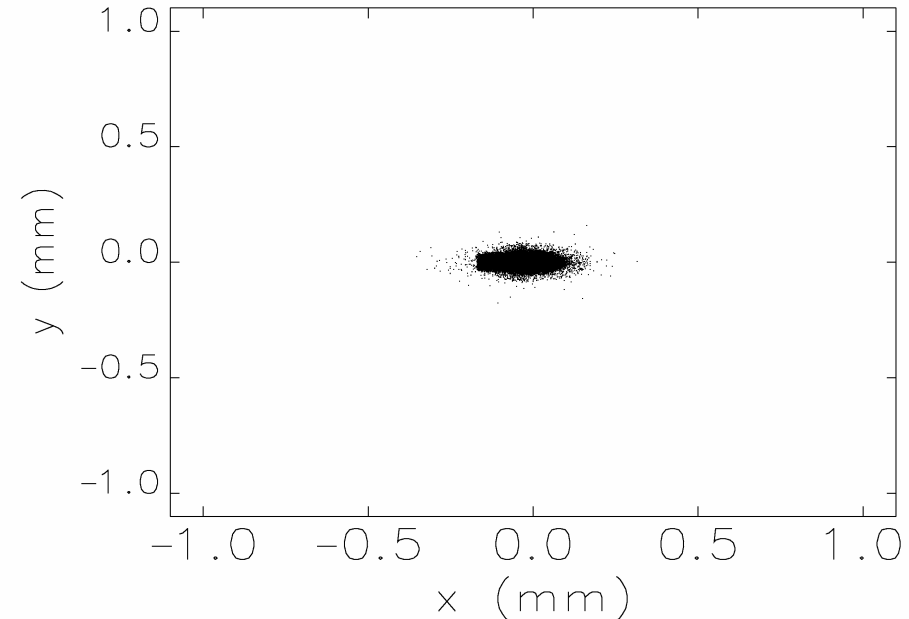


With TTF2 Injector, $\varepsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



dt (s)
End of LINAC



x (mm)
End of LINAC

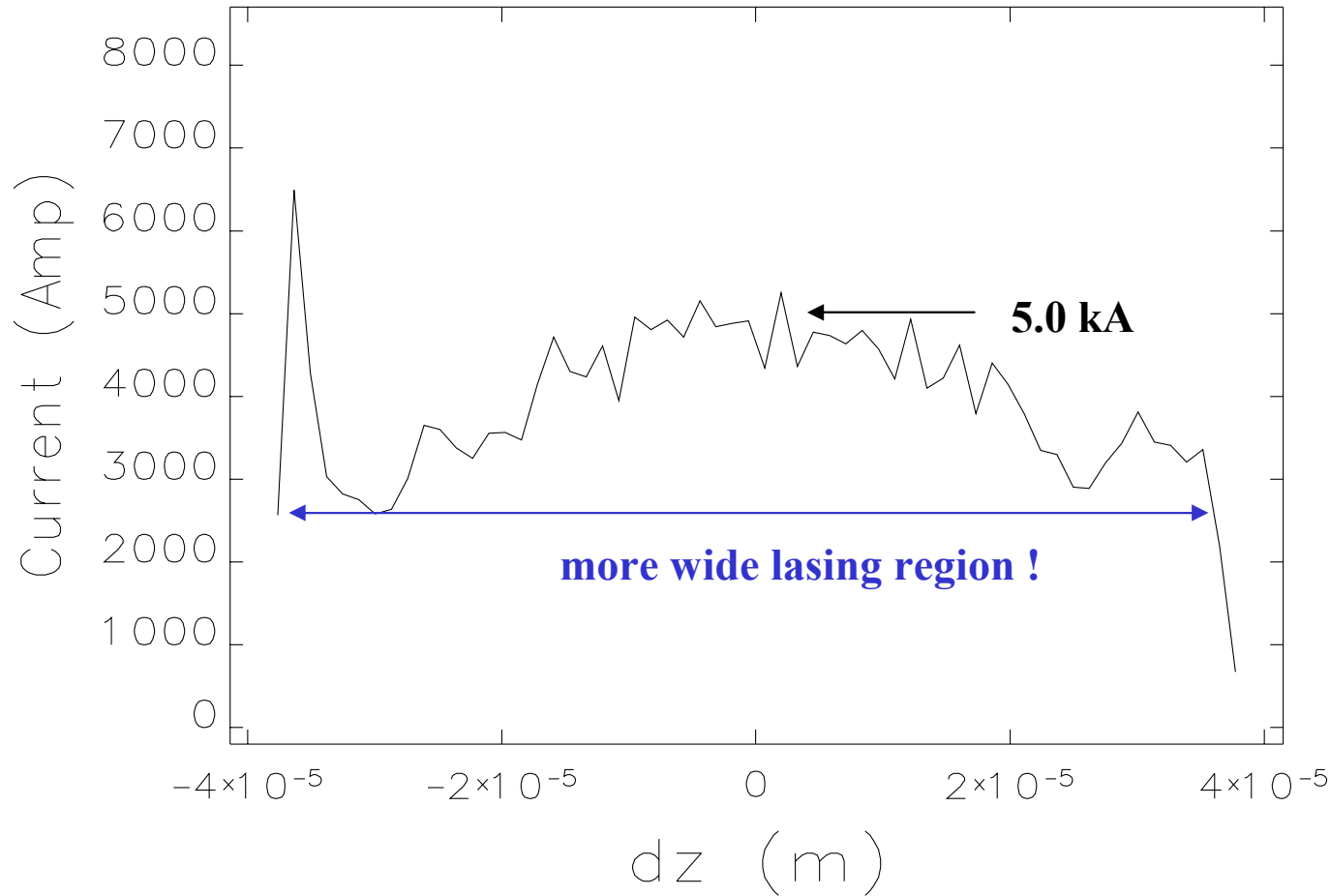
Although we use only 200000 particles, we did not meet any strong artificial modulation after BC2. This means that the amplification action of our new lattice is much weaker than that of our old lattice.

