

Introduction to Accelerator Physics

Part 4

Pedro Castro / Accelerator Physics Group (MPY)
Introduction to Accelerator Physics
DESY, 29th July 2014

Differences between proton and electron accelerators

HERA (Hadron Electron Ring Accelerator) tunnel:

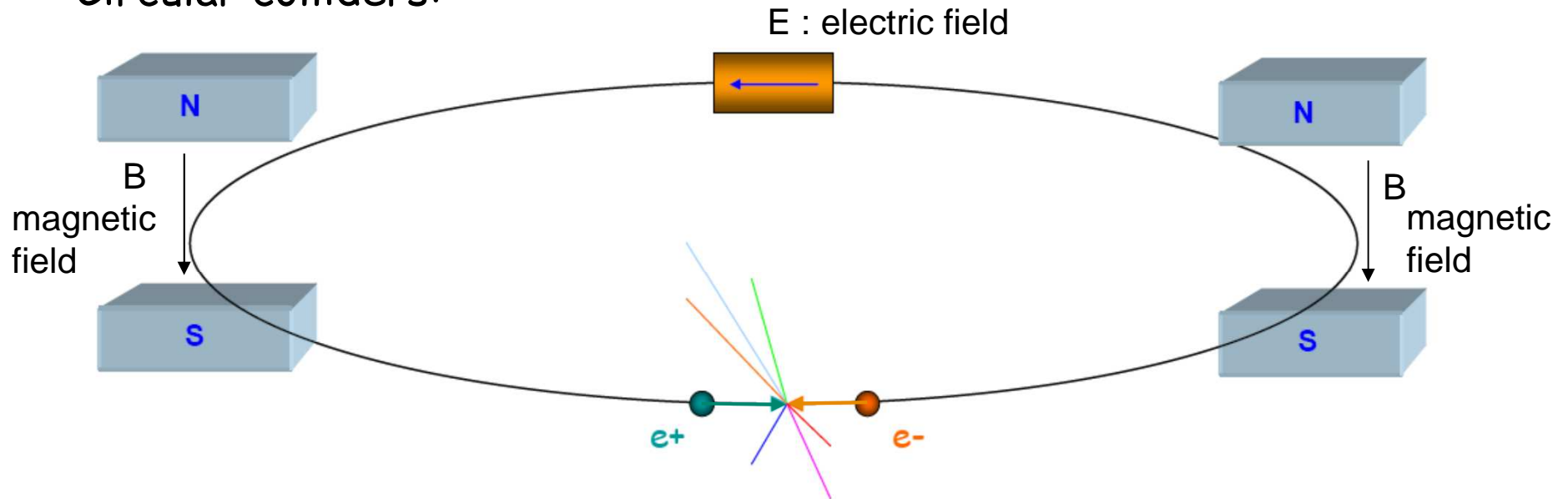


proton
accelerator
920 GeV

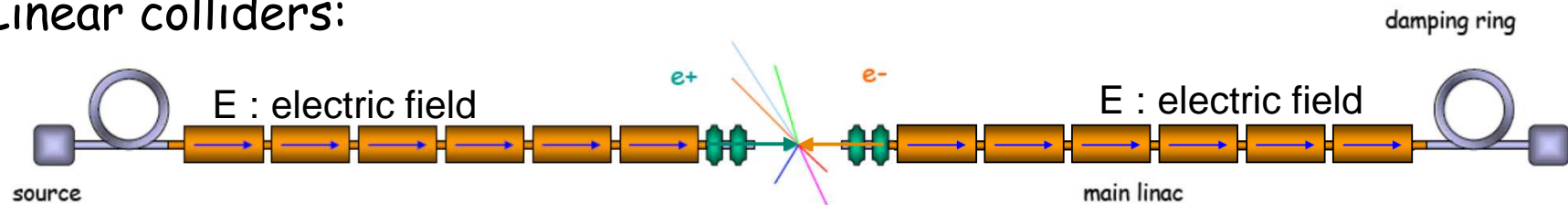
electron accelerator
27.5 GeV

Which collider is better?

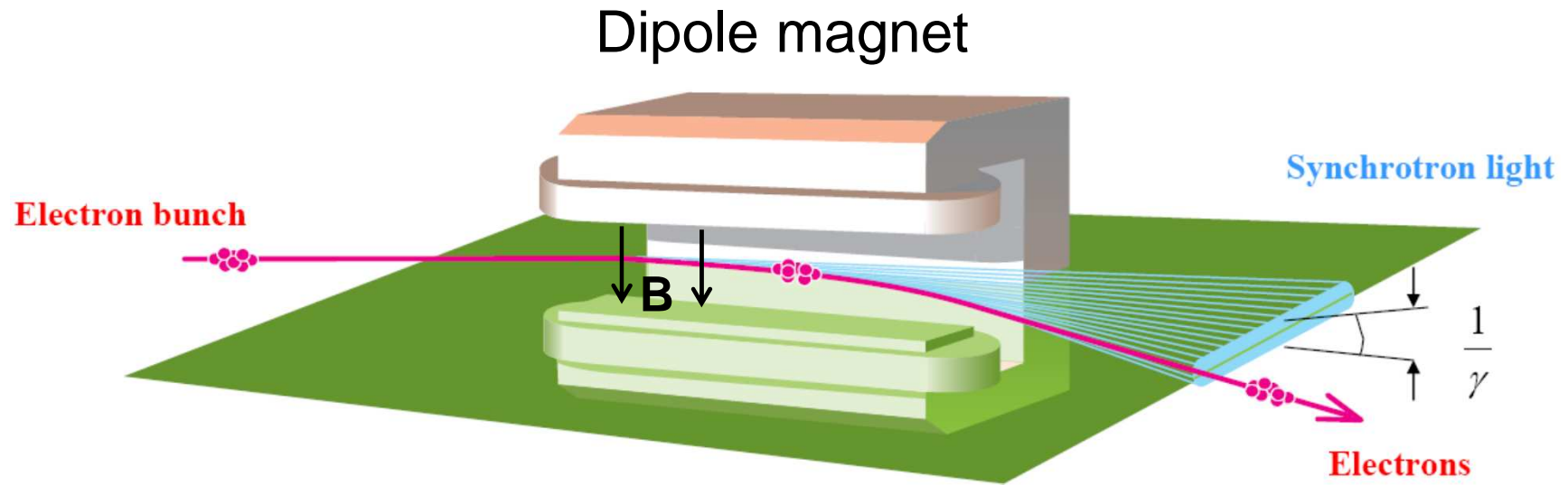
Circular colliders:



Linear colliders:



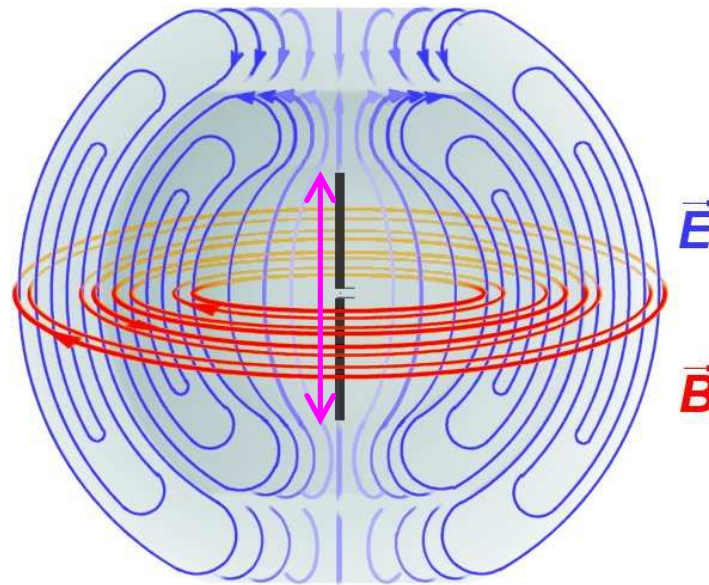
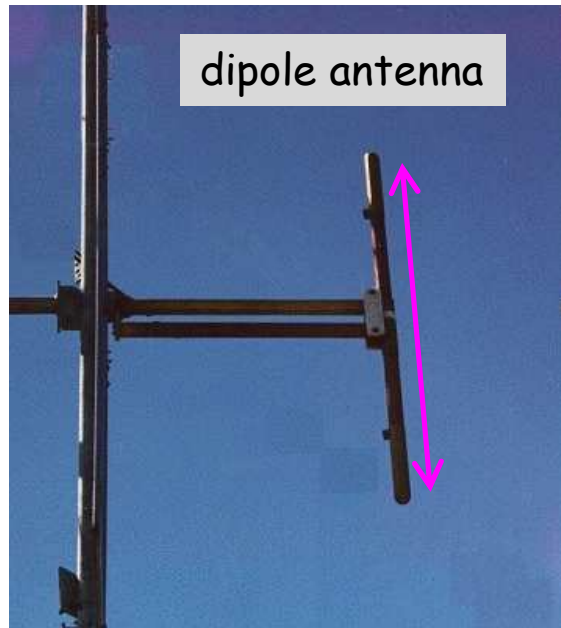
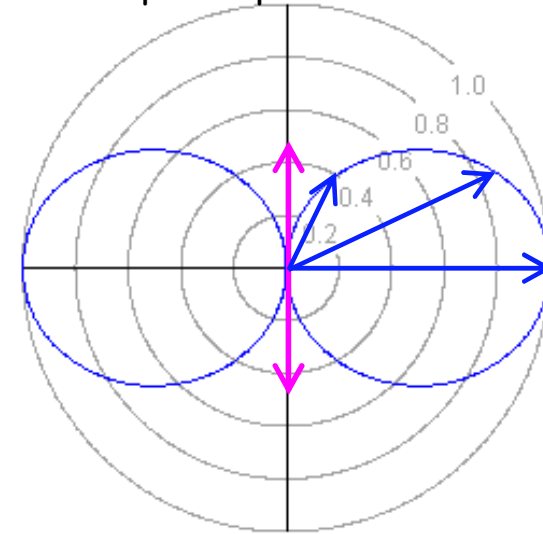
Synchrotron radiation



Radio antenna



radiation power pattern:



Radiation of a dipole antenna

local oscillator:

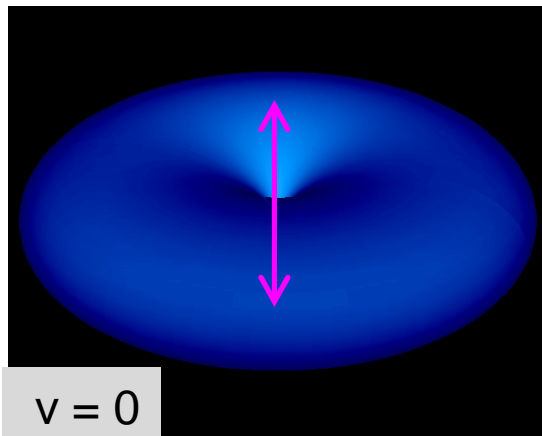
$$P = \frac{q^2 a^2}{12\pi\epsilon_0 c^3} \omega^4$$

(oscillation amplitude: $a < \lambda$)

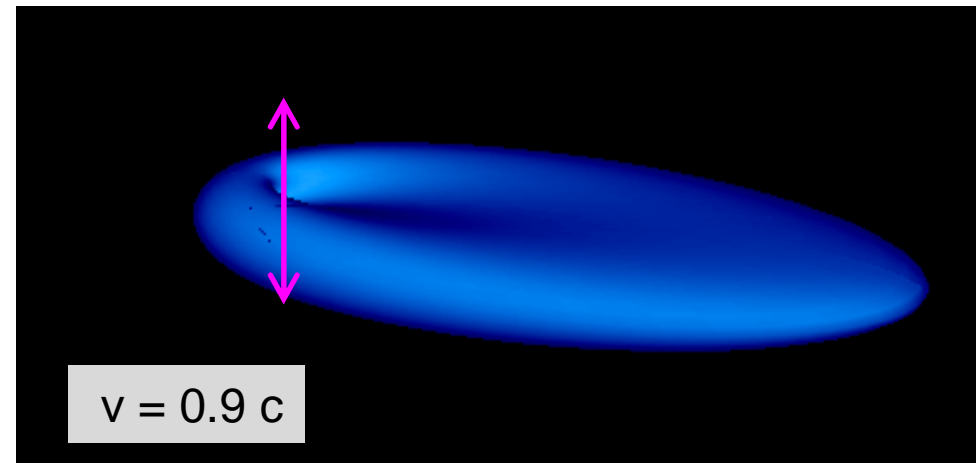
moving oscillator:

$$P = \frac{q^2 a^2}{12\pi\epsilon_0 c^3} \gamma^4 \omega^4$$
$$\gamma = \frac{E}{m_0 c^2}$$

