

TESLA Collaboration Meeting

ISTITUTO NAZIONALE FISICA NUCLEARE
Laboratori Nazionali di Legnaro
Legnaro (Padua) ITALY

Enzo Palmieri

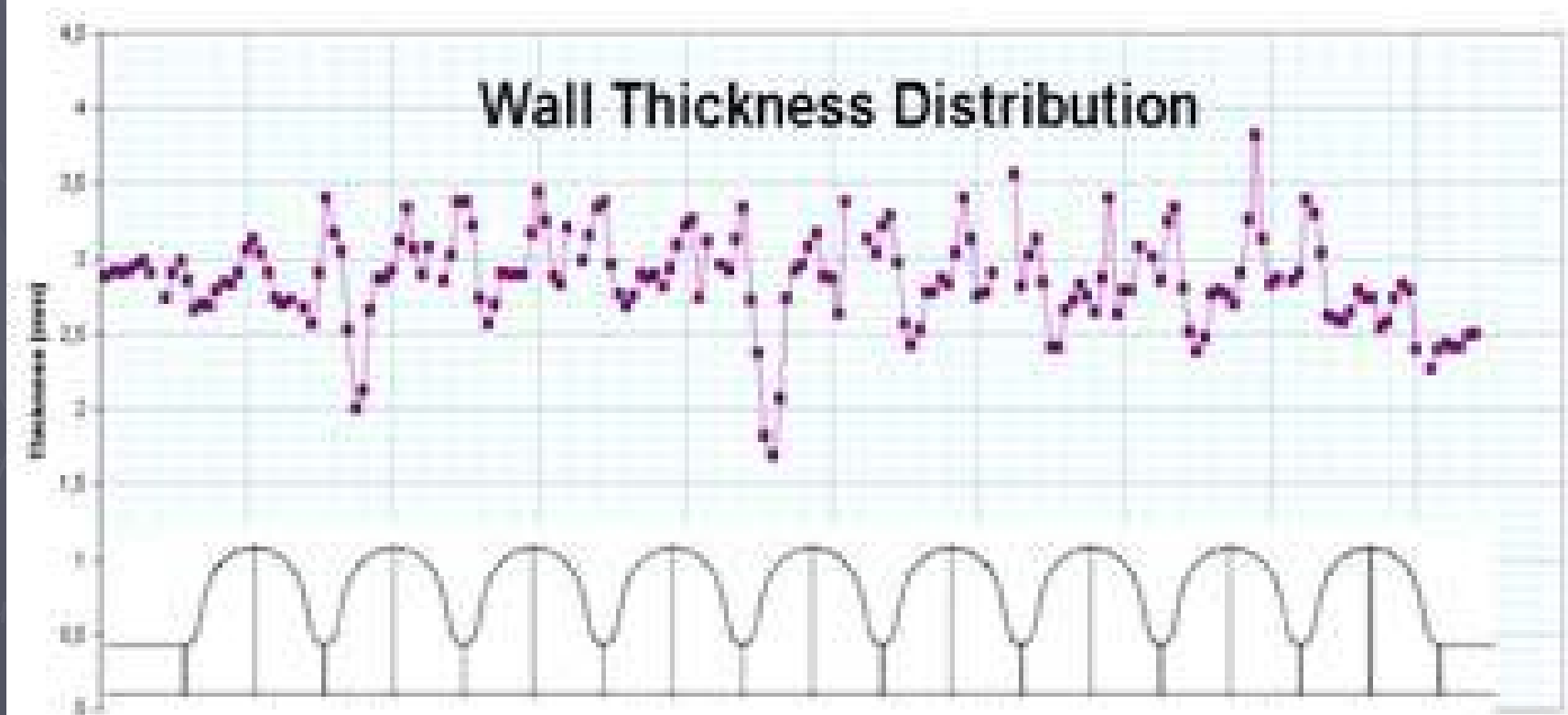
Desy-Zeuten January 21-21, 2004

Contents:

- ▶ Seamless cavity spinning from circular blanks
- ▶ Nb/Cu thin film investigation: Problem of Q-decay vs E_{acc}

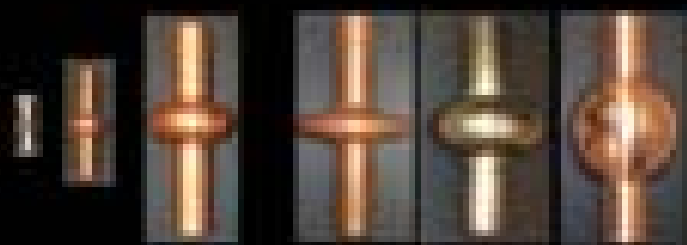


The first Niobium seamless 9-cell cavity ever fabricated



Advantages of Spinning technology

- Short fabrication times
- No welds
- No intermediate annealing
- Almost no swarf
- Low fabrication costs
- Machine adaptability to almost any cavity shape and size without further investments in equipment