At the end of LINAC - Slice Parameters



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



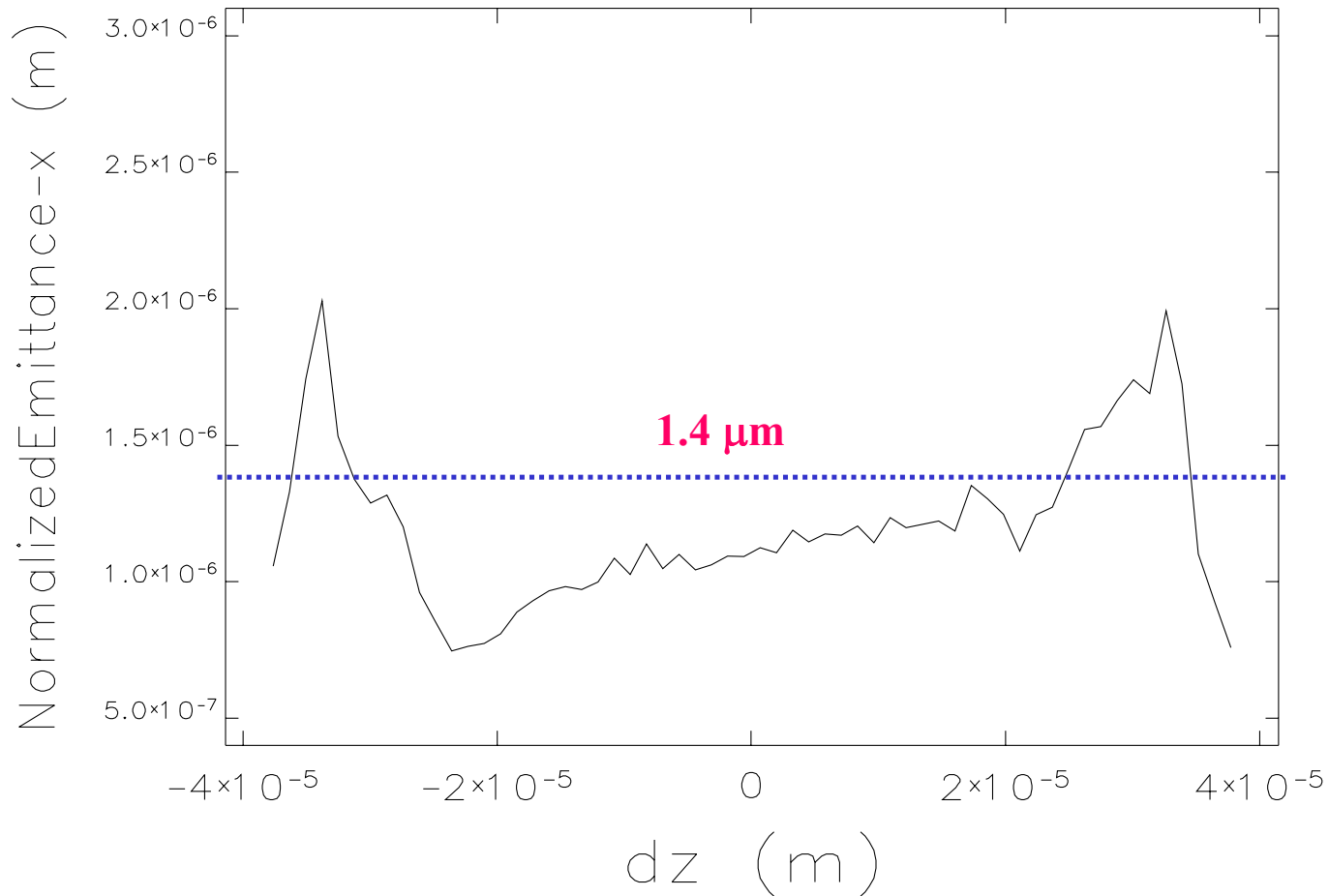
END of LINAC with 60 slices

At the end of LINAC - Slice Parameters



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



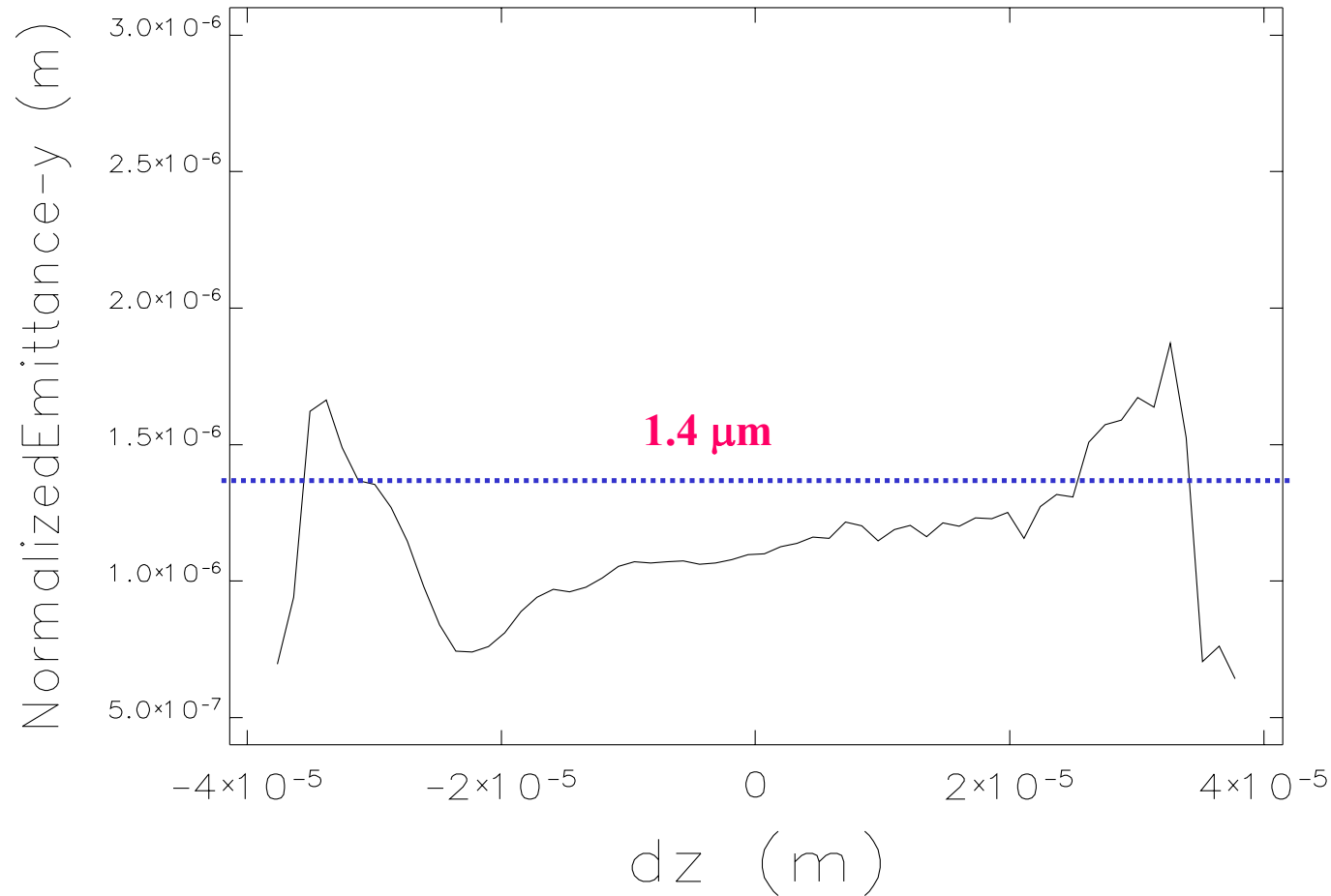
END of LINAC with 60 slices

At the end of LINAC - Slice Parameters



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



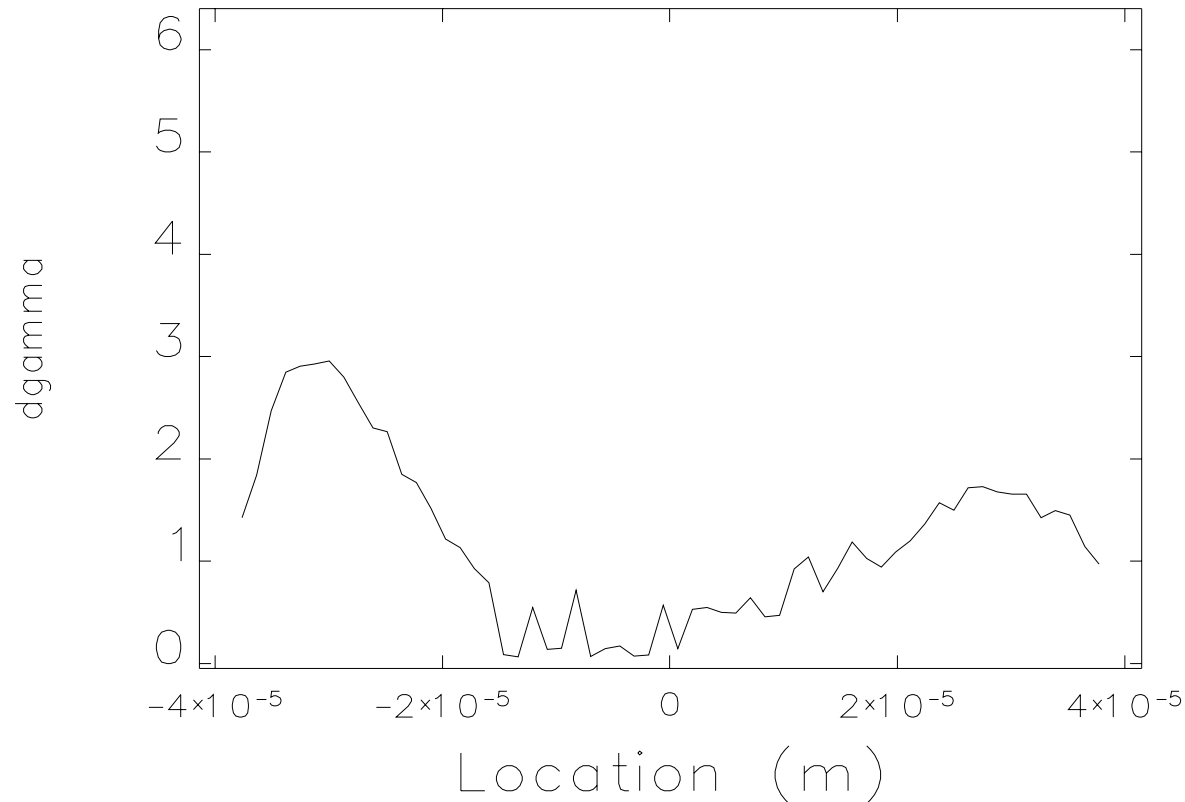
END of LINAC with 60 slices

At the end of LINAC - Slice Parameters



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000



END of LINAC with 60 slices

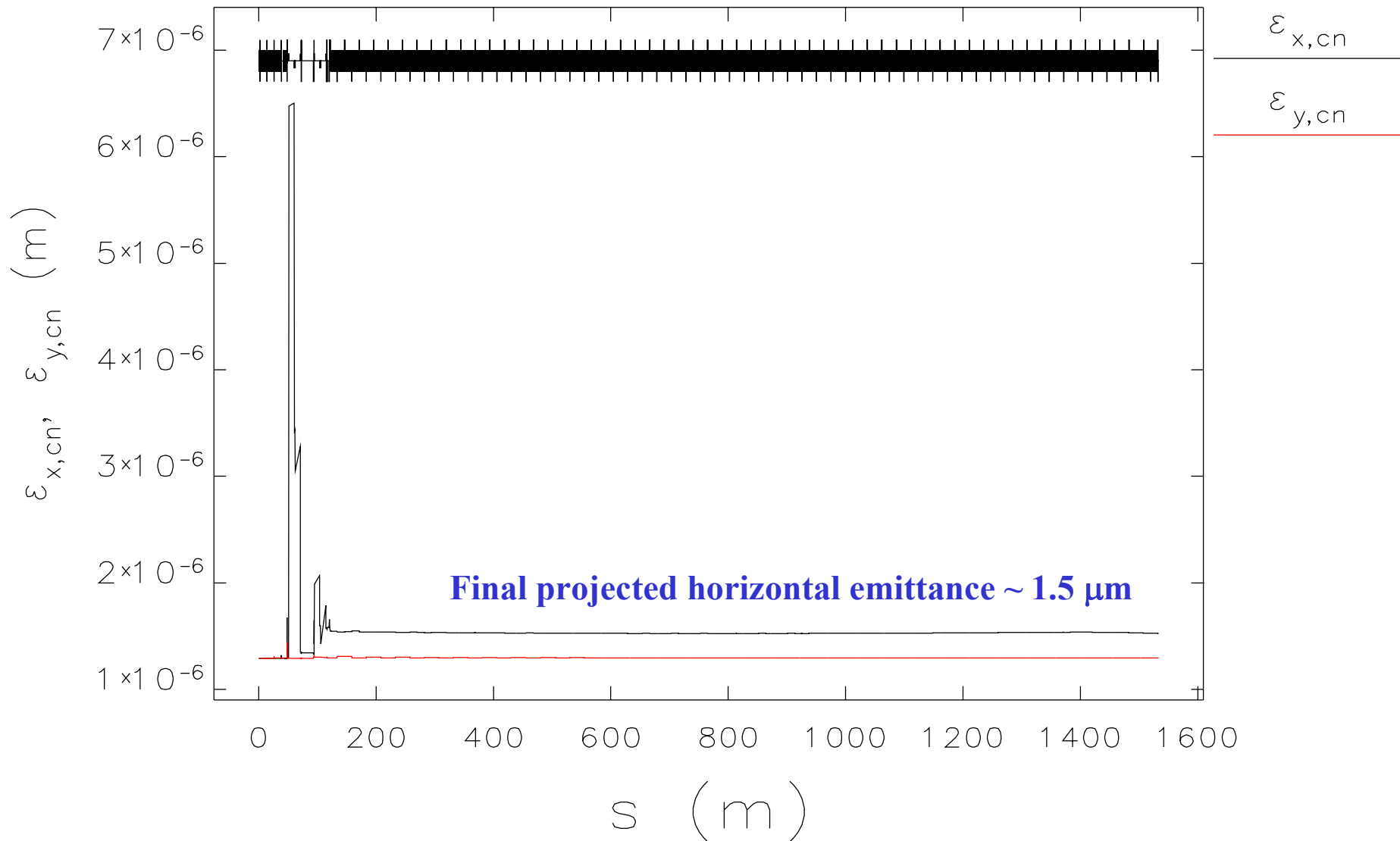
Maximum uncorrelated energy spread : 1.533 MeV @ 20.0 GeV ~ 0.0077% < 0.0125%
Fortunately, the uncorrelated energy spread in the center region is around 0.0013%