Radiation of an oscillating dipole



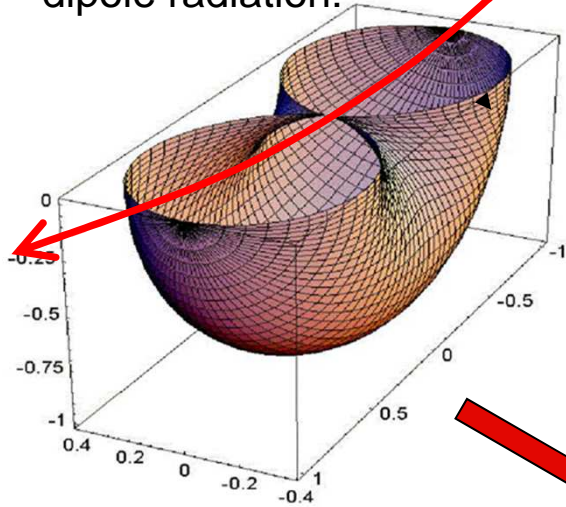
Radiation of a moving oscillating dipole



Radiation of a oscillating dipole under relativistic conditions

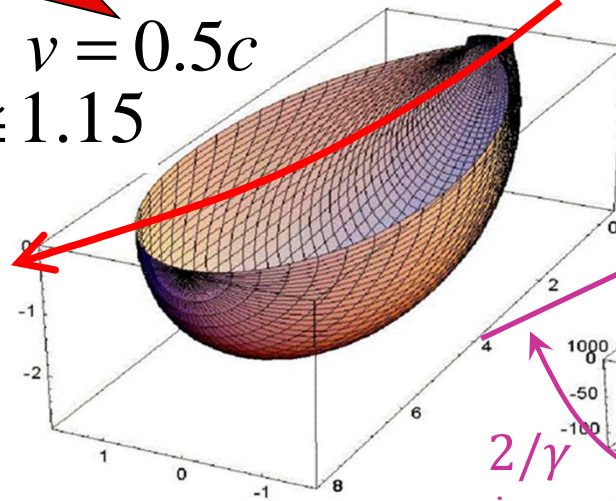
dipole radiation:

electron trajectory

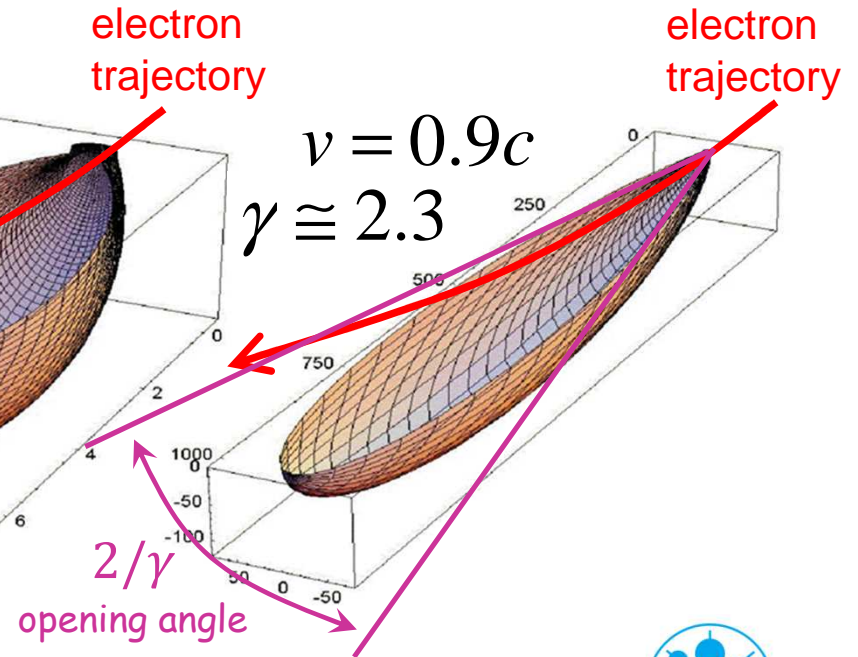


PETRA: $\gamma = 12000$

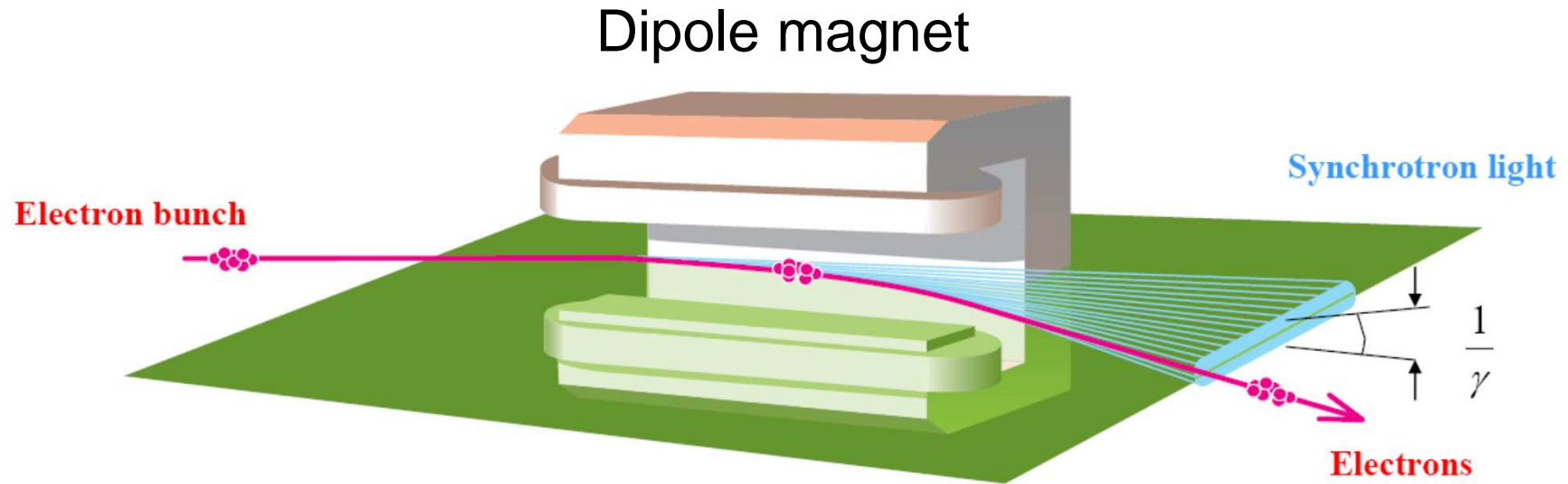
$v = 0.5c$
 $\gamma \cong 1.15$



$v = 0.9c$
 $\gamma \cong 2.3$



Synchrotron radiation



Power radiated by one electron in a dipole field B :

$$\left(\frac{2\pi r}{c} \right) \times$$

$$P = \frac{c q^2}{6\pi \epsilon_0} \frac{\gamma^4}{r^2}$$

$$\gamma = \frac{E}{m_0 c^2}$$

$$\frac{1}{r} = \frac{q B}{p}$$

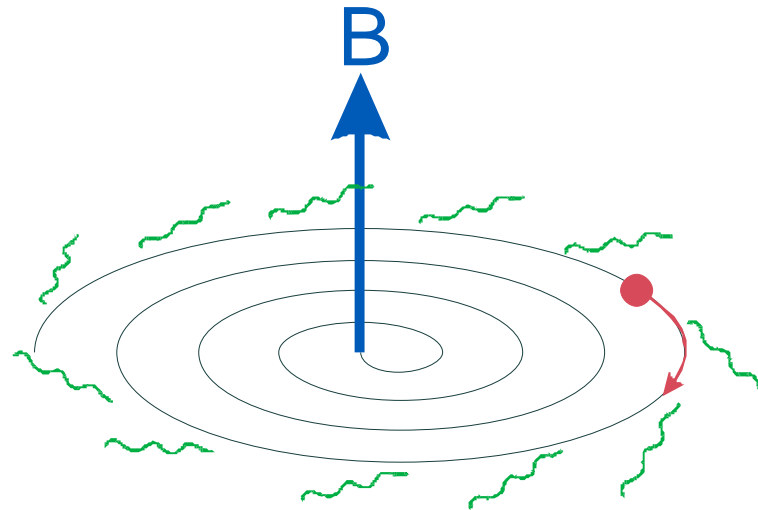
vacuum permittivity



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \Rightarrow \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \quad \Rightarrow \quad \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

$$\gamma = 54000$$

$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

HERA proton ring:

$$r = 580 \text{ m}$$

$$E = 920 \text{ GeV}$$

$$\gamma = 980$$

$$\Delta E_{\text{turn}} \cong 10 \text{ eV (10}^{-9}\text{\%)}$$

← same →

need acceleration = 87 MV per turn



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \quad \Rightarrow \quad \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

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HERA proton ring:

$$r = 580 \text{ m}$$

$$E = 920 \text{ GeV}$$

$$\gamma = 980$$

← same →

the limit is the max. dipole field = 5.5 Tesla

$$\frac{1}{r} = \frac{qB}{p} \quad \Rightarrow \quad p_{\text{max}} = RqB_{\text{max}}$$



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \Rightarrow \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

$$\gamma = 54000$$

$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

need acceleration = 87 MV per turn

LEP collider:

$$r = 2800 \text{ m}$$

$$E = 105 \text{ GeV}$$

$$\gamma = 205000$$

$$\Delta E_{\text{turn}} \cong 4 \text{ GeV (4\%)}$$

need 4 GV per turn !!



Summing-up

Basics of synchrotron radiation

particle type	limitation
<ul style="list-style-type: none">• proton synchrotrons	dipole magnet
<ul style="list-style-type: none">• electron synchrotrons	synchrotron radiation



International Linear Collider (ILC)

Colliding beams with $E_{cm} = 500 \text{ GeV}$ (update to 1 TeV possible)

(not to scale)

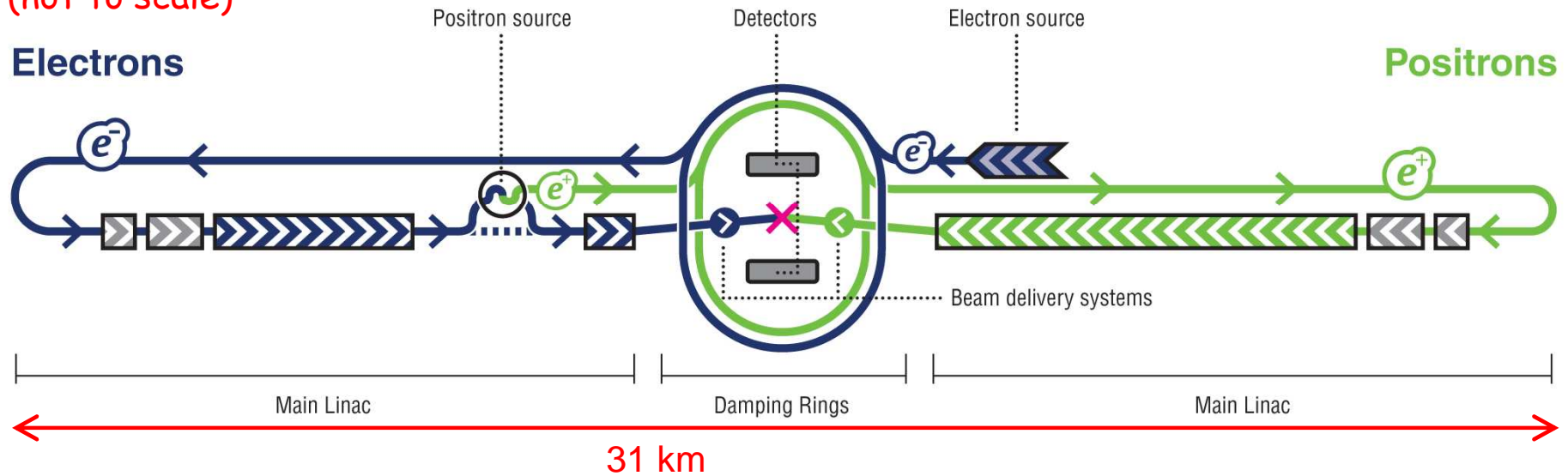
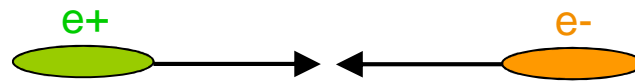
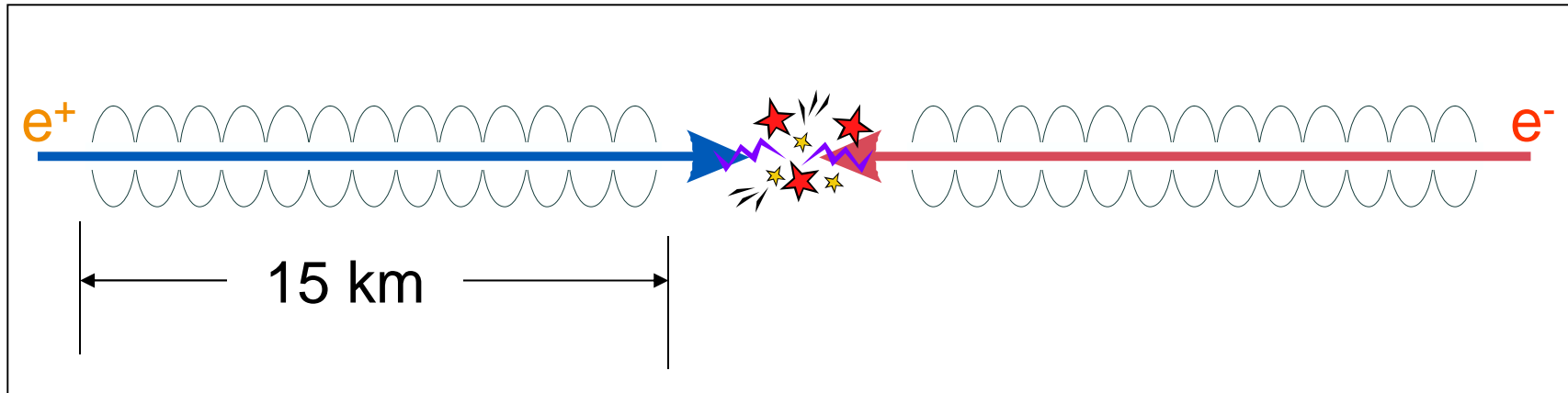