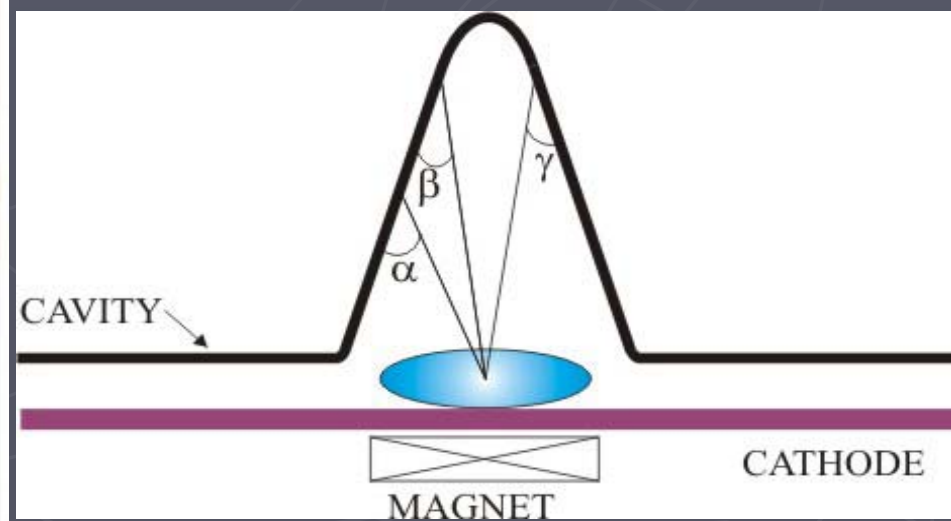
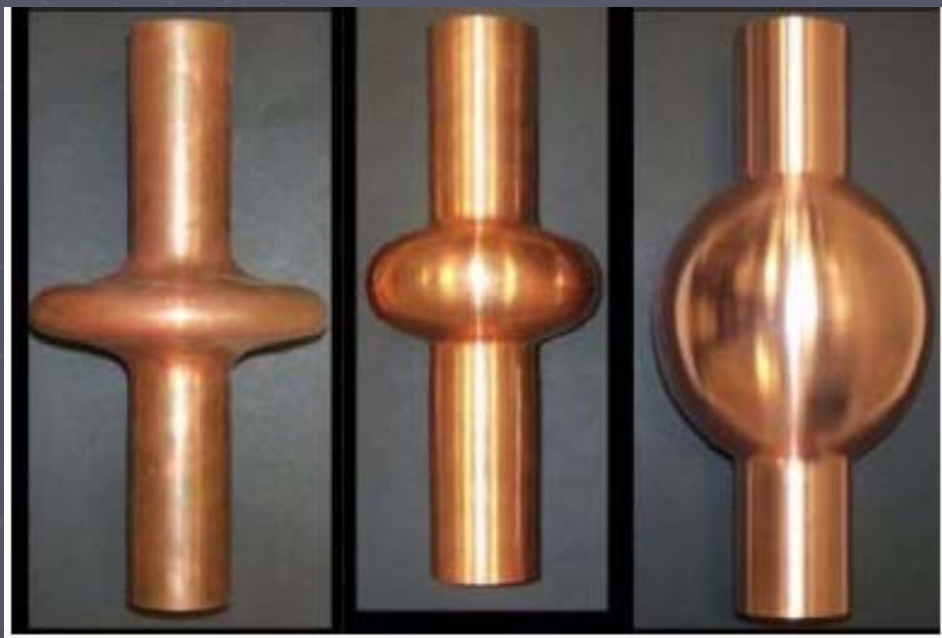
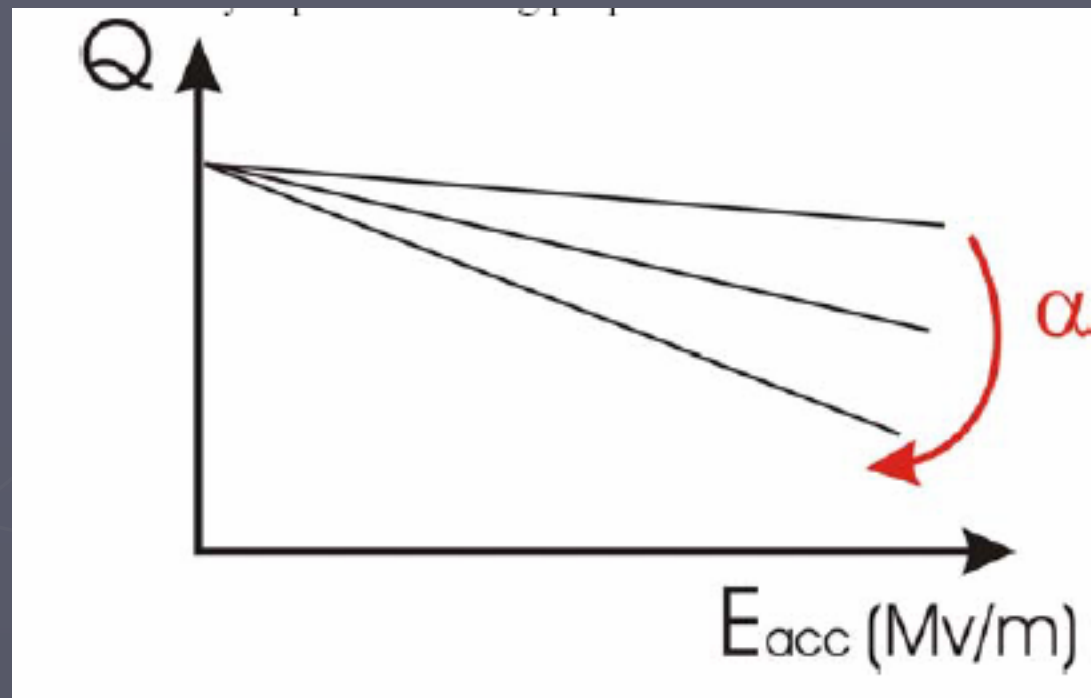


6 GHz



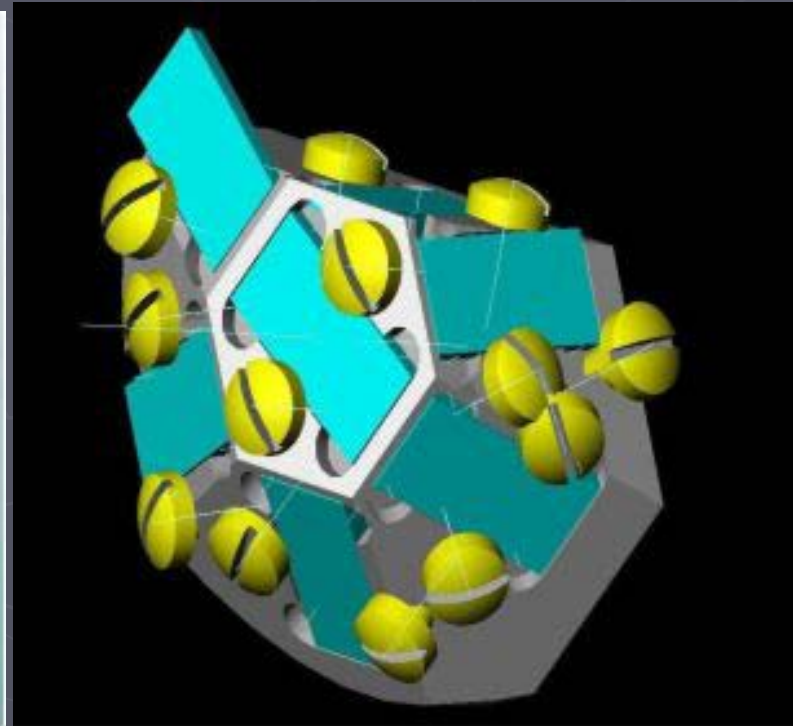
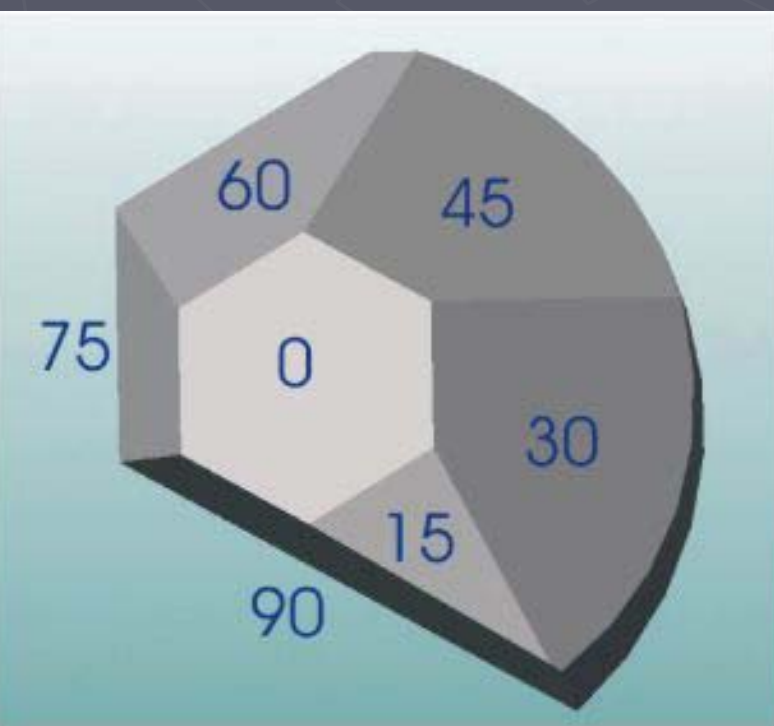
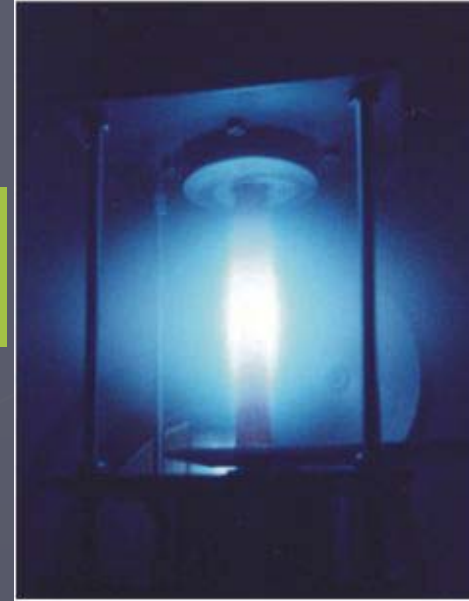
500 MHz