Projected Parameters Along Beamline



With TTF2 Injector, $\varepsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

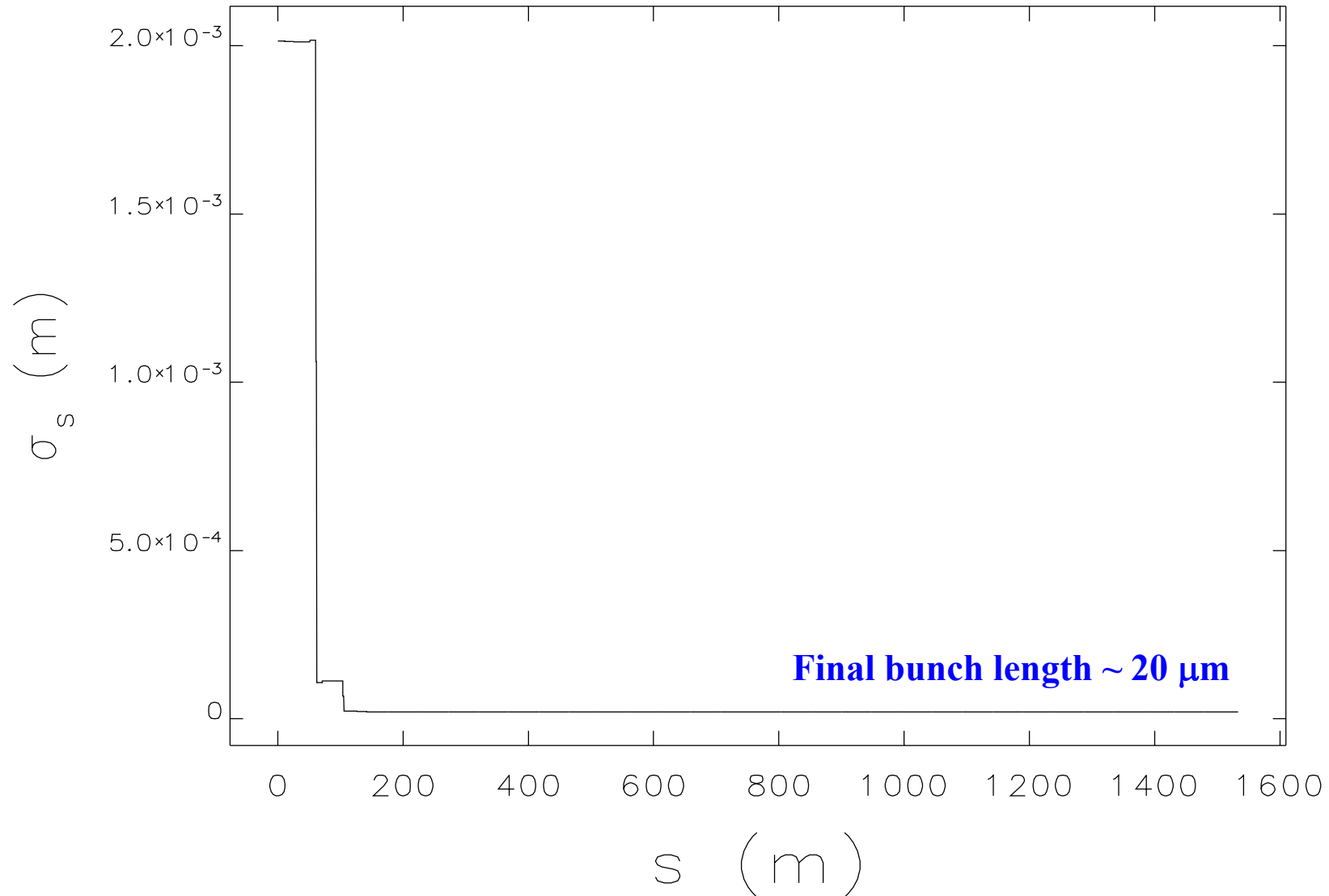


Projected Parameters Along Beamline



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

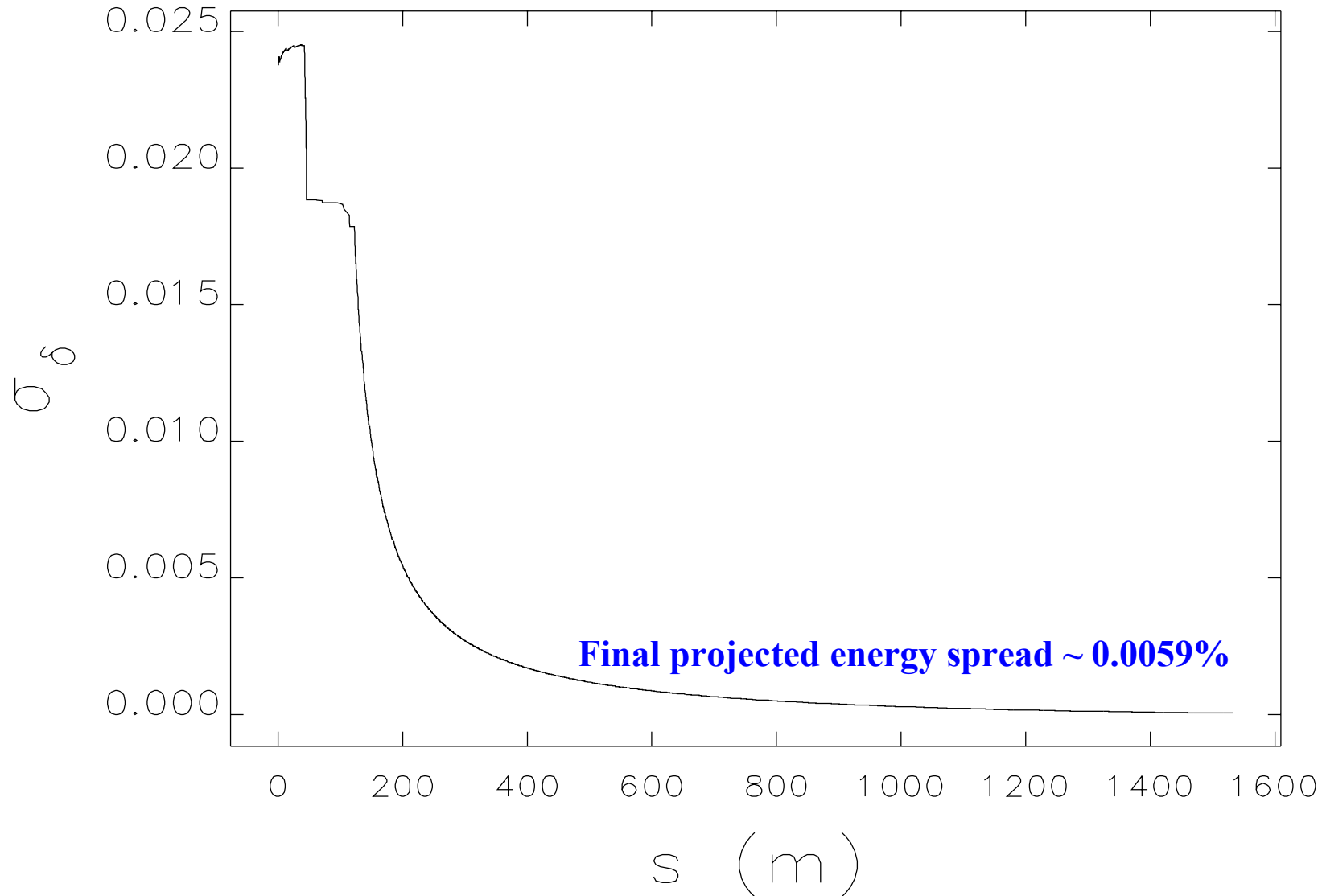


Projected Parameters Along Beamline



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Simulated Particles = 200000