Figure of merit: Luminosity



production rate of a given event (for example, Z particle production):

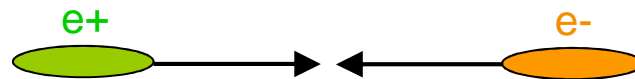
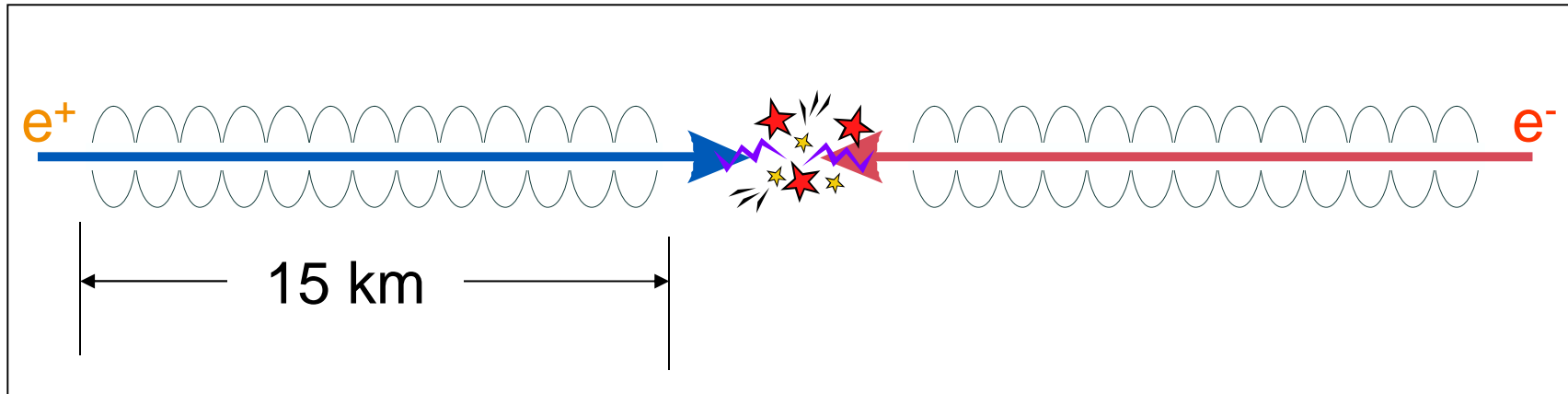
$$R_Z = \frac{dN_Z}{dt} = \sum_z \sigma_z \cdot L$$

number of events

cross section of Z production

luminosity (independent of the event type)

Luminosity



number of colliding bunches per second

number of positrons per bunch

number of electrons per bunch

$$R_Z = \sum_z \cdot L = \sum_z \cdot \frac{f_b N_{e^+} N_{e^-}}{4\pi \sigma_x^* \sigma_y^*} \cdot H_D$$

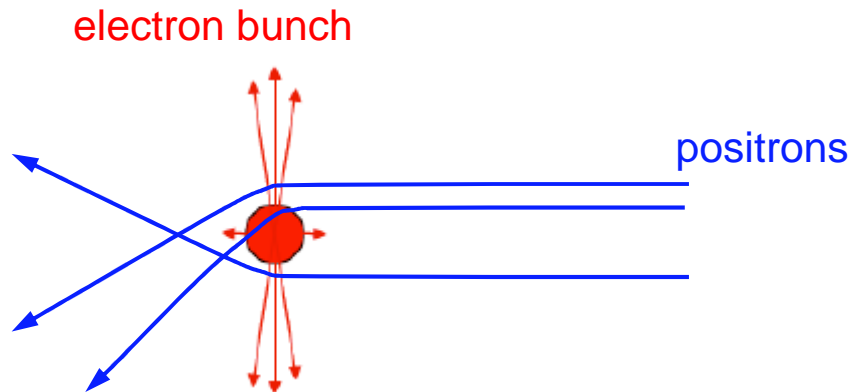
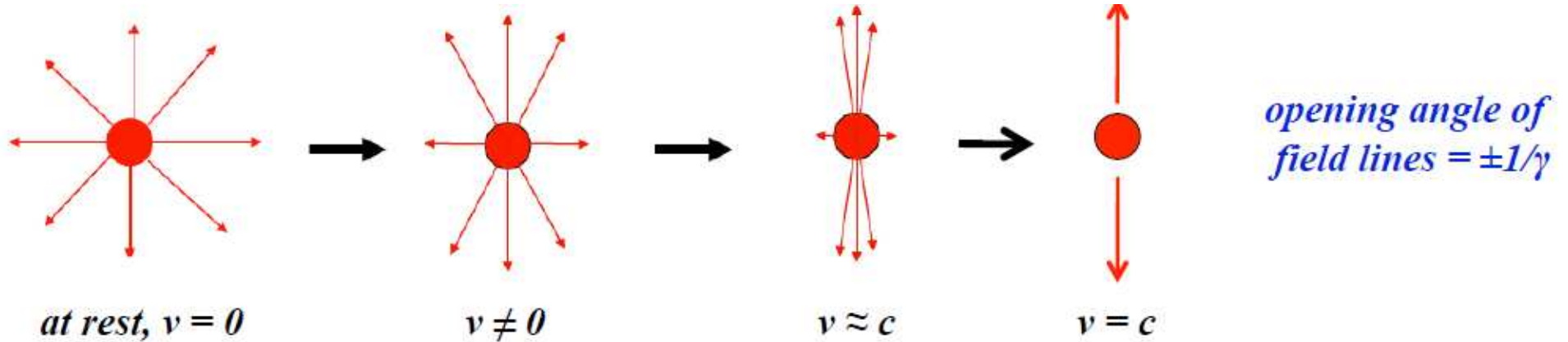
luminosity

beam-beam enhancement factor

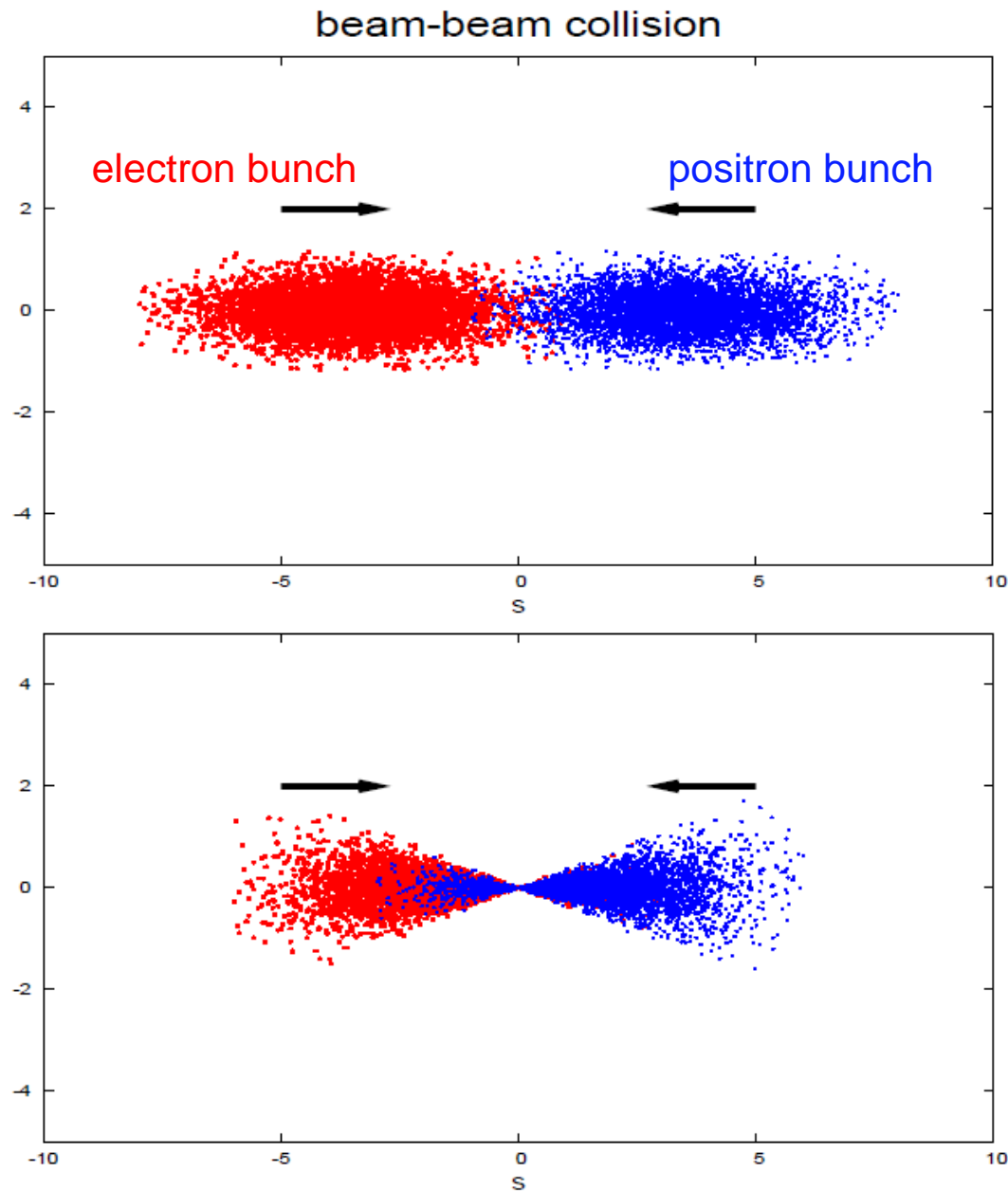
transverse bunch sizes (at the collision point *)

Luminosity enhancement factor H_D due to focusing of opposite beam

electric field of a charged particle (or bunch)



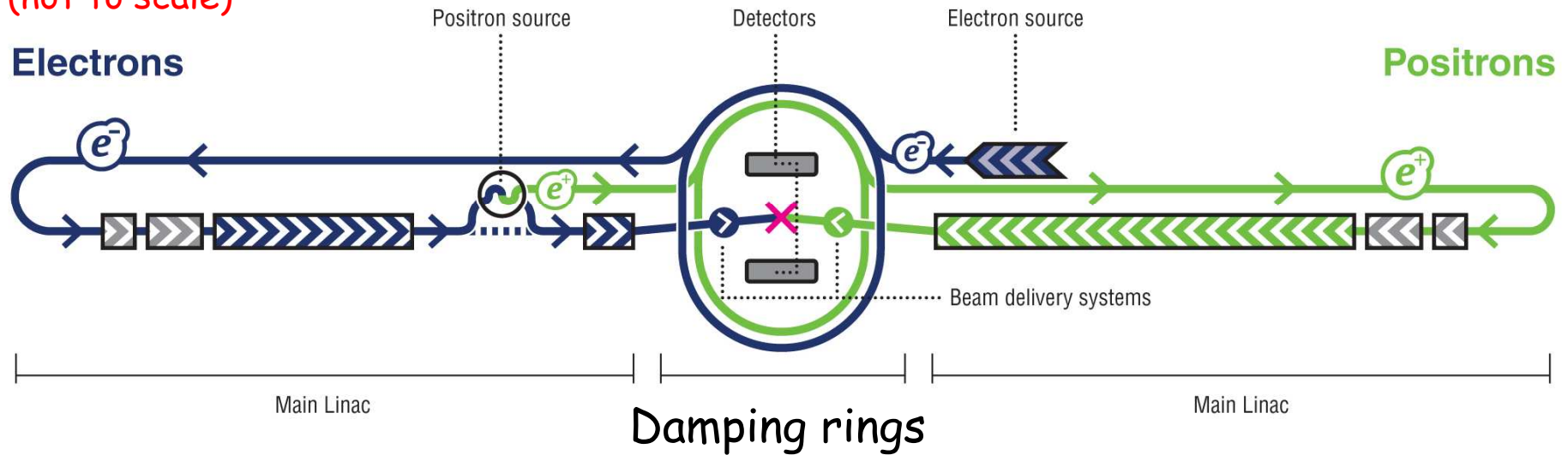
Luminosity enhancement factor H_D due to focusing of opposite beam