Thin film Research

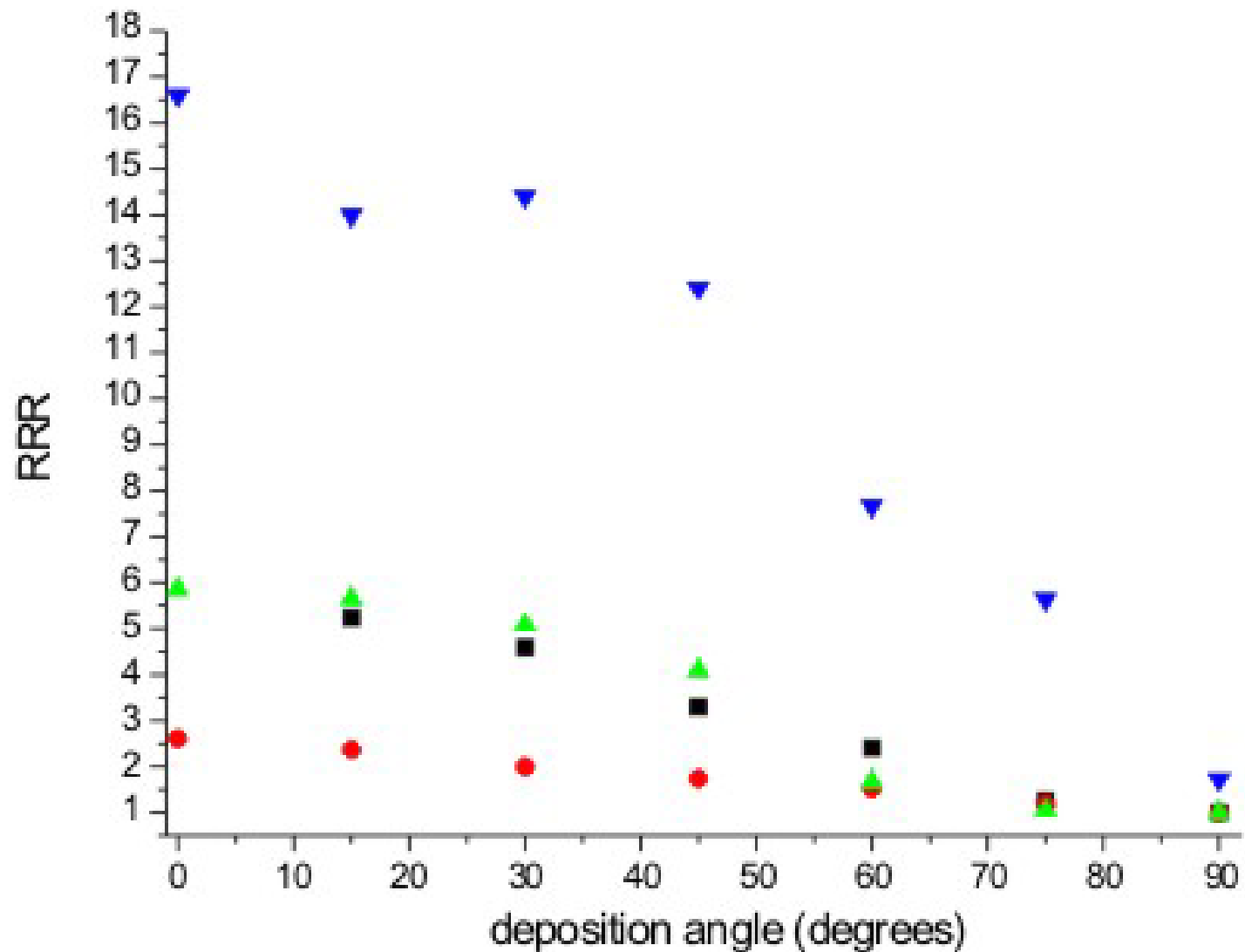


Multi-angle substrate holders

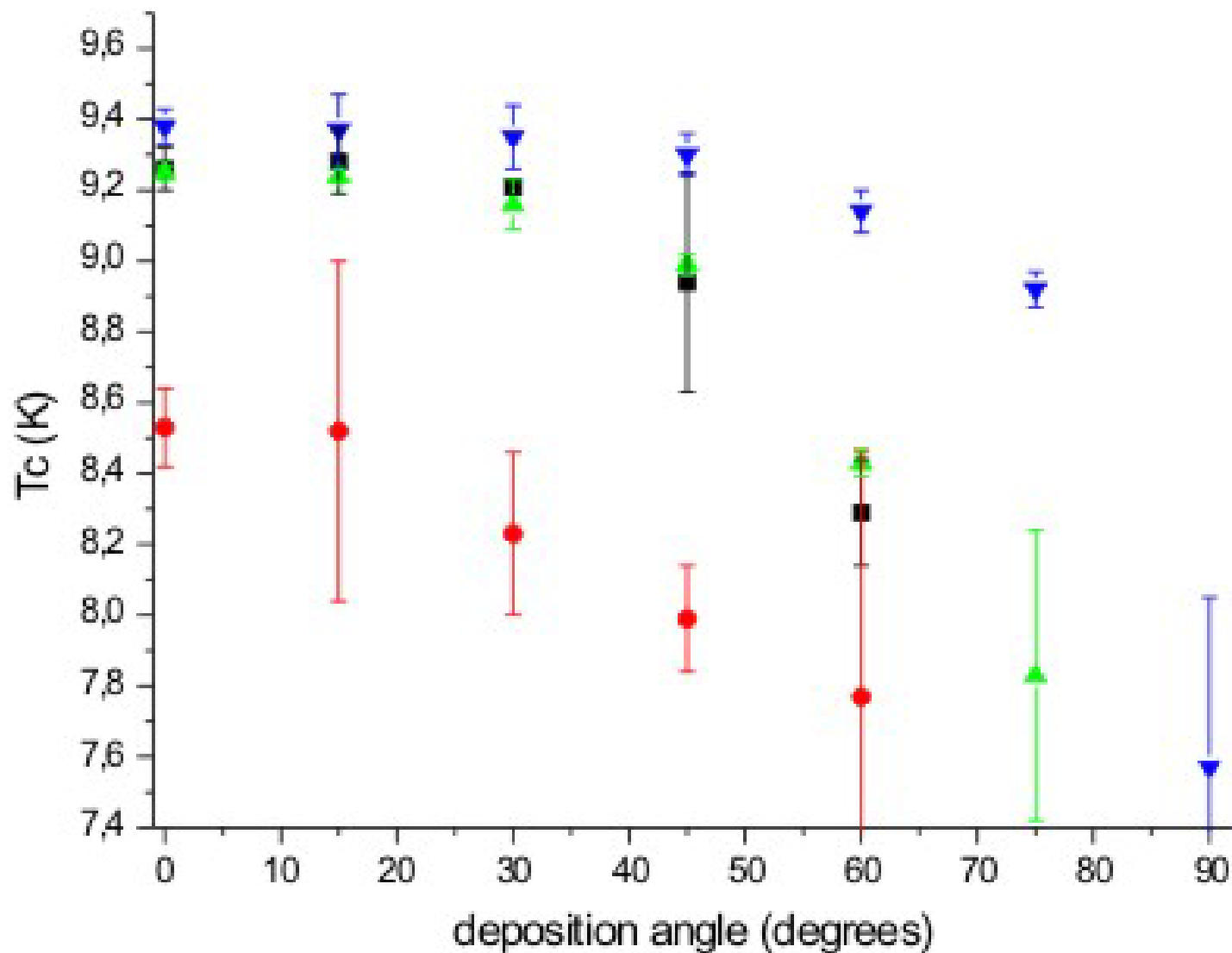