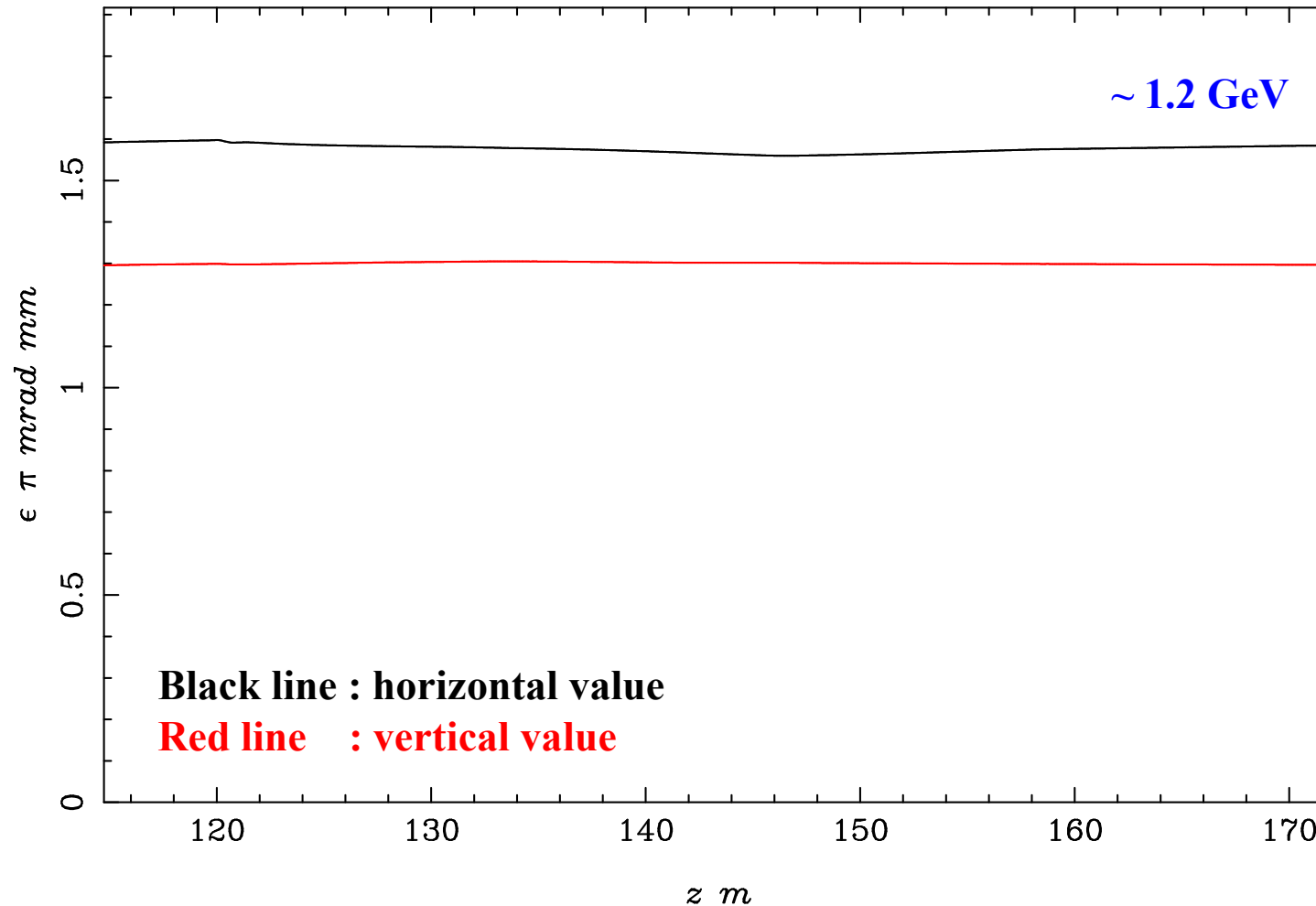


ASTRA Tracking after BC2 with Space Charge



With TTF2 Injector, $\epsilon_n = 1.3 \mu\text{m}$

Transverse Emittance



Projected transverse emittances are almost constant under space charge force !!!

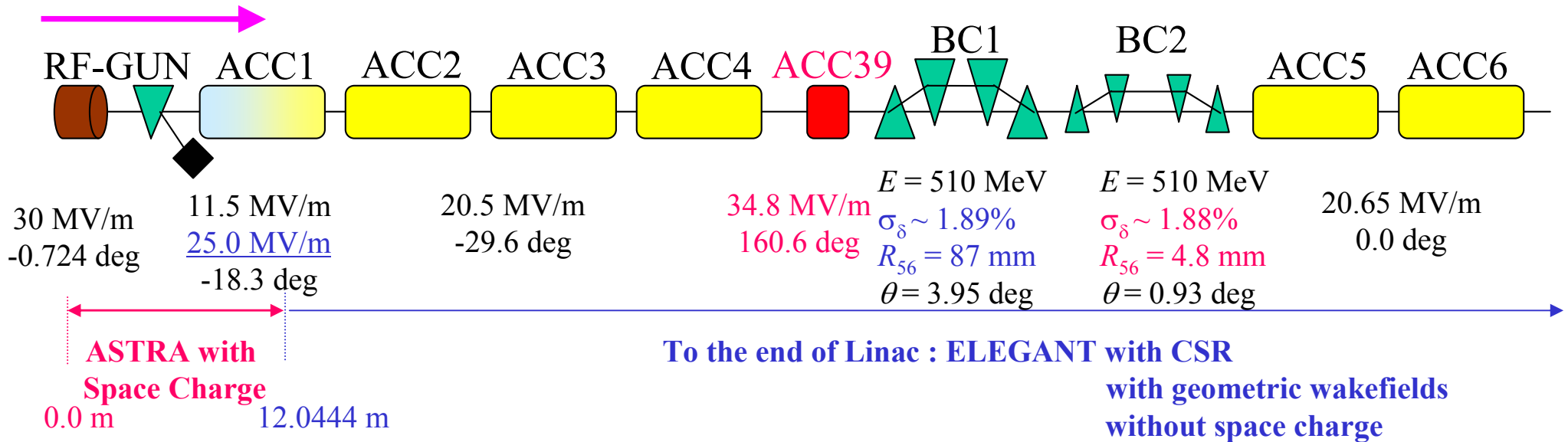
New Lattice for TESLA XFEL – 4th Version



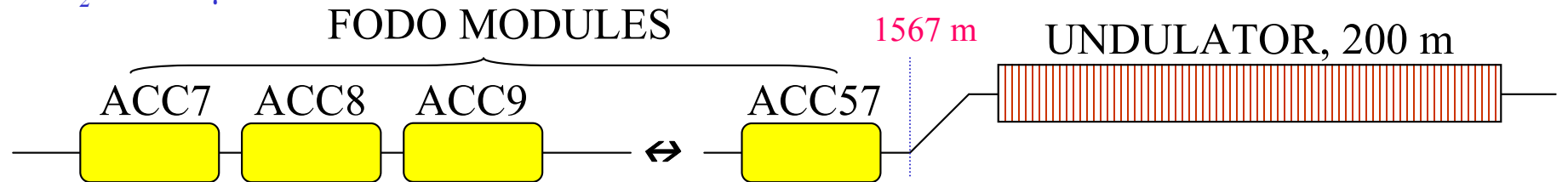
With TESLA XFEL Injector, $\epsilon_n = 0.9 \mu\text{m}$

Q=1.0 nC
e-beam

$\sigma_z = 1.76 \text{ mm} \longrightarrow 113 \mu\text{m} \longrightarrow 23 \mu\text{m}$