International Linear Collider (ILC)

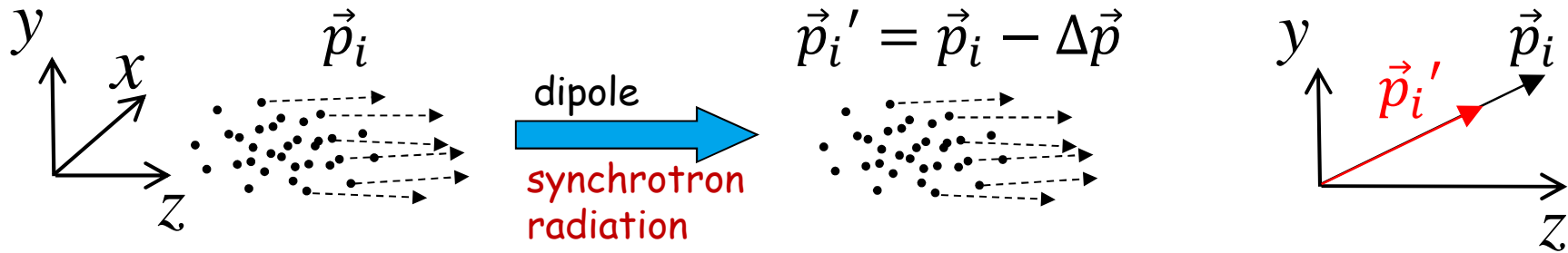
Colliding beams with $E_{cm} = 500 \text{ GeV}$ (update to 1 TeV possible)

(not to scale)

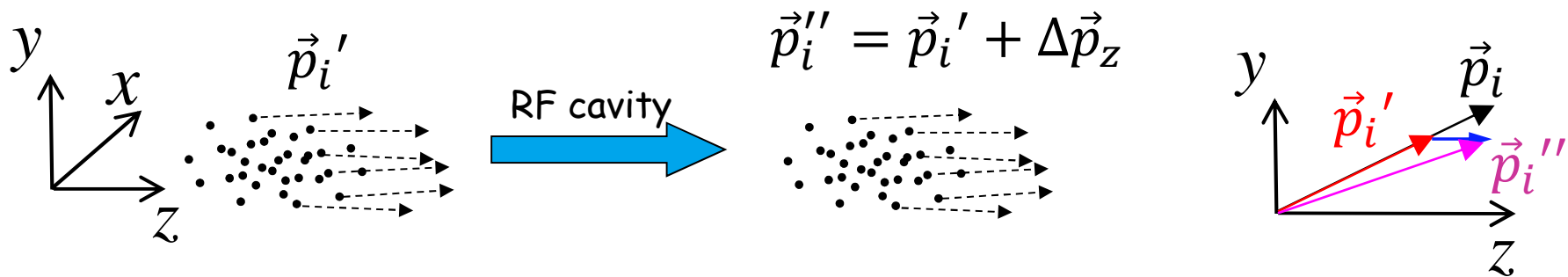


Damping rings

Radiation damping:



Longitudinal acceleration:

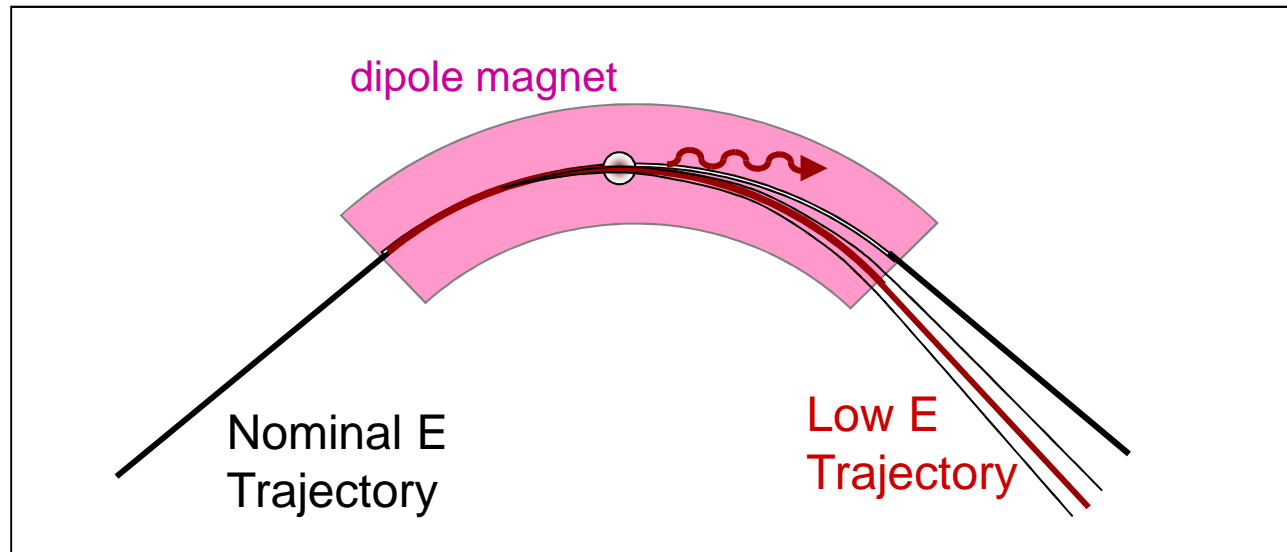


acceleration (only in z direction)

$\frac{p_x}{|\vec{p}|}$ and $\frac{p_y}{|\vec{p}|}$ get smaller



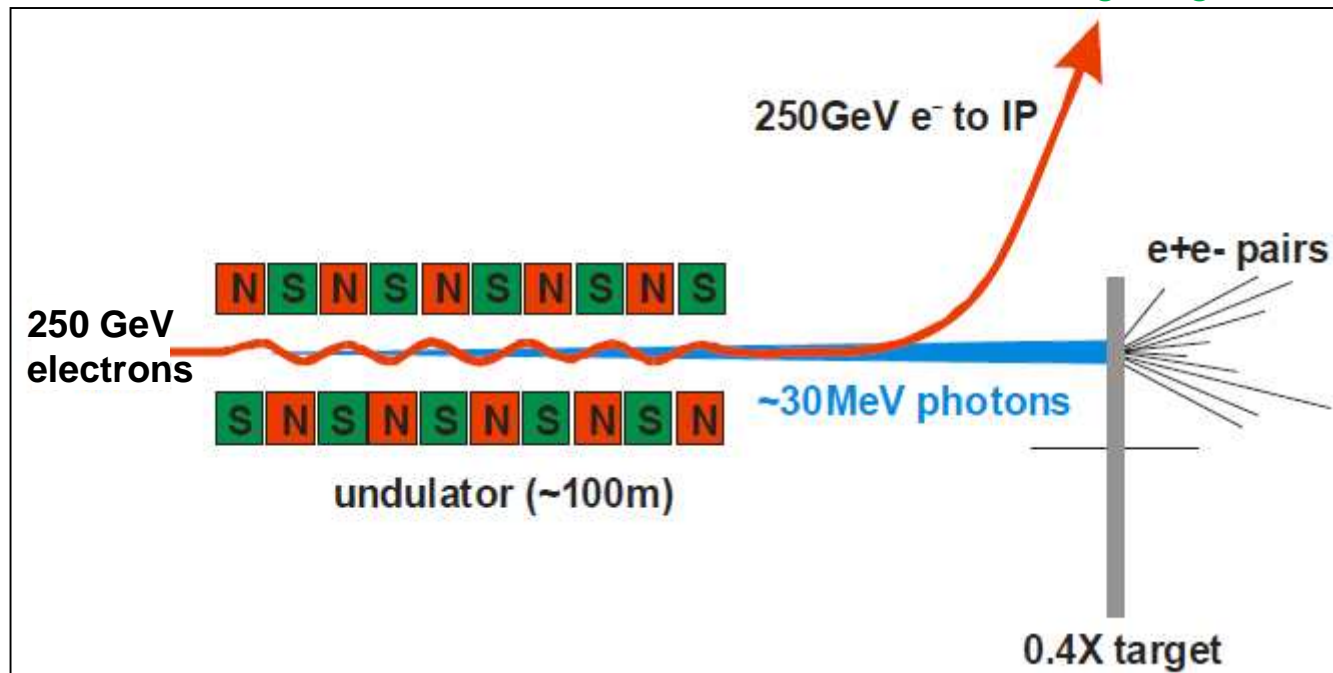
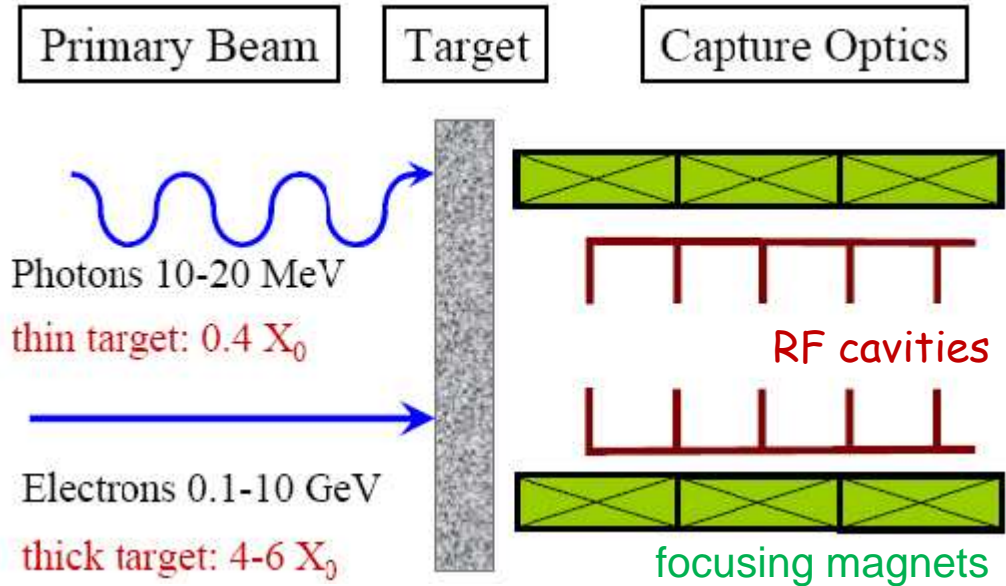
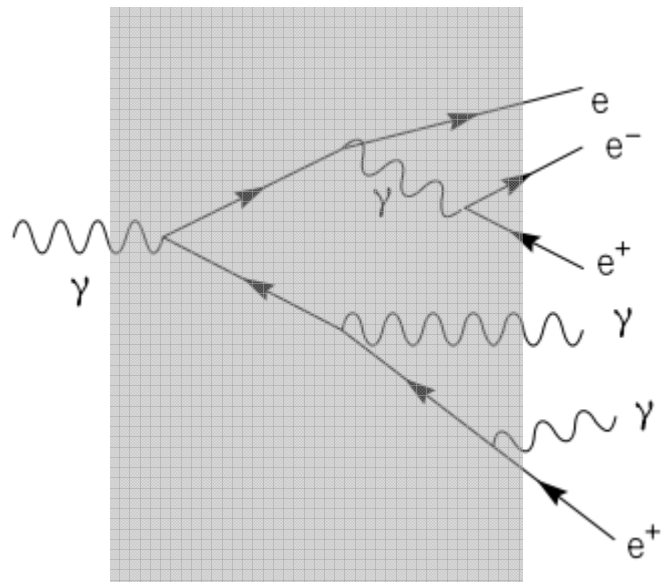
Quantum excitation