faces at 0° , 15° , 30° , 45° , 60° , 75° , 90°



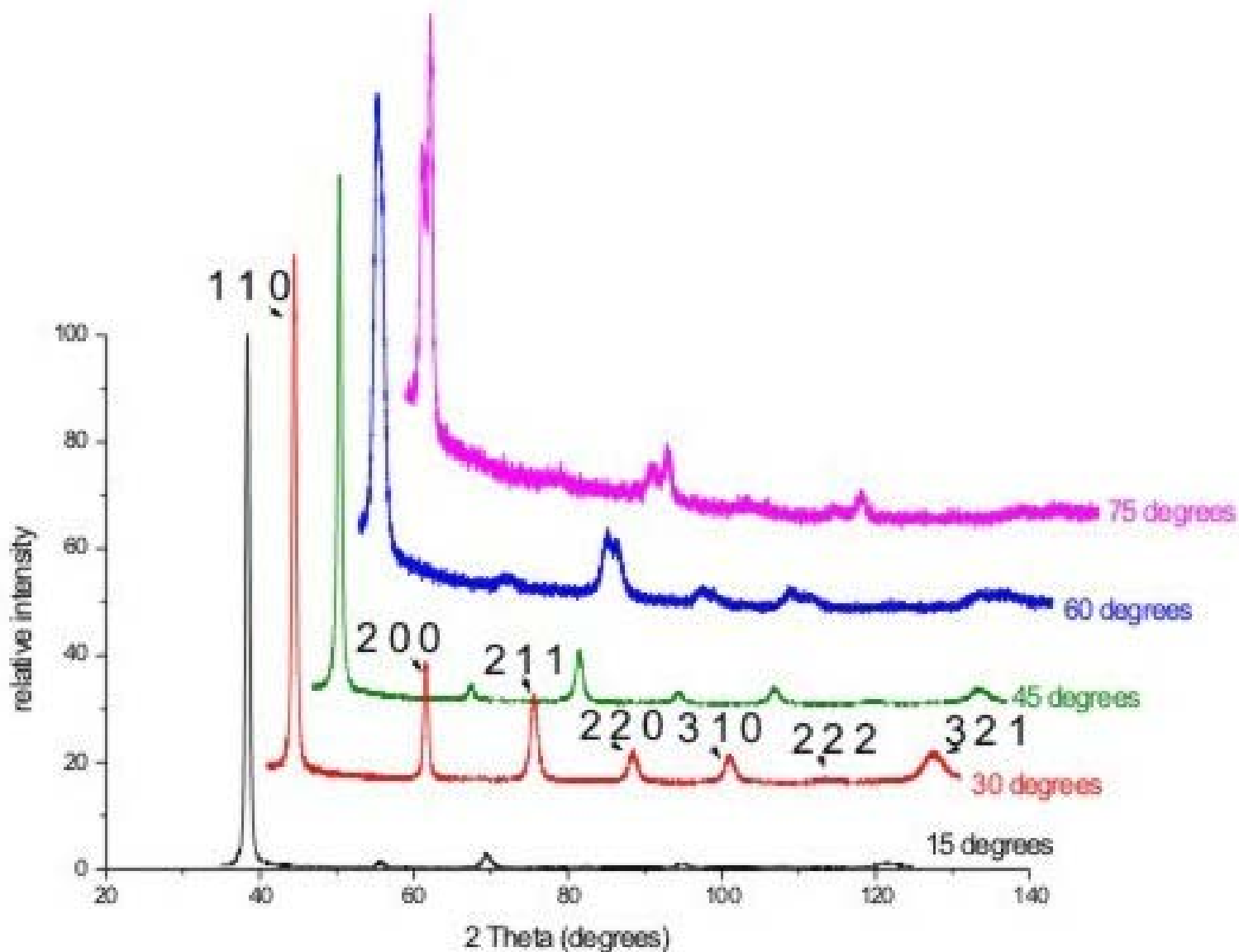
RRR vs. deposition angle for samples at different level of contamination