$\sigma_z = 20.5 \mu\text{m}$



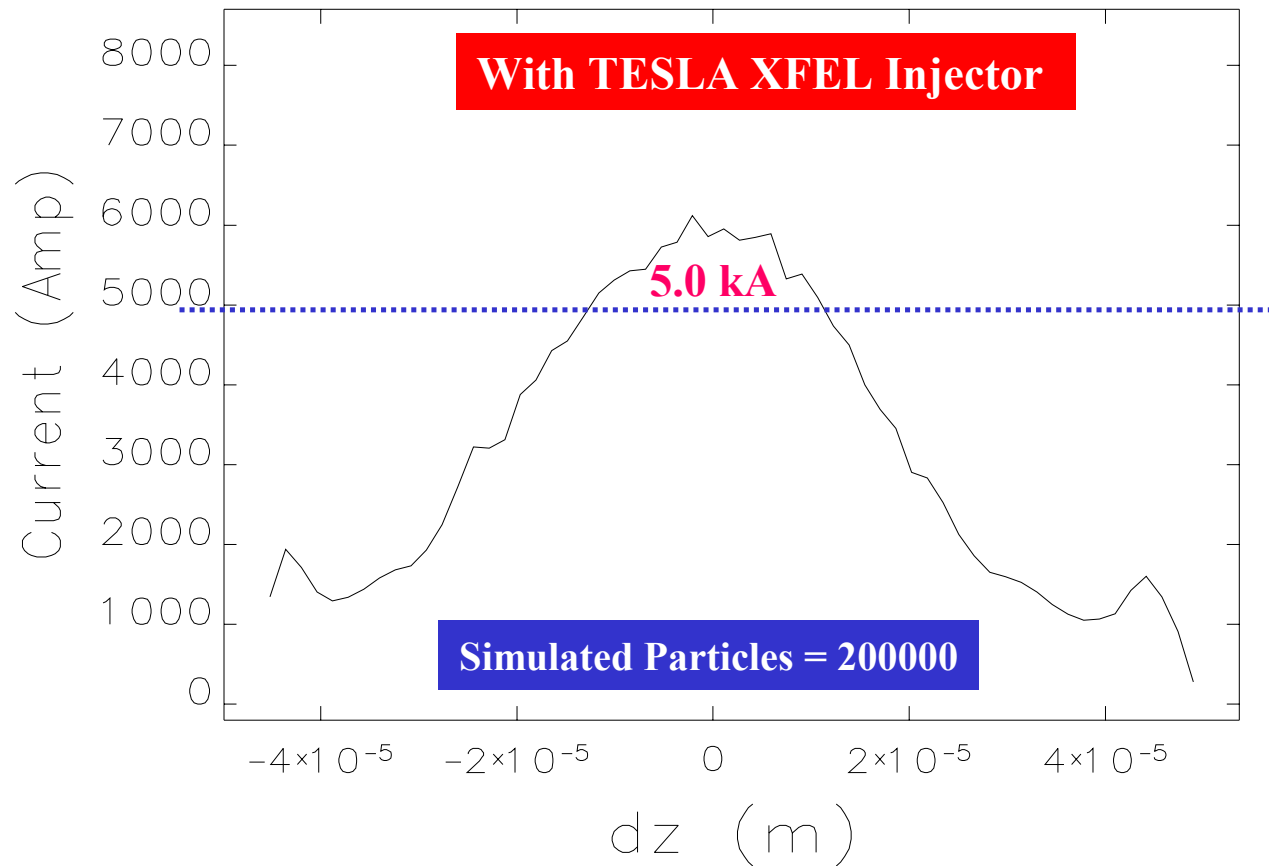
20.65 MV/m
0.0 deg

$E = 20.0 \text{ GeV}$ **All projected parameters !**
 $\sigma_\delta = 0.008\%$
 $\sigma_x = 37.3 \mu\text{m}, \sigma_y = 31.6 \mu\text{m}, \sigma_z = 20.5 \mu\text{m}$
 $\epsilon_{nx} = 1.5 \mu\text{m}, \epsilon_{ny} = 0.94 \mu\text{m}$