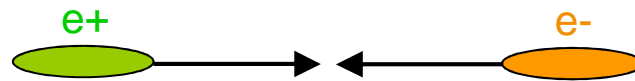
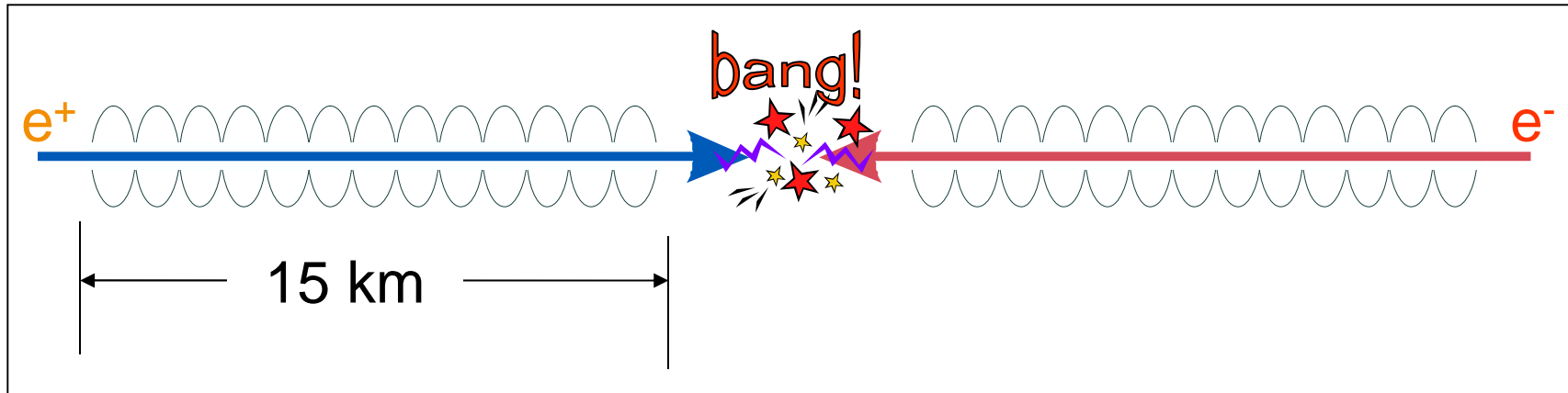
> Quantum excitation

- Radiation is emitted in discrete quanta
- Number and energy distribution etc. of photons obey statistical laws
- → Increase beam size

Positron source



Luminosity



$$L = \frac{f_b N^2}{4\pi \sigma_x^* \sigma_y^*} \cdot H_D = \frac{\eta_{RF \rightarrow beam} P_{RF} N}{4\pi \sigma_x^* \sigma_y^* E_{cm}} \cdot H_D$$

beam power

RF power

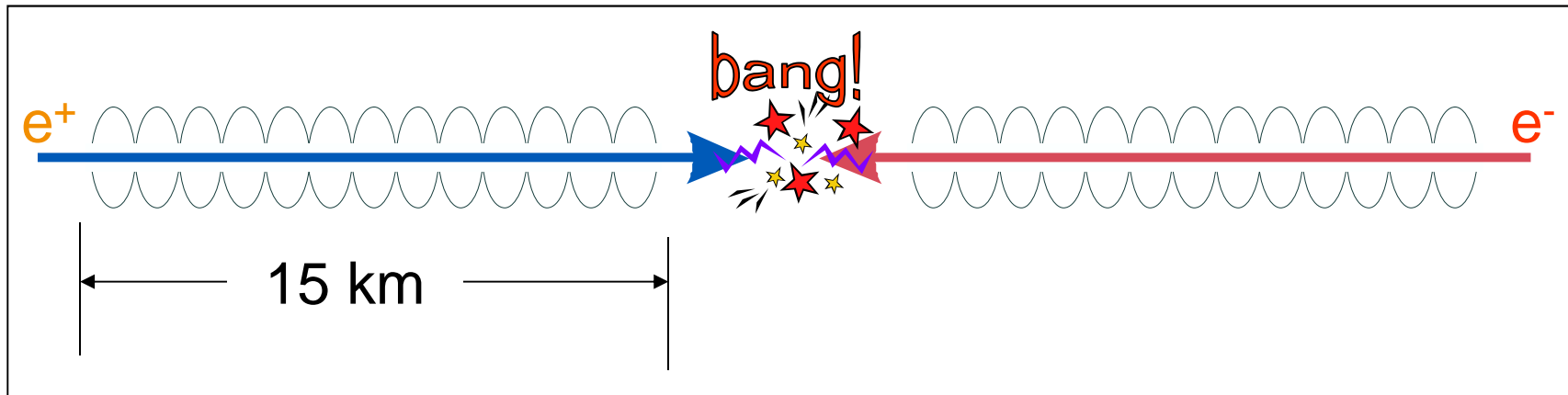
$$f_b N E_{cm} = P_{beam} = \eta_{RF \rightarrow beam} P_{RF}$$

conversion efficiency



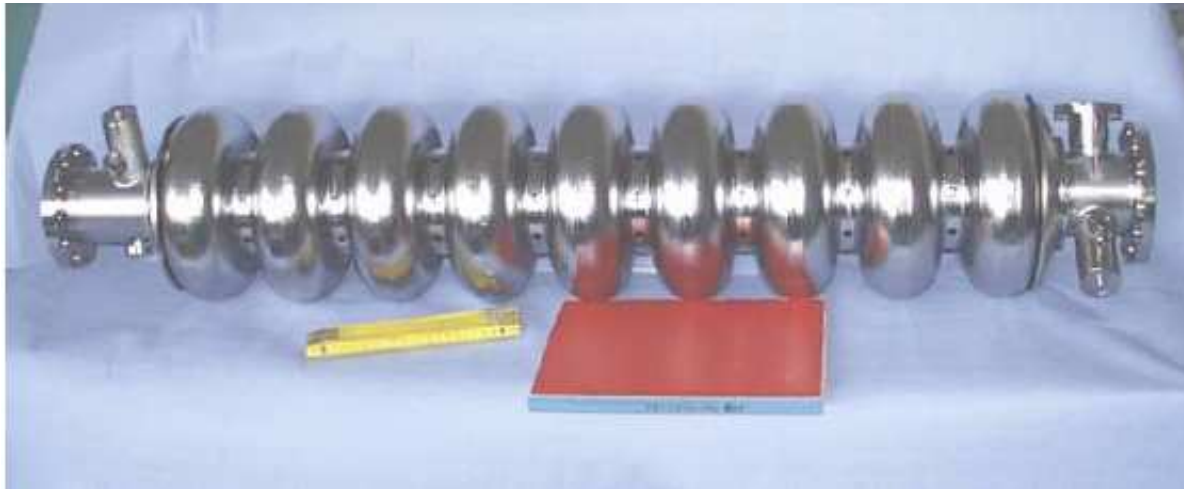
Project for a future e^-e^+ collider: ILC

The International Linear Collider

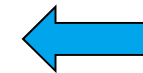
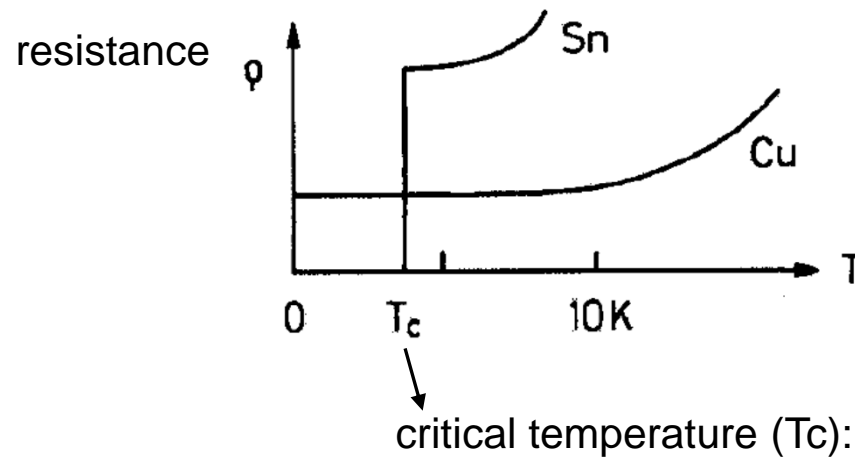


Colliding beams with $E_{CM} = 500 \text{ GeV}$

using superconducting cavities for acceleration:



Advantages of RF superconductivity



for DC currents !

at radio-frequencies, there is a "microwave surface resistance"
which typically is 5 orders of magnitude lower than R of copper

Advantages of RF superconductivity

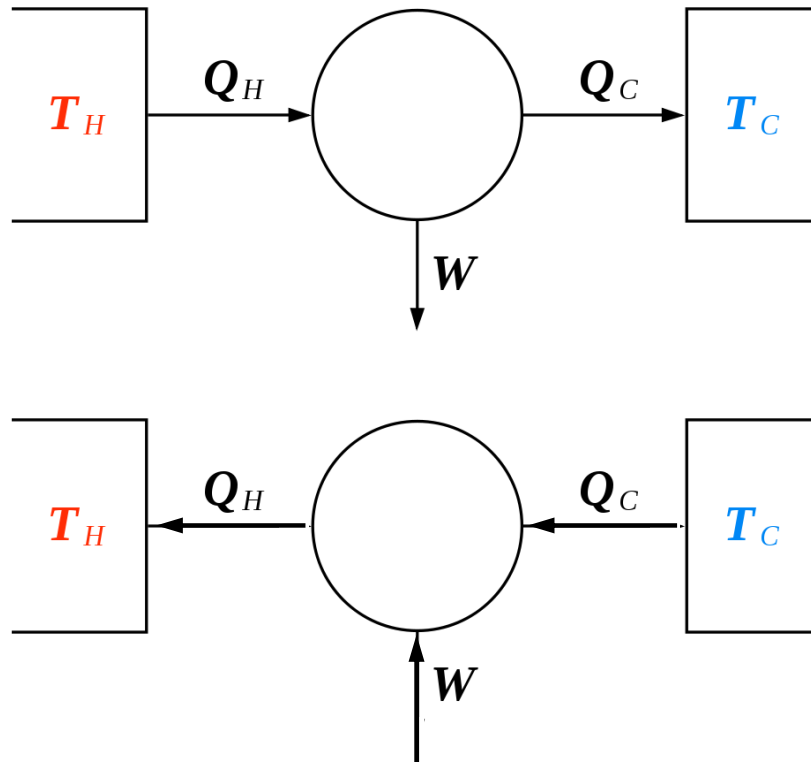
Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
dynamic losses			
for $E = 1$ MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
static losses			



2nd law of Thermodynamics

“Heat cannot spontaneously flow from a colder location to a hotter location”



max. efficiency

$$\eta_c = \frac{T_H - T_C}{T_H}$$

Carnot efficiency:

$$\eta_c = \frac{T_C}{T_H - T_C}$$

most common applications

thermal power stations,
cars, ...

air conditioners,
refrigerators, ...



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for E = 1 MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
Carnot efficiency: $\eta_c = \frac{T}{300 - T} = 0.007$ x			cryogenics efficiency 20-30%
for E = 1 MV/m	1 kW/m	0.4 W/m	60 kW / m
	electric power	RF power lost	



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for $E = 1$ MV/m	$0.4 + 1$ W / m at 2 K	60 kW / m	dissipated at the cavity walls
for $E = 1$ MV/m	1 kW/m 0.4 W/m	60 kW / m	
for $E = 1$ MV/m	1 kW/m 0.8 W/m	120 kW / m	including RF generation efficiency (50%)
	e-power	e-power for RF	

reduction factor of >100 in (electrical) power



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for E = 1 MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
for E = 1 MV/m	1 kW/m (e-power) 0.4 W/m (RF power)	60 kW / m (RF power)	
for E = 1 MV/m and 20 mA beam	1 kW/m (e-power) 0.4 W/m (RF power) + 20 kW/m (beam)	60 kW / m (RF power) + 20 kW / m (beam)	
$\eta_{RF \rightarrow beam}$	~ 100 %	~ 25 %	

$$P_{beam} = \eta_{RF \rightarrow beam} P_{RF}$$



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for E = 1 MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
for E = 1 MV/m	1 kW/m (e-power) 0.4 W/m (RF power)	60 kW / m (RF power)	
for E = 1 MV/m and 20 mA beam	1 kW/m (e-power) 0.4 W/m (RF power) + 20 kW/m (beam)	60 kW / m (RF power) + 20 kW / m (beam)	
for E = 1 MV/m and 20 mA beam	1 kW/m (e-power) 40 kW/m (e-power for RF)	160 kW / m (RF power)	including RF generation efficiency (50%)

reduction factor of 4 in (electrical) power





Accelerator Research at DESY

SC Technology

(R&D on CW, cryo module test bench, ongoing)

SC Processes

(large series production & testing of sc cavities and modules, AMTF, ongoing)

Surface Technology

(cavity surfaces, CRISP, ongoing)

FEL Seeding

(sFLASH, FLASH2 seeding, ongoing)

LAOLA at FLASH2

(beam-driven plasma acceleration, 2016+)