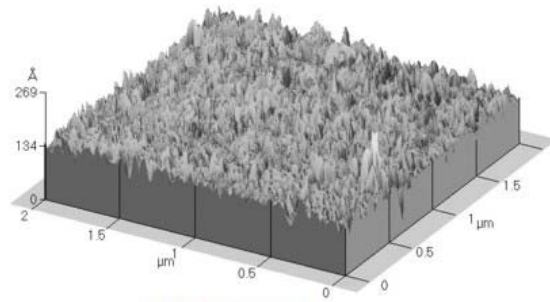
T_c vs. deposition angle for samples at different level of contamination



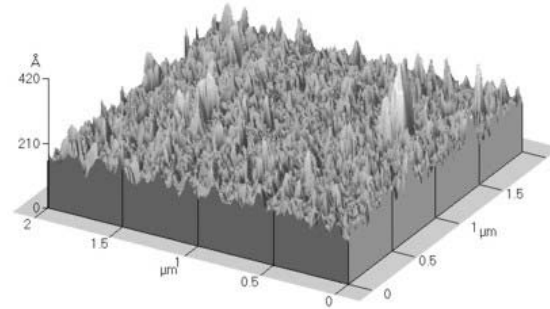
Normalized X-ray diffraction spectra at various deposition angles



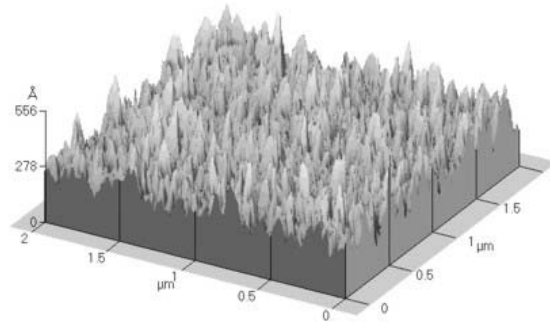
AFM topographic images of niobium films at different target-substrate angle