4th Lattice with TESLA XFEL Injector Layout



With TESLA XFEL Injector, $\epsilon_n = 0.9 \mu\text{m}$

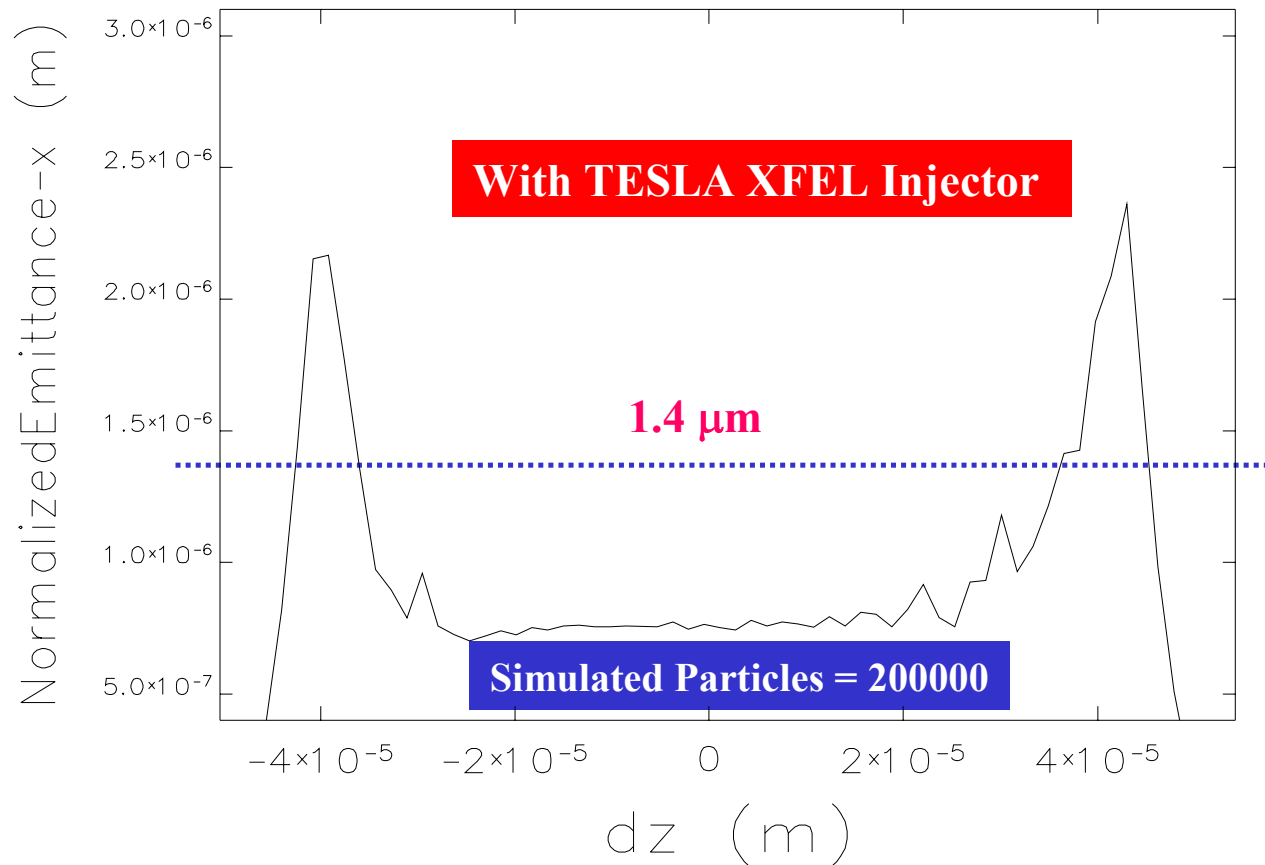


END of LINAC with 60 slices

4th Lattice with TESLA XFEL Injector Layout



With TESLA XFEL Injector, $\epsilon_n = 0.9 \mu\text{m}$

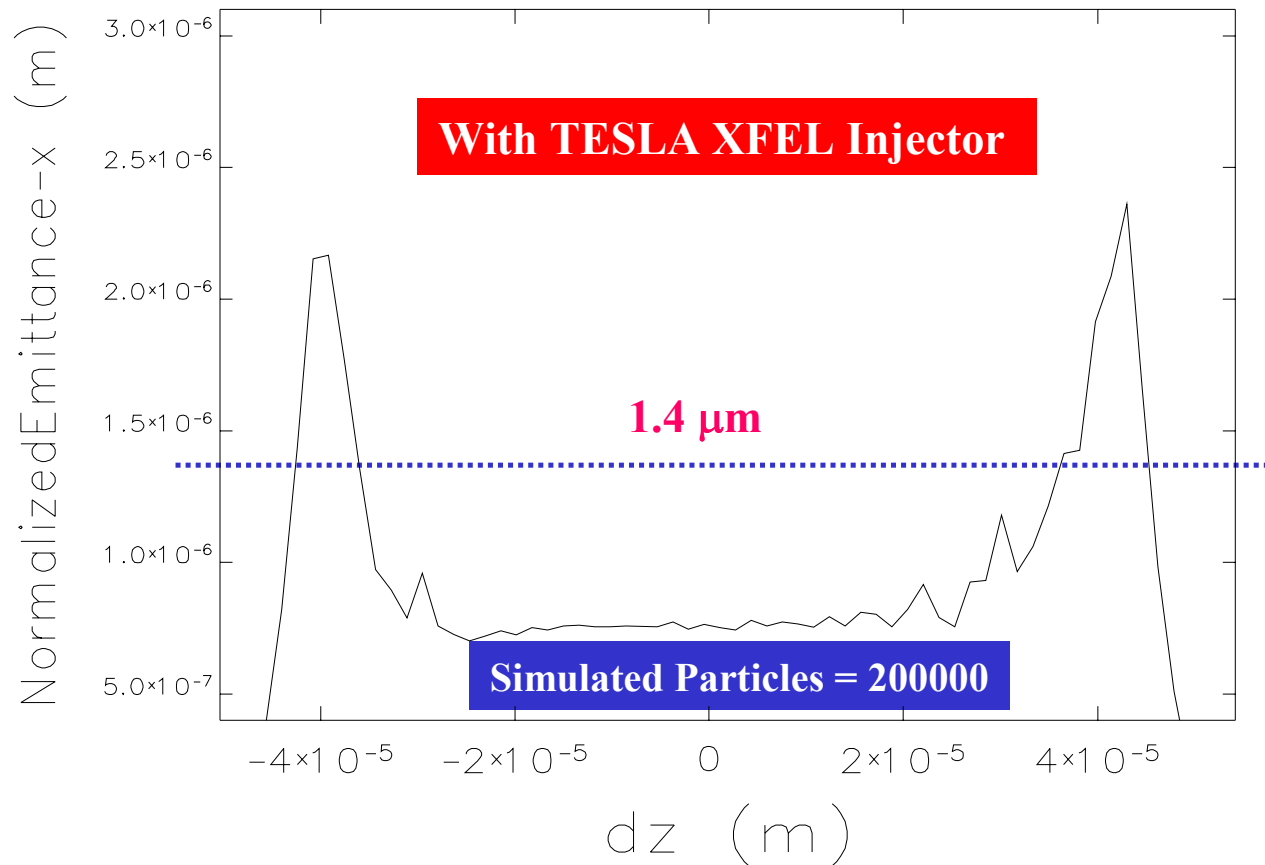


END of LINAC with 60 slices

4th Lattice with TESLA XFEL Injector Layout



With TESLA XFEL Injector, $\epsilon_n = 0.9 \mu\text{m}$



END of LINAC with 60 slices

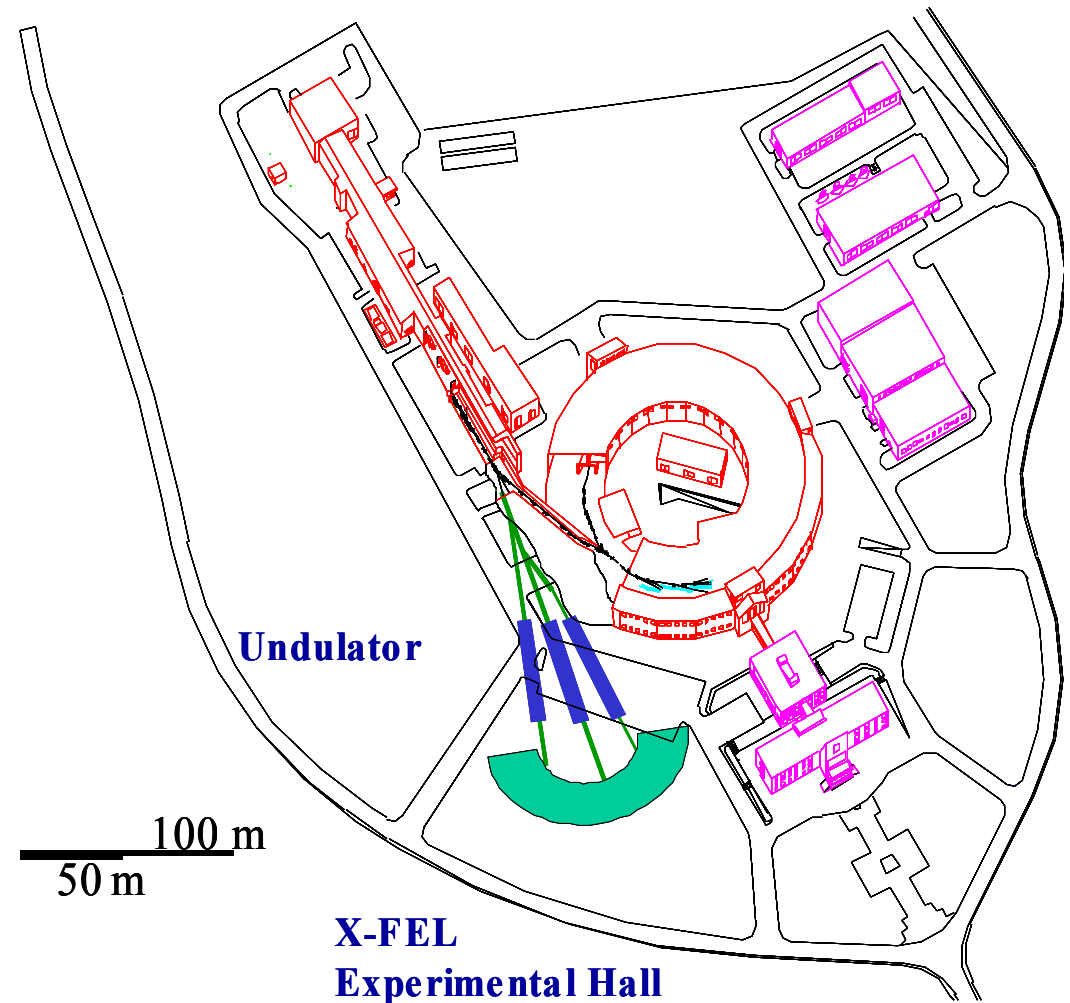
PAL XFEL Project



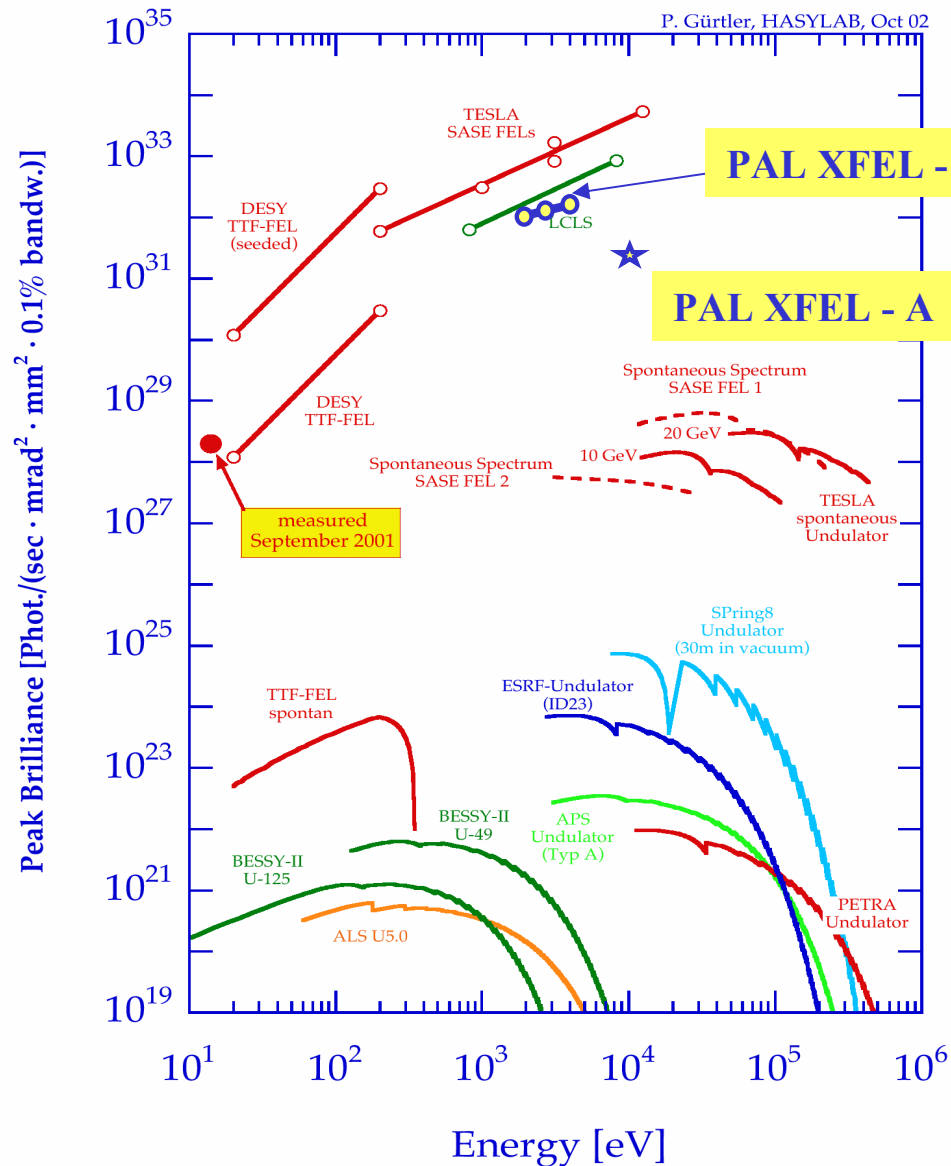
For 3.0 Å PAL XFEL Project:

- Existing 2.5 GeV S-band PLS Linac
- + New S-band Photoinjector
- + New 0.5 GeV Linac
- + New Two Bunch Compressors
- + New ~ 60 m long Undulator

Courtesy of J. S. Oh



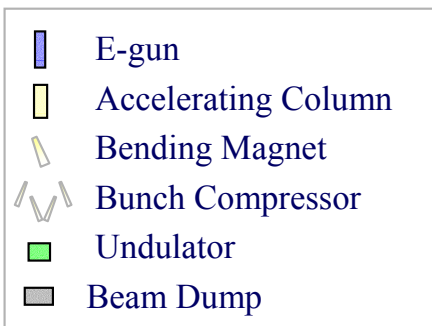
Courtesy of J. S. Oh



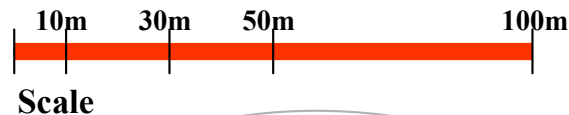
	A	B
λ_x (Å)	1.0	3.0
ϵ_n (μm)	1.0	1.5
τ (fs)	20	233
λ_u (mm)	6.5	12.5
gap (mm)	3.0	3.0
K	0.4	1.14
β (m)	30	15
L_u (m)	22.5 + 36	58.5
PB ($\times 10^{32}$)	0.4	1.4