SINBAD

(ultra-short bunches, LAOLA, prototype table-top FEL, 2017+)

AXSIS

(atto-second bunches, ICS, 2014+)

LAOLA at REGAE

(laser-driven plasma acceleration, 2013-2016)

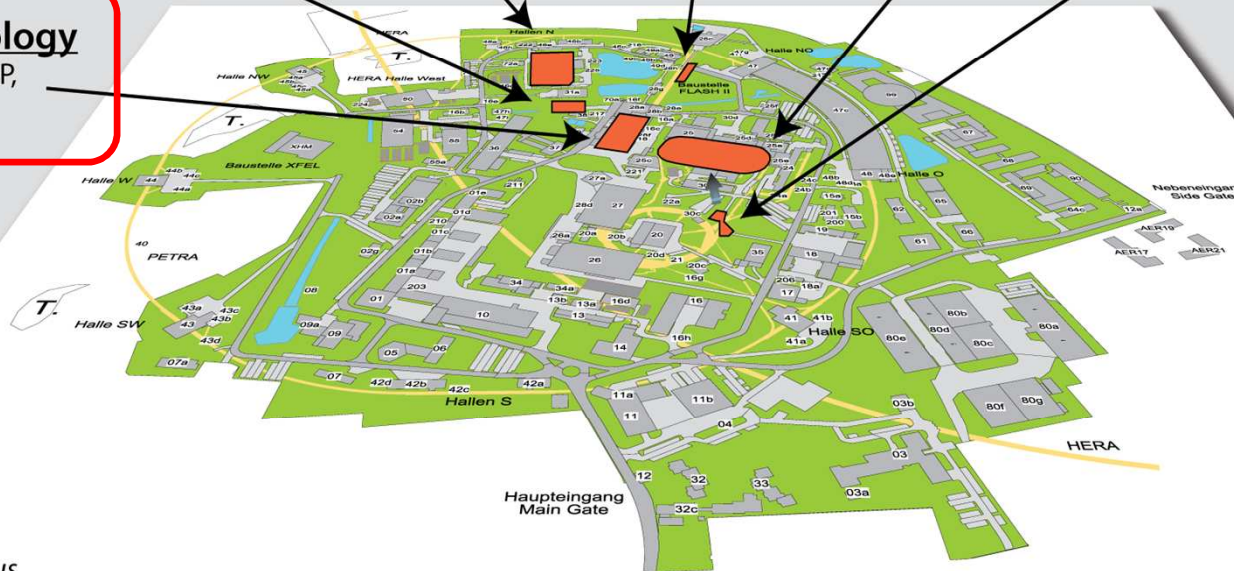
Zeuthen Campus

Photo-Injector

(ongoing)

LAOLA at PITZ

(bunch modulation in plasma, 2013+)



Hamburg Campus





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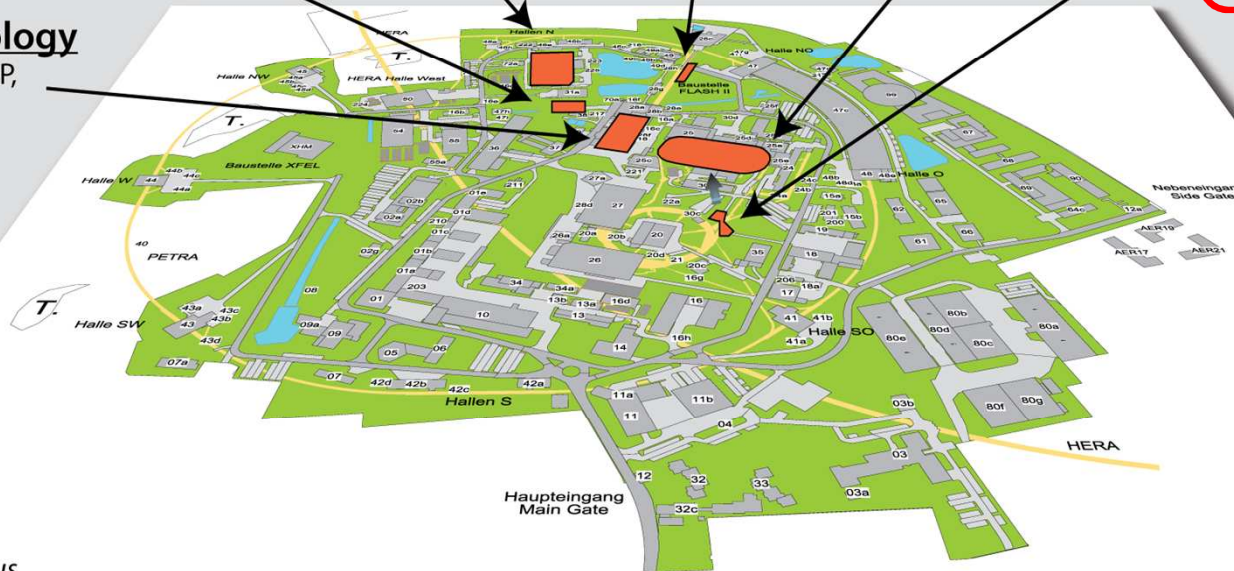
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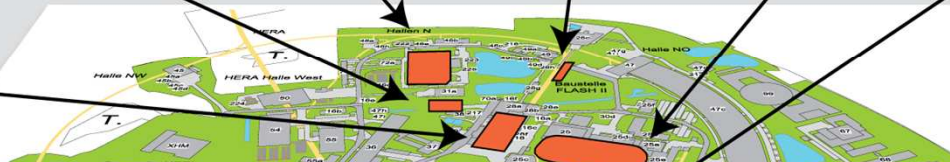
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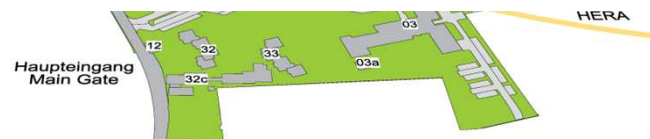


Laboratory for **L**aser- and beam-driven plasma **A**cceleration (**L**ao**L**A)
Relativistic **E**lectron **G**un for **A**tomic **E**xploration (**R**EGAE)
Short **I**nnovative **B**unches and **A**ccelerators at **D**oris (**S**IN**B**AD)
Attosecond **X**-ray **S**cience **I**maging and **S**pectroscopy (**A**X**S**IS)
Photo **I**njector **T**est Facility at DESY, Location **Z**euthen (**P**ITZ)

Zeuthen Campus

Photo-Injector
(ongoing)

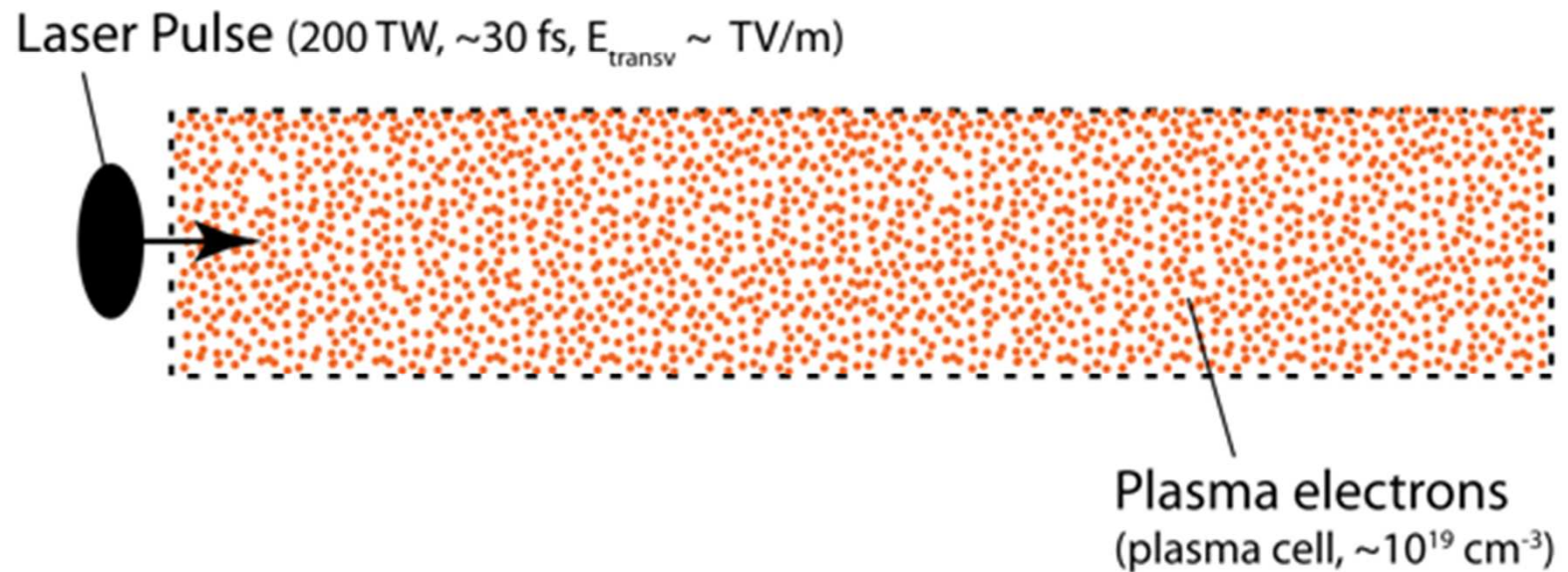
LAOLA at PITZ
(bunch modulation in plasma, 2013+)