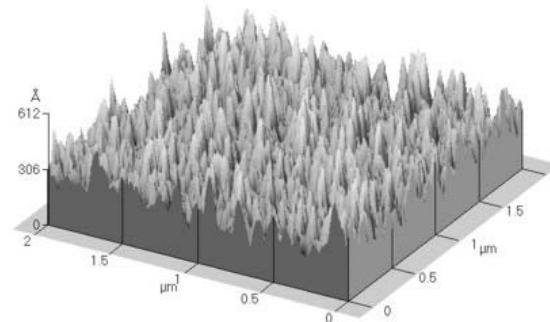
15 degrees



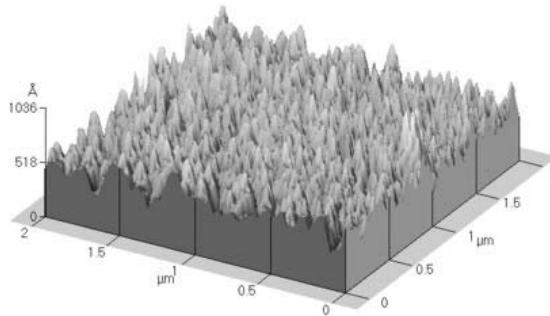
30 degrees



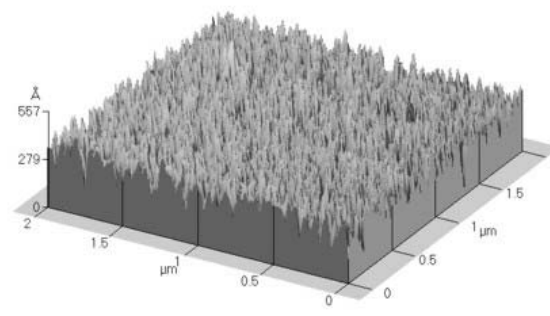
45 degrees



60 degrees

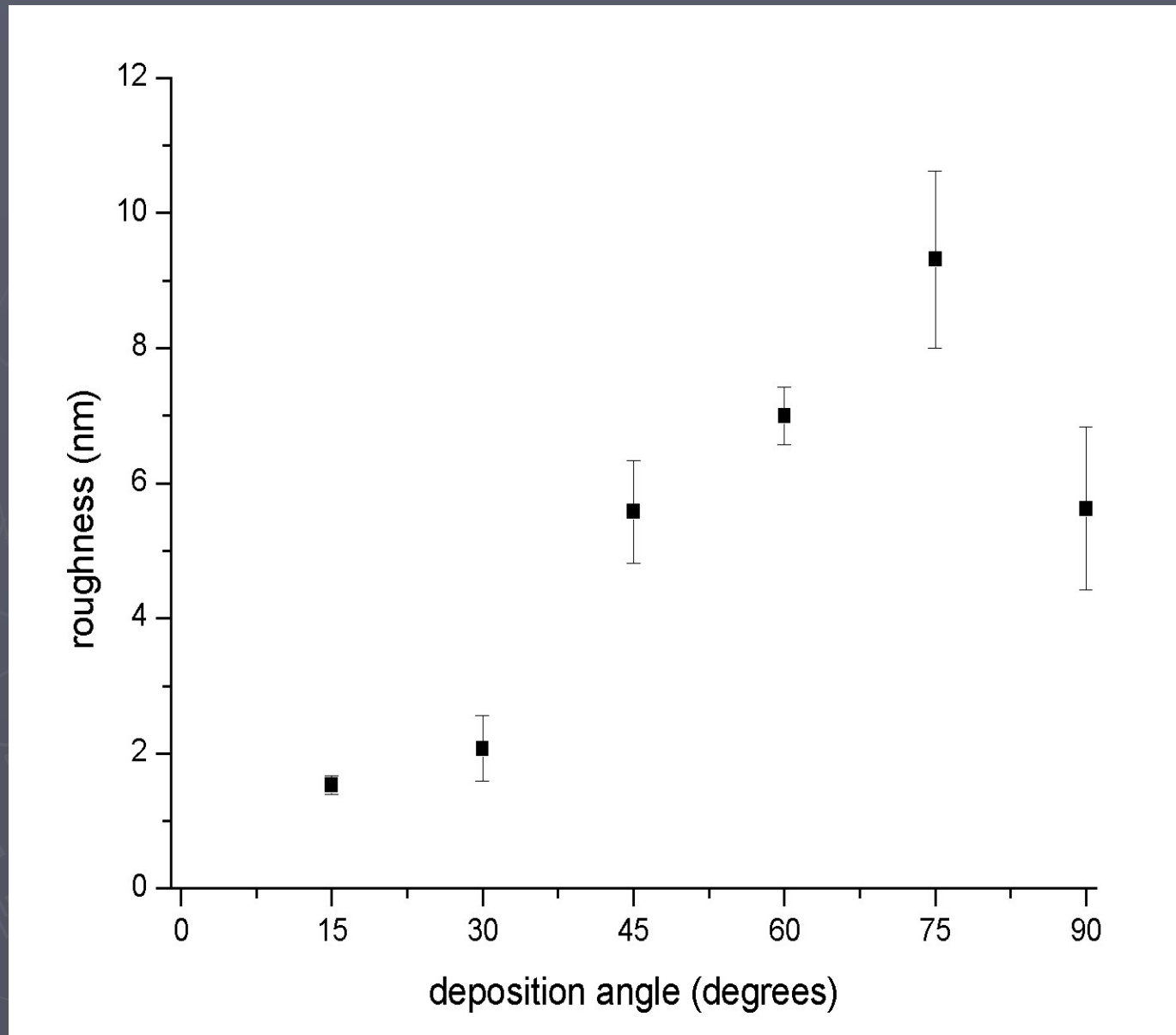


75 degrees

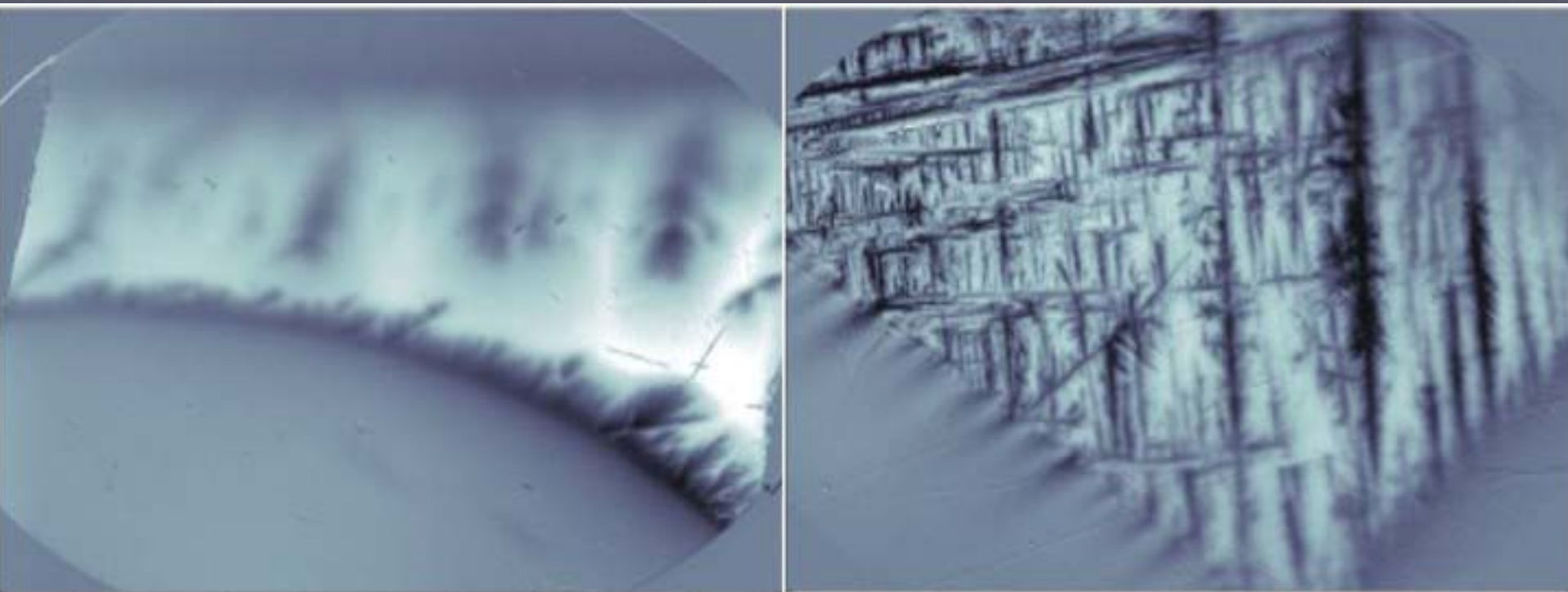


90 degrees

AFM average roughness vs. deposition angle



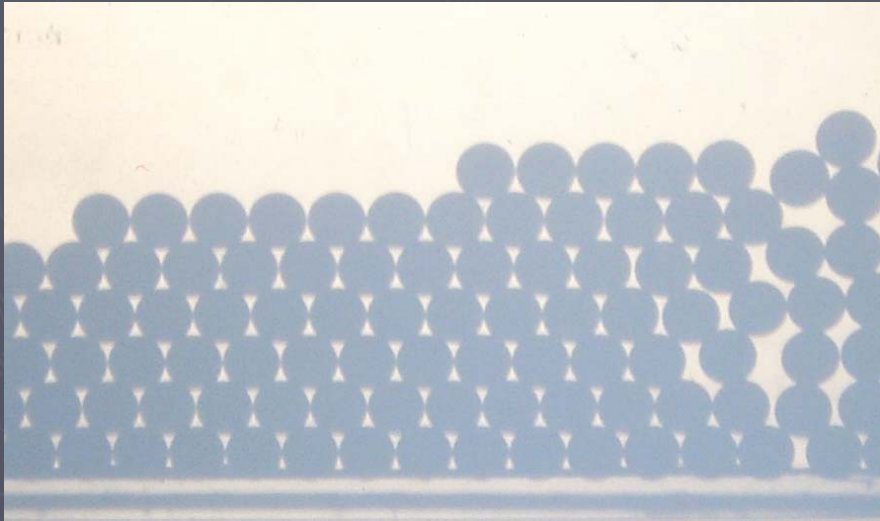
Magneto-optical images of Nb film deposited on Cu for parallel target - substrate (left) and at 45 (right) degrees