A: HGHG layout with B undulator

Original Linac Layout for PAL XFEL Project

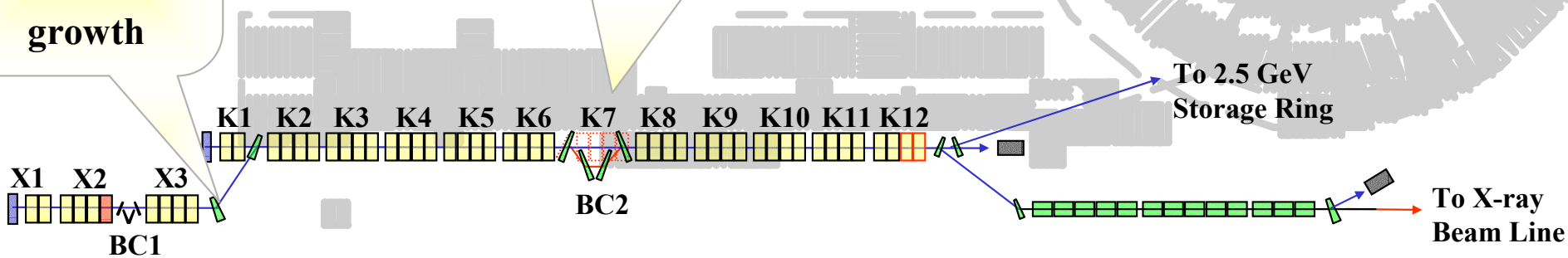
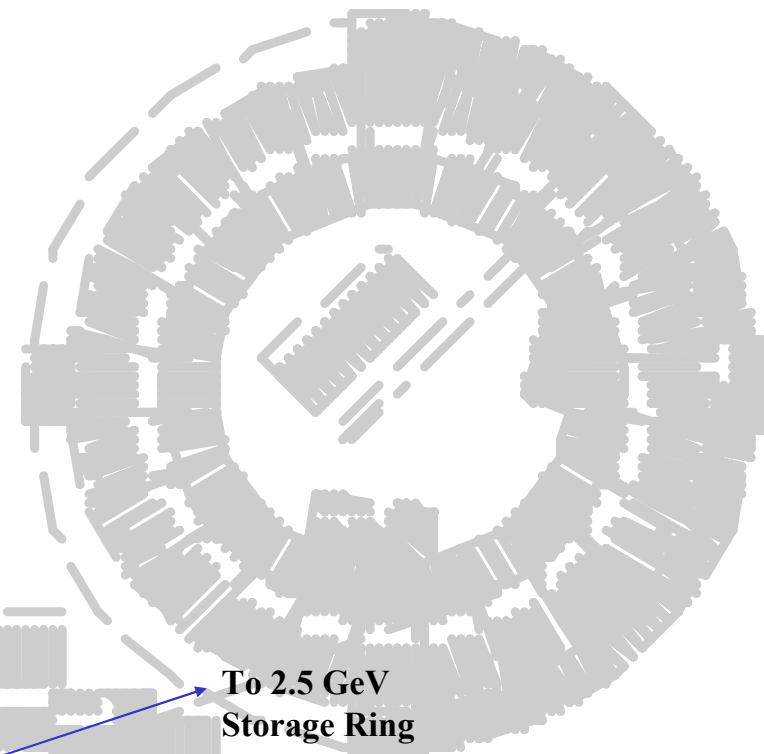


Courtesy of J. S. Oh



Let's remove dogleg to reduce emittance growth

Let's move BC2 to a new linac to control energy spread & microbunching instability



New 0.5 GeV Linac
 X1=150 MeV
 X2=100 MeV
 X3=250 MeV

Existing 2.5 GeV PLS Linac
 K1= 100 MeV Injector
 K2~K6, K8~K12=250 MeV
 BC2 @ K7 + 2 A/C @ K12

New Undulator ($\lambda_x=0.3$ nm)
 $E_0=3.0$ GeV, $\epsilon_n=1.5\mu\text{m-rad}$
 $I_p=4.0$ kA, $\sigma_E=0.02\%$
 $\lambda_U=1.25$ cm, $g=3.0$ mm, $L_U=58.5$ m

New Layout for PAL XFEL - 05DEC03 Version



With S-band Photoinjector, $\epsilon_n = 0.9 \mu\text{m}$

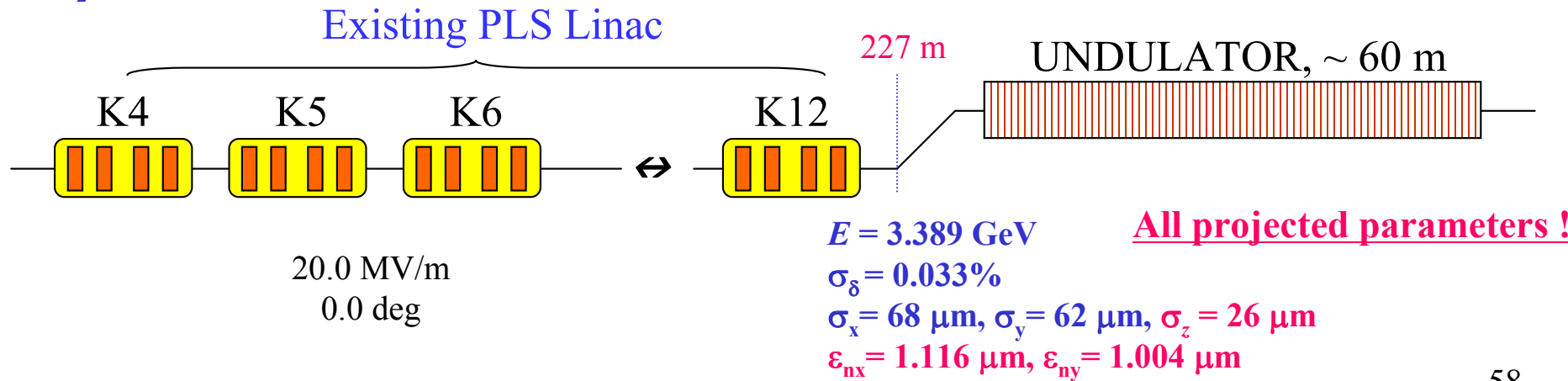
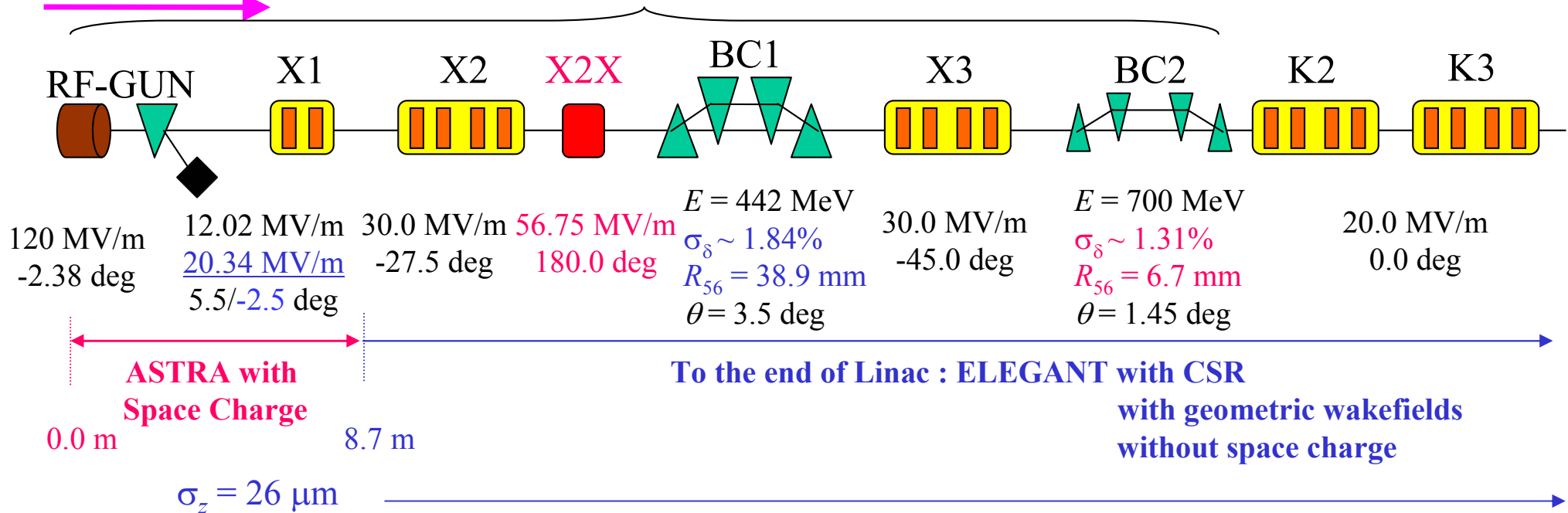
$Q = 1.0 \text{ nC}$

$\sigma_z = 829 \mu\text{m} \longrightarrow 114 \mu\text{m} \longrightarrow 26 \mu\text{m}$