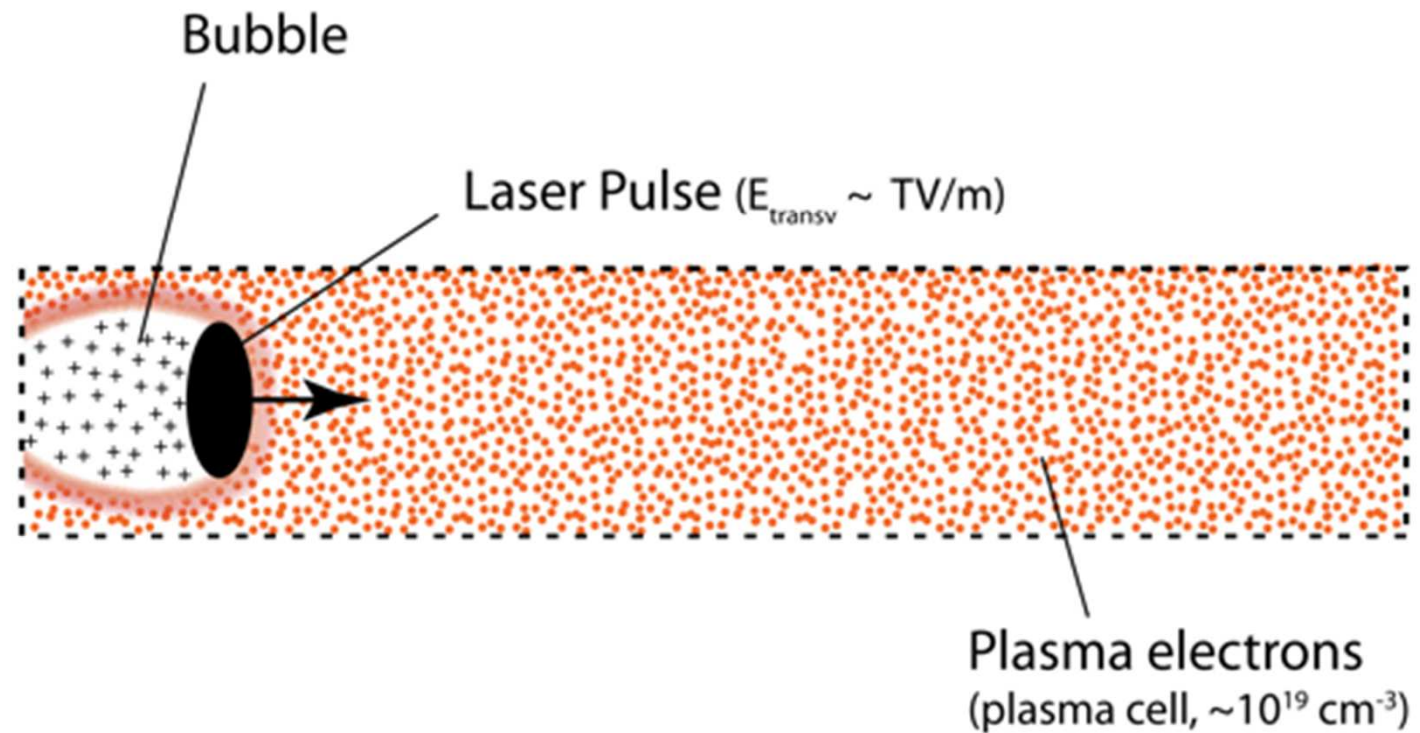
Hamburg Campus



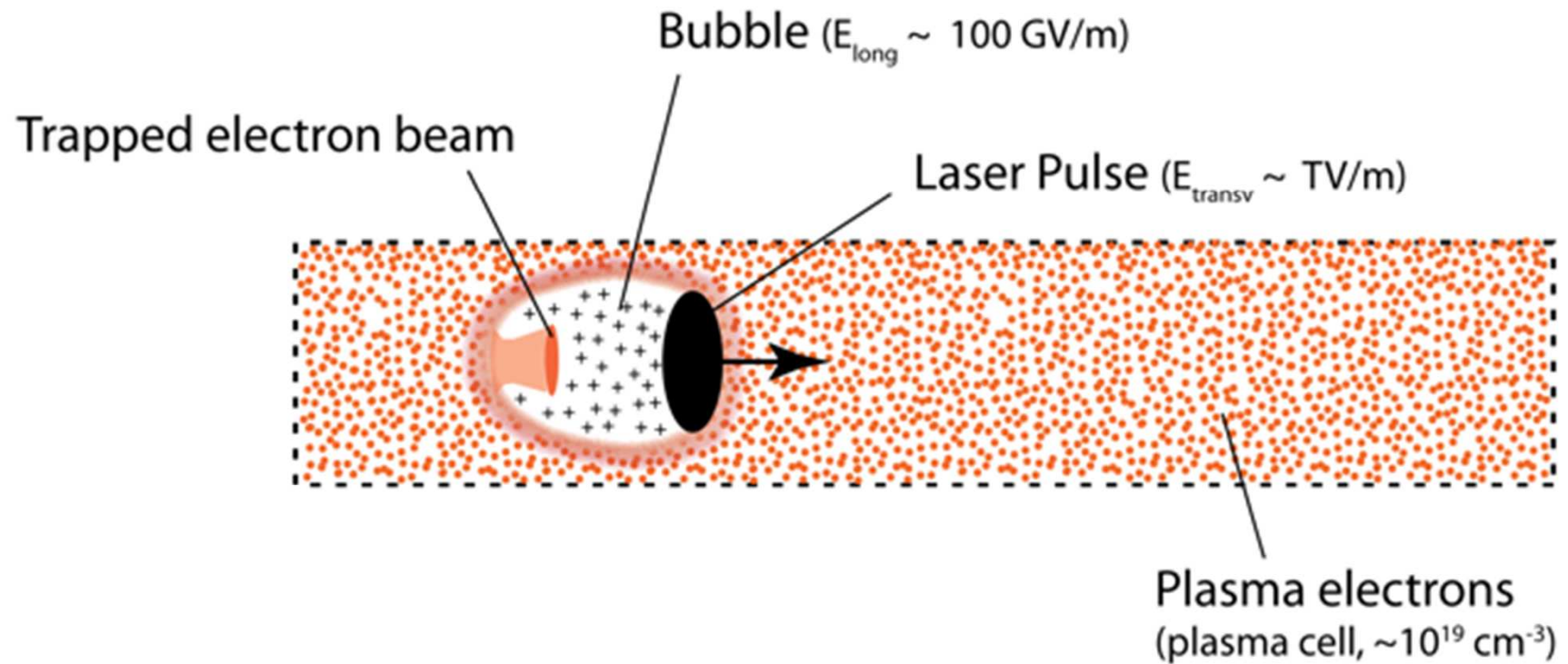
Laser-driven plasma wakefield acceleration



Laser-driven plasma wakefield acceleration



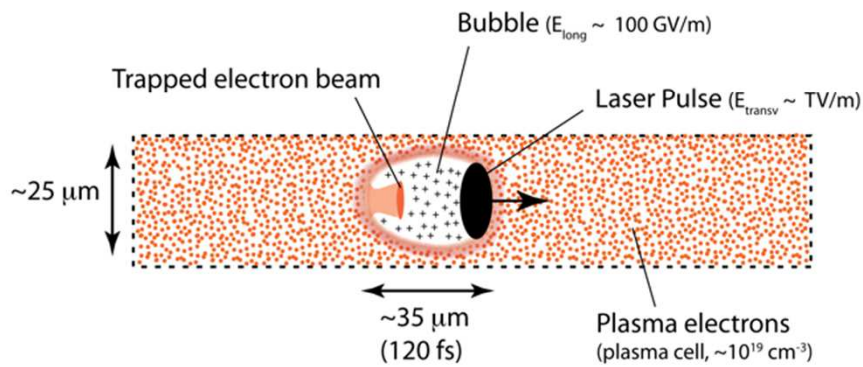
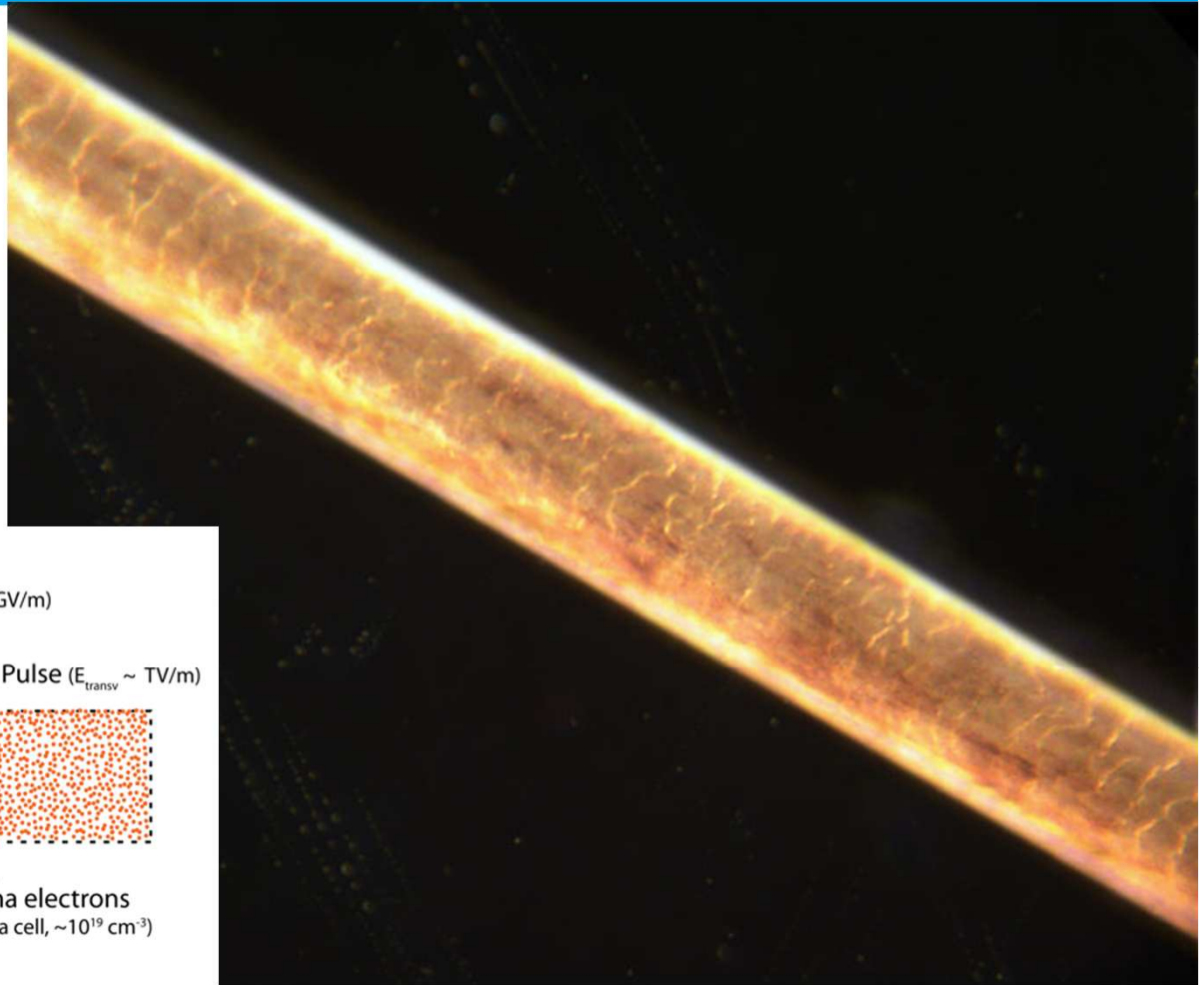
Laser-driven plasma wakefield acceleration



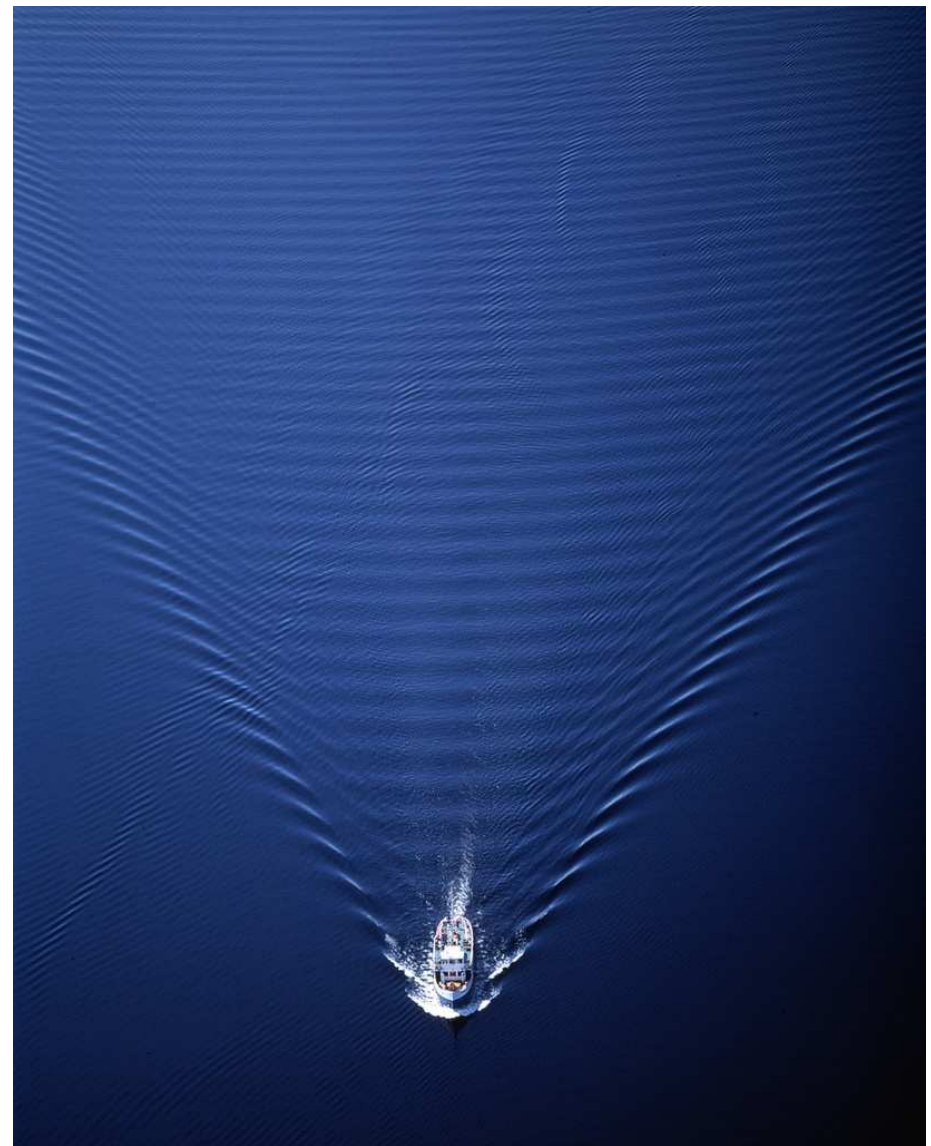
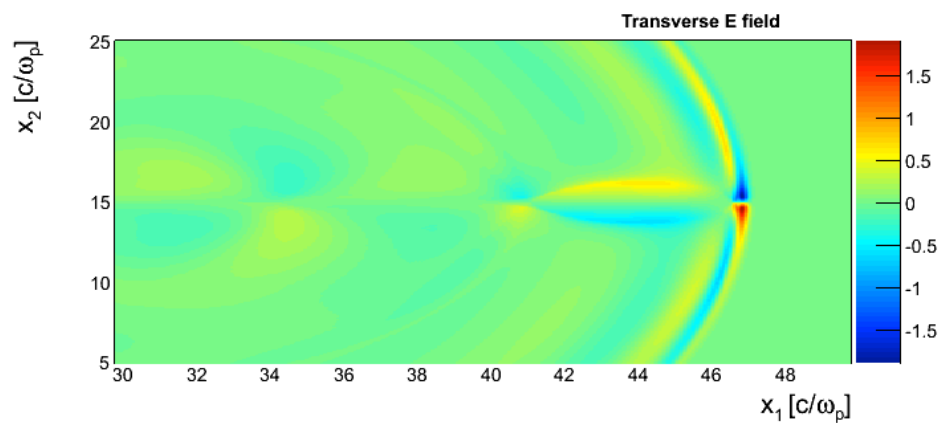
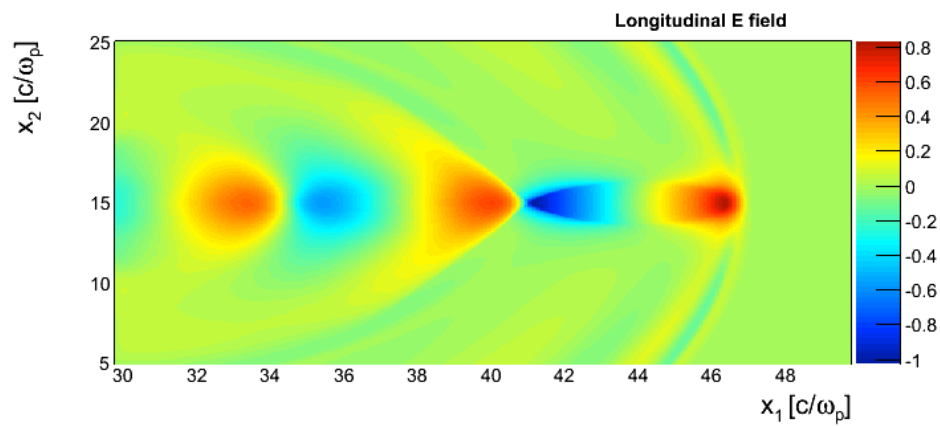
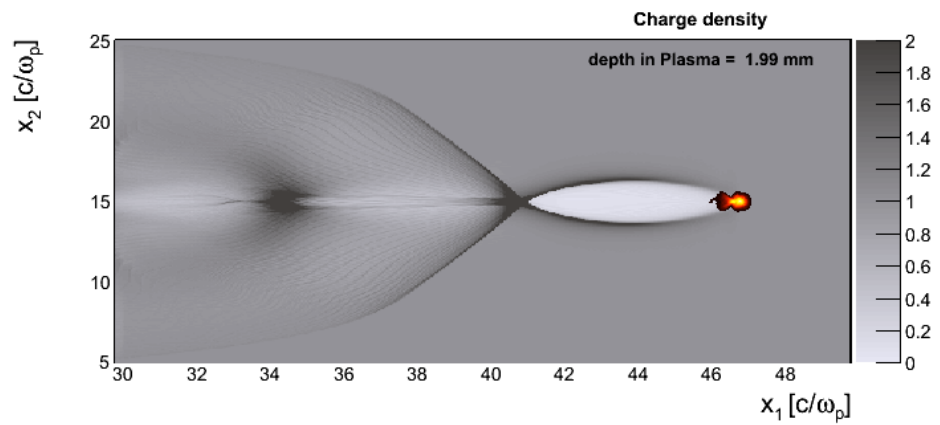
Acceleration (and focusing) within a hair width ...