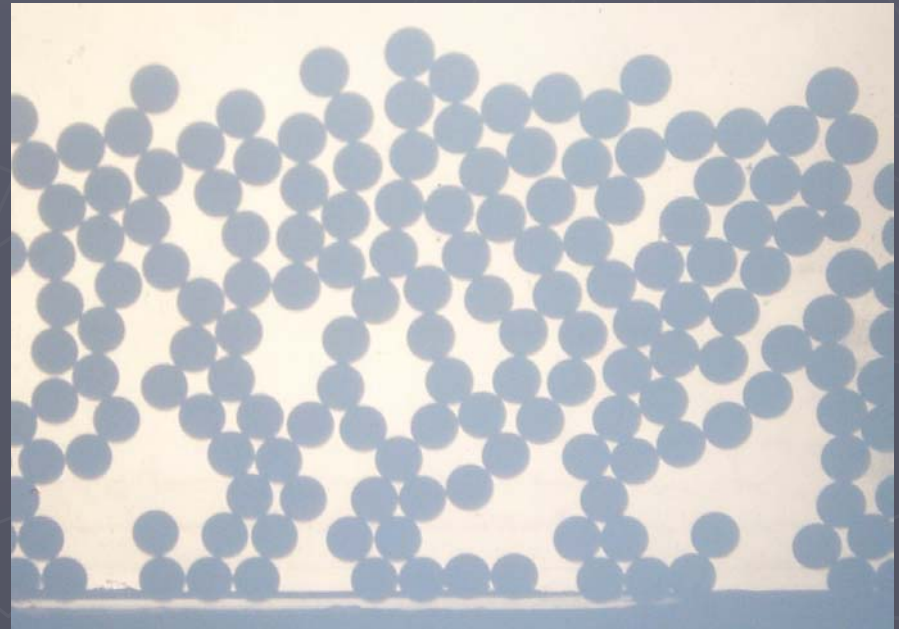
External applied field is 176 mT at $T = 5\text{K}$

Simulation of film growth for:

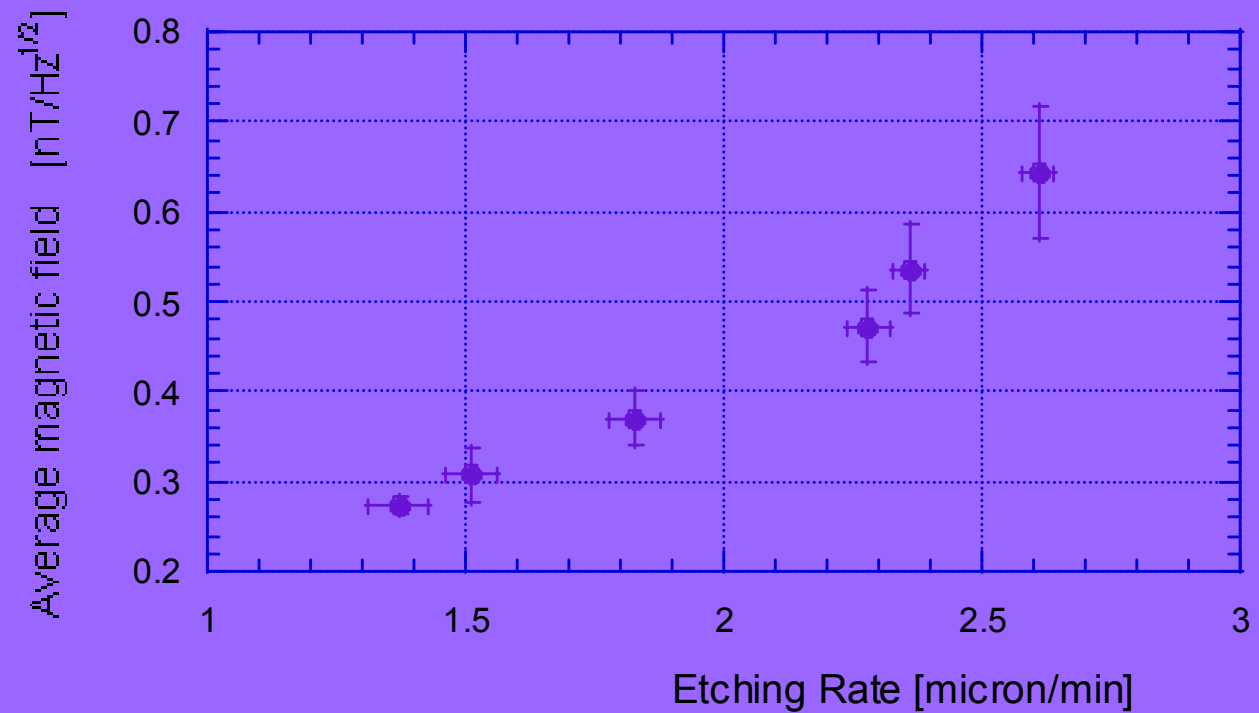
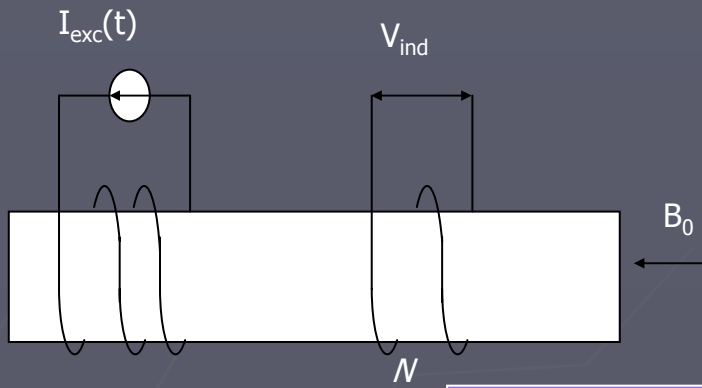
target parallel to substrate