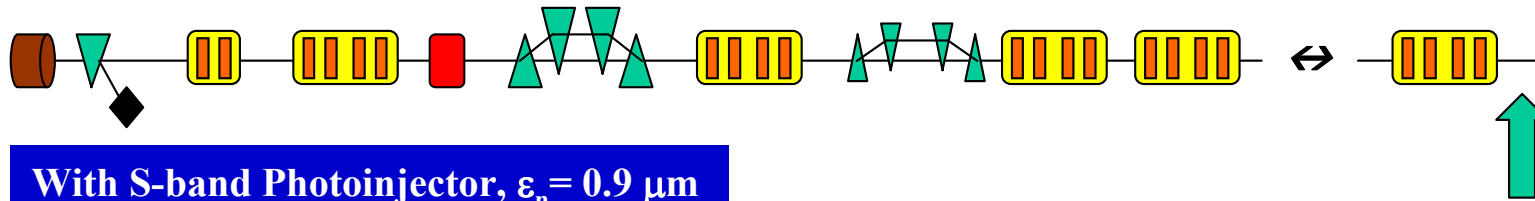
e-beam

New Linac & BCs

Existing PLS Linac

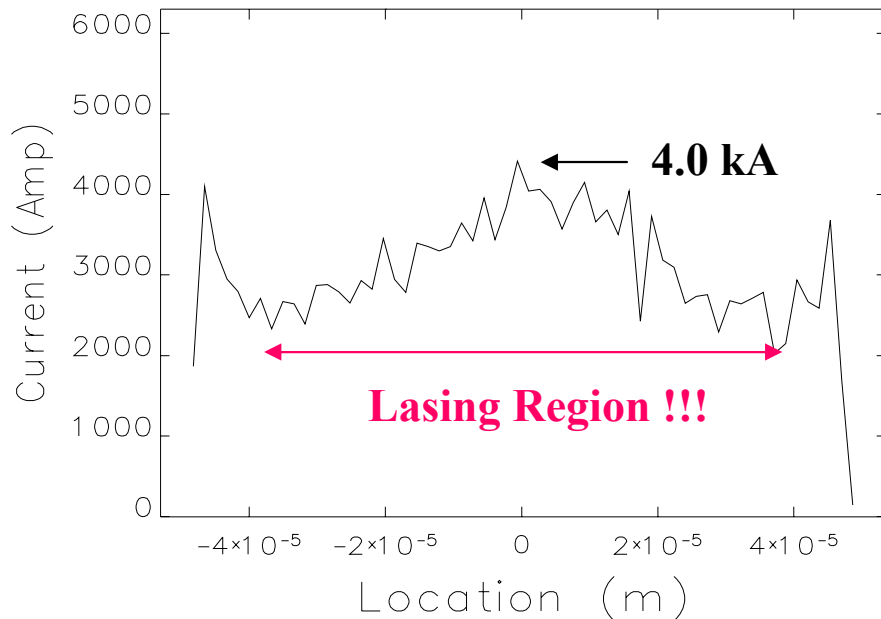


Current & Slice Energy Spread at End of Linac

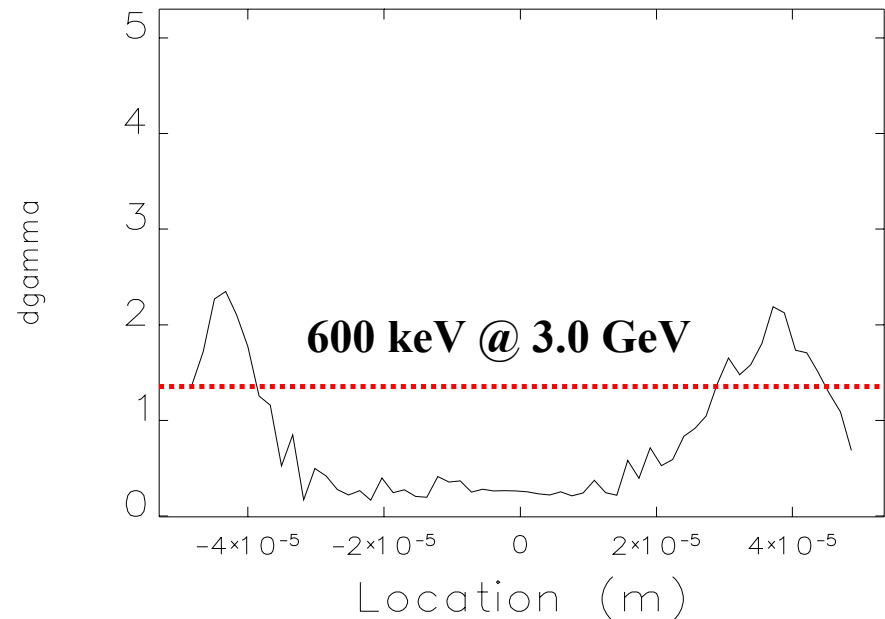


With S-band Photoinjector, $\epsilon_n = 0.9 \mu\text{m}$

Simulated Particles = 200000



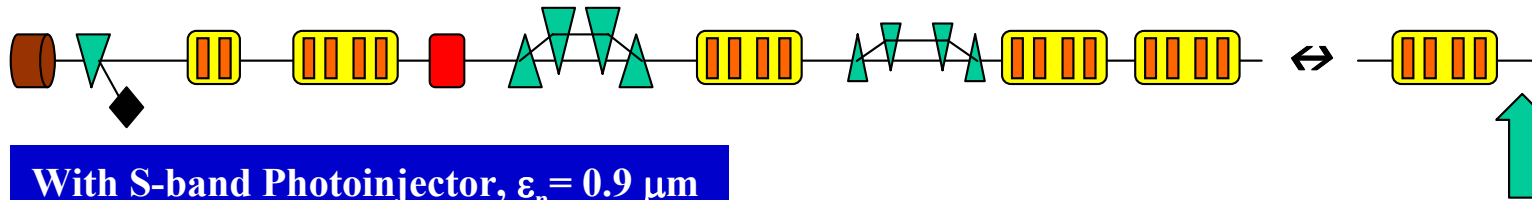
End of Linac with 200000 particles and 60 slices



End of Linac with 200000 particles and 60 slices

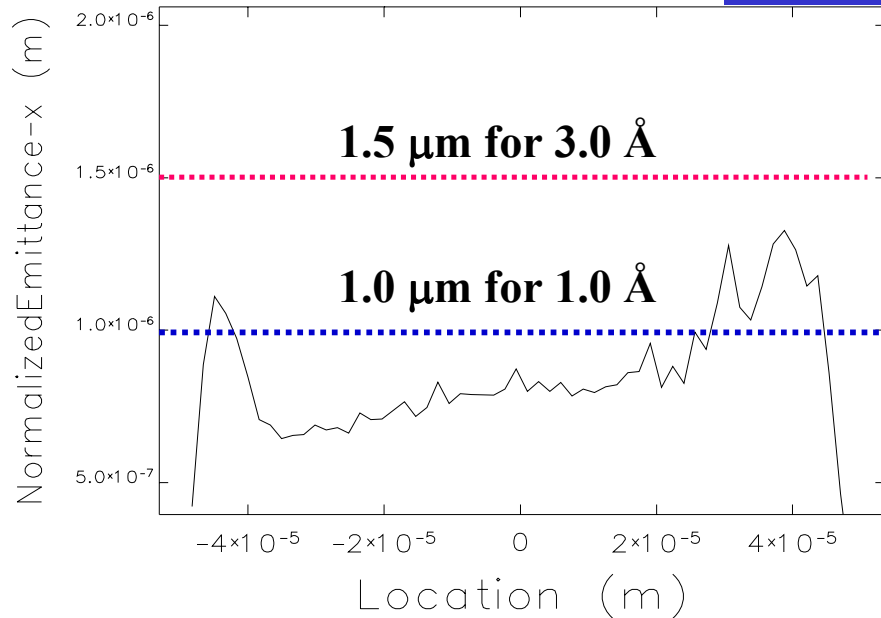
$E = 3.389 \text{ GeV}$
 $\sigma_\delta = 0.033\%$
 $\sigma_x = 68.1 \mu\text{m}, \sigma_y = 61.9 \mu\text{m}, \sigma_z = 26 \mu\text{m}$
 $\epsilon_{nx} = 1.116 \mu\text{m}, \epsilon_{ny} = 1.004 \mu\text{m}$

Slice Emittances at the End of PAL XFEL Linac

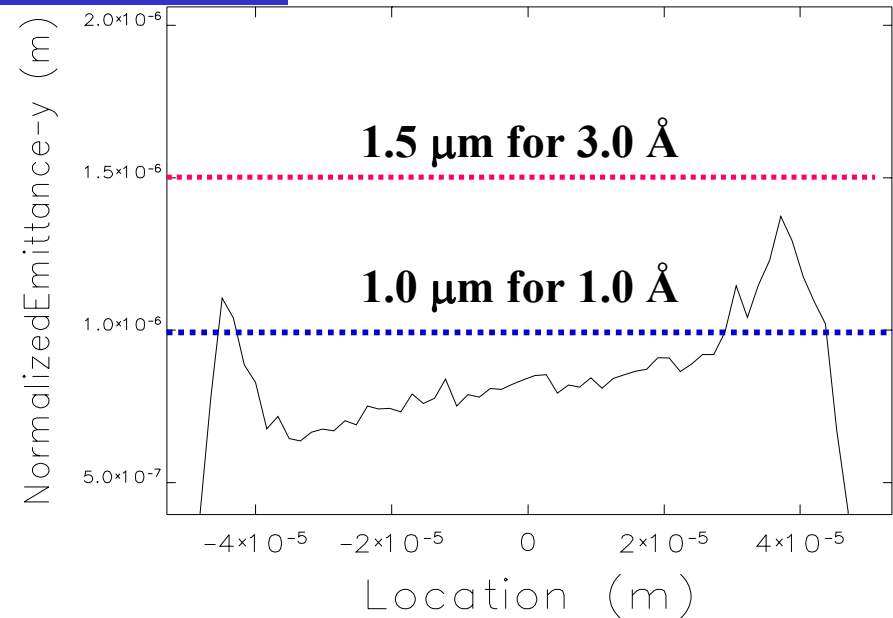


With S-band Photoinjector, $\epsilon_n = 0.9 \mu\text{m}$

Simulated Particles = 200000



End of Linac with 200000 particles and 60 slices



End of Linac with 200000 particles and 60 slices

$E = 3.389 \text{ GeV}$

$\sigma_\delta = 0.033\%$

$\sigma_x = 68.1 \mu\text{m}, \sigma_y = 61.9 \mu\text{m}, \sigma_z = 26 \mu\text{m}$

$\epsilon_{nx} = 1.116 \mu\text{m}, \epsilon_{ny} = 1.004 \mu\text{m}$

We have optimized a new lattice for TESLA XFEL and PAL XFEL projects to reduce the slice parameter growths due to the microbunching instability.

By removing two S-type chicanes and by using only one BC stage with double chicanes, the total number of dipoles in BCs is reduced from 16 to 8, which certainly helped in reducing the microbunching instability at TESLA XFEL.

Although we did not use any laser beam heater in the low energy region or any superconducting wiggler before BC2, we have achieved promising beam parameters for TESLA XFEL and PAL XFEL projects only by optimizing BC layout.

We do not worry about too much on the emittance growth in BCs any more if we optimize BC well with a large energy spread and a long space for BCs.

Acknowledgments



Y. Kim sincerely thanks **Z. Huang, M. Borland, P. Emma, E. A. Schneidmiller, M. Dohlus, J. S. Oh, Professor J. Rossbach, Professor K.-J Kim, Professor I. S. Ko, and Professor W. Namkung** for their encouragements of this work and many useful comments and discussions on the BC layout and on the microbunching instability.