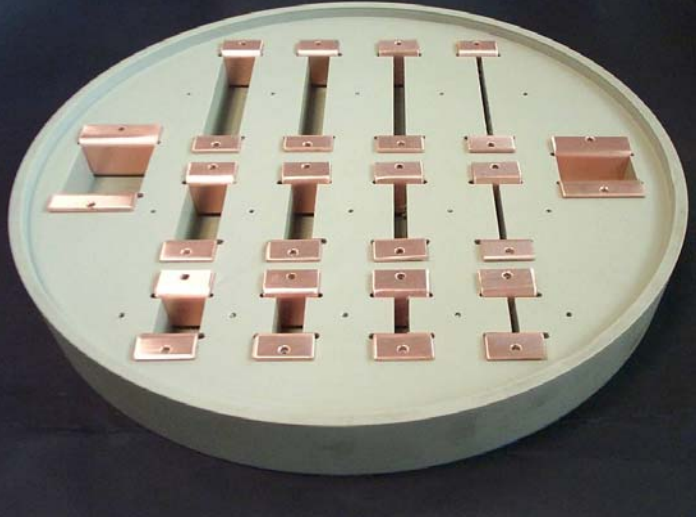


45 degrees orientation

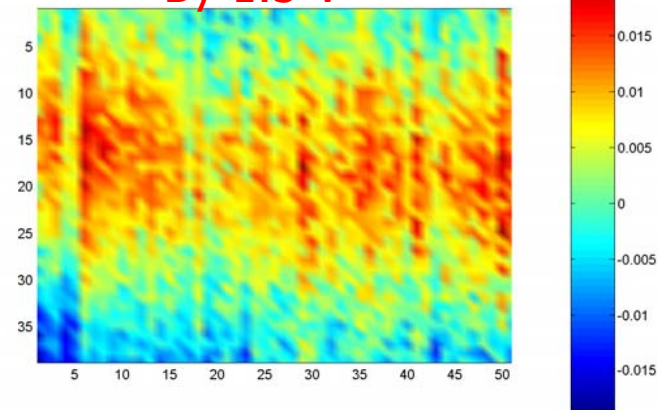


Flux gate Magnetometer

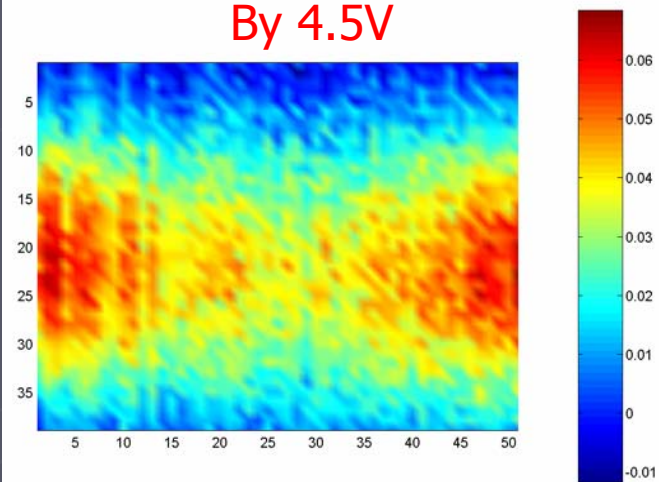




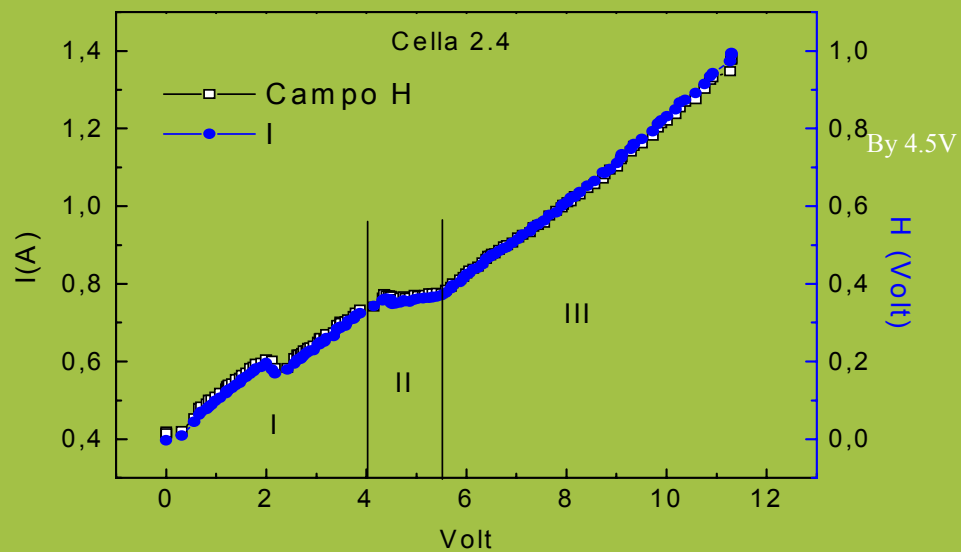
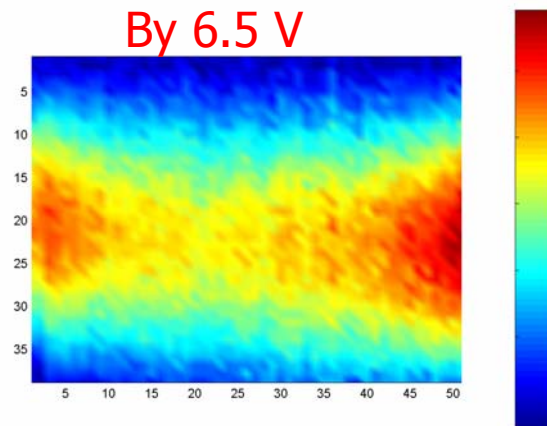
By 1.5 V



By 4.5V



By 6.5 V



Inverting the Biot-Savart law

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}') \wedge (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} d^3 r'$$

assuming the continuity equation and the condition $\mathbf{J}_z = 0$,
it is possible to deduce the current density induced by the magnetic field

$$\begin{aligned} \mu_0 \mathbf{H}_{ind} \cdot \hat{\mathbf{i}} &= \mu_0 H_x(x, y, z) = \\ &= \frac{\mu_0}{4\pi} \iiint \frac{J_y(x', y', z') \cdot (z - z') - J_z(x', y', z') \cdot (y - y')}{\sqrt{[(x - x')^2 + (y - y')^2 + (z - z')^2]^3}} dx' dy' dz' \end{aligned}$$

$$\begin{aligned} \mu_0 \mathbf{H}_{ind} \cdot \hat{\mathbf{j}} &= \mu_0 H_y(x, y, z) = \\ &= \frac{\mu_0}{4\pi} \iiint \frac{J_z(x', y', z') \cdot (x - x') - J_x(x', y', z') \cdot (z - z')}{\sqrt{[(x - x')^2 + (y - y')^2 + (z - z')^2]^3}} dx' dy' dz' \end{aligned}$$

$$\begin{aligned} \mu_0 \mathbf{H}_{ind} \cdot \hat{\mathbf{k}} &= \mu_0 H_z(x, y, z) = \\ &= \frac{\mu_0}{4\pi} \iiint \frac{J_x(x', y', z') \cdot (y - y') - J_y(x', y', z') \cdot (x - x')}{\sqrt{[(x - x')^2 + (y - y')^2 + (z - z')^2]^3}} dx' dy' dz' \end{aligned}$$