E_{long} **100 GV/m**

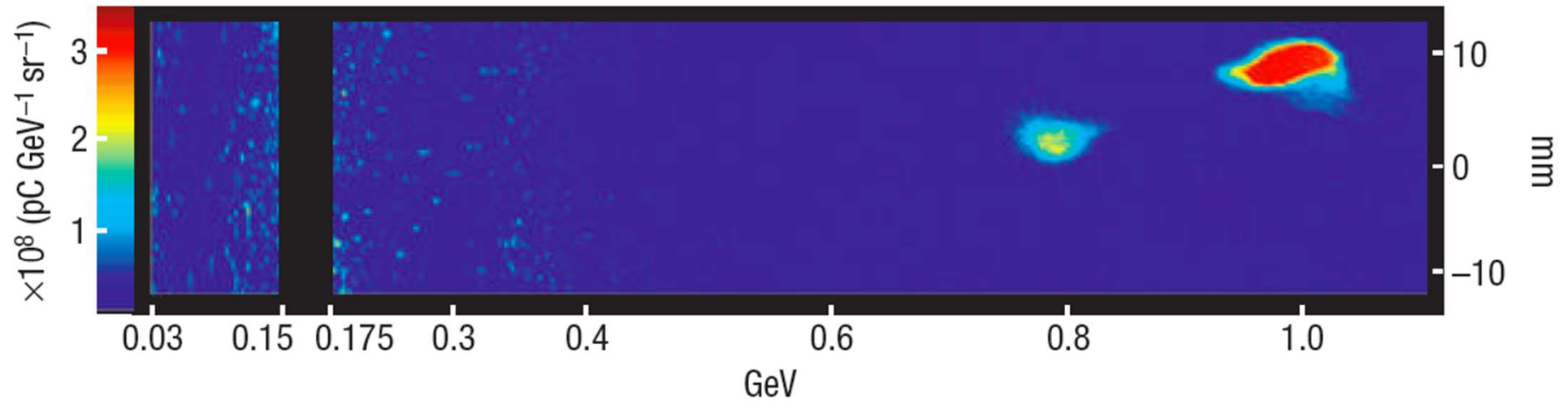
g_{transv} **300 MT/m**



Simulations



Breakthrough in the lab



LETTERS

GeV electron beams from a centimetre-scale accelerator

W. P. LEEMANS^{1*}, B. NAGLER¹, A. J. GONSALVES², Cs. TÓTH¹, K. NAKAMURA^{1,3}, C. G. R. GEDDES¹, E. ESAREY^{1*}, C. B. SCHROEDER¹ AND S. M. HOOKER²

¹Lawrence Berkeley National Laboratory, 1 Cyclotron Road, Berkeley, California 94720, USA

²University of Oxford, Clarendon Laboratory, Parks Road, Oxford OX1 3PU, UK

³Nuclear Professional School, University of Tokyo, 22-2 Shirane-shirakata, Tokai, Naka, Ibaraki 319-1188, Japan

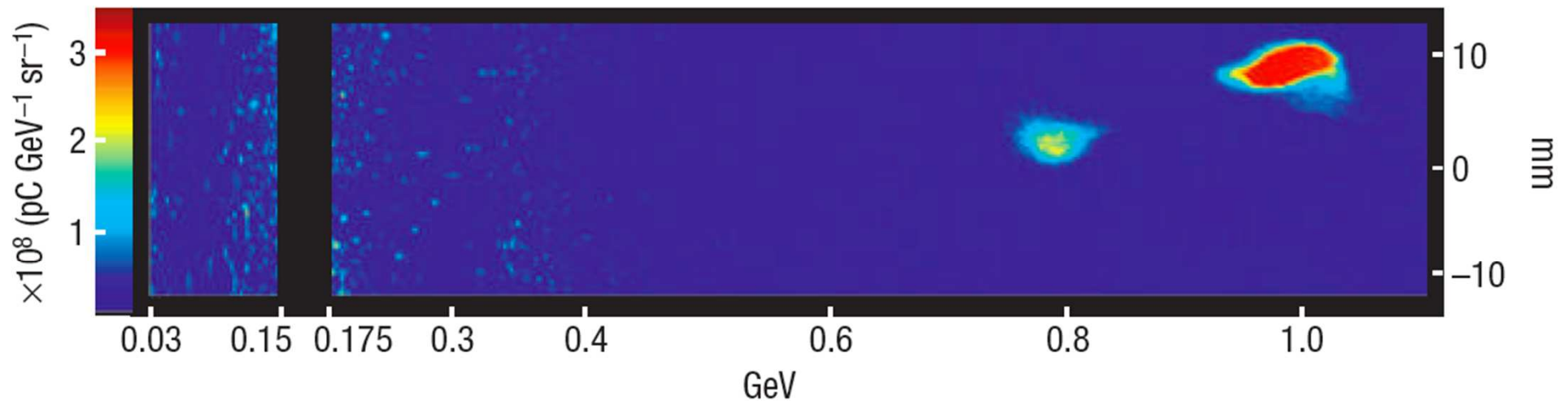
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nature physics | VOL 2 | OCTOBER 2006



Breakthrough in the lab



LETTERS

GeV electron beams from a centimetre-scale accelerator

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2-3 orders of magnitude higher acceleration fields than conventional RF cavities

smaller and cheaper accelerators are possible

nature physics | VOL 2 | OCTOBER 2006



Challenges of laser-driven plasma wakefield acceleration

- plasma self-injection is very unstable in energy, timing ...
- very low efficiency “wall-plug power” to “beam power”

(reminder: power efficiency in RF cavity accelerators)

$$P_{beam} = \eta_{RF \rightarrow beam} P_{RF}$$
$$L = \frac{\eta_{RF \rightarrow beam} P_{RF} N}{4\pi \sigma_x^* \sigma_y^* E_{cm}} \cdot H_D$$

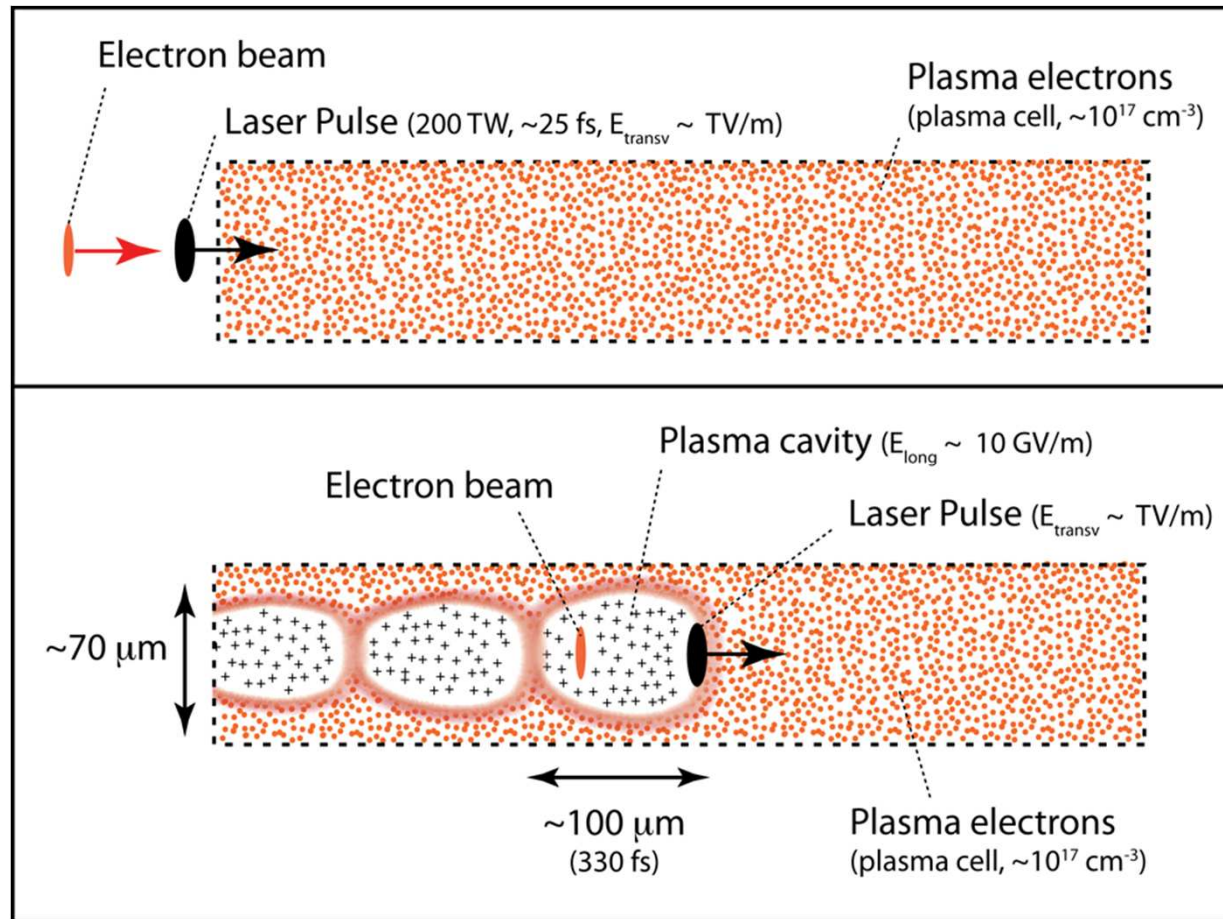


Challenges of laser-driven plasma wakefield acceleration

- plasma self-injection is very unstable in energy, timing ...
- very low efficiency “wall-plug power” to “beam power”
- maximum energy limitation at 1 GeV or few GeV (in self-injection mode)
- multiple-stage plasma acceleration not yet demonstrated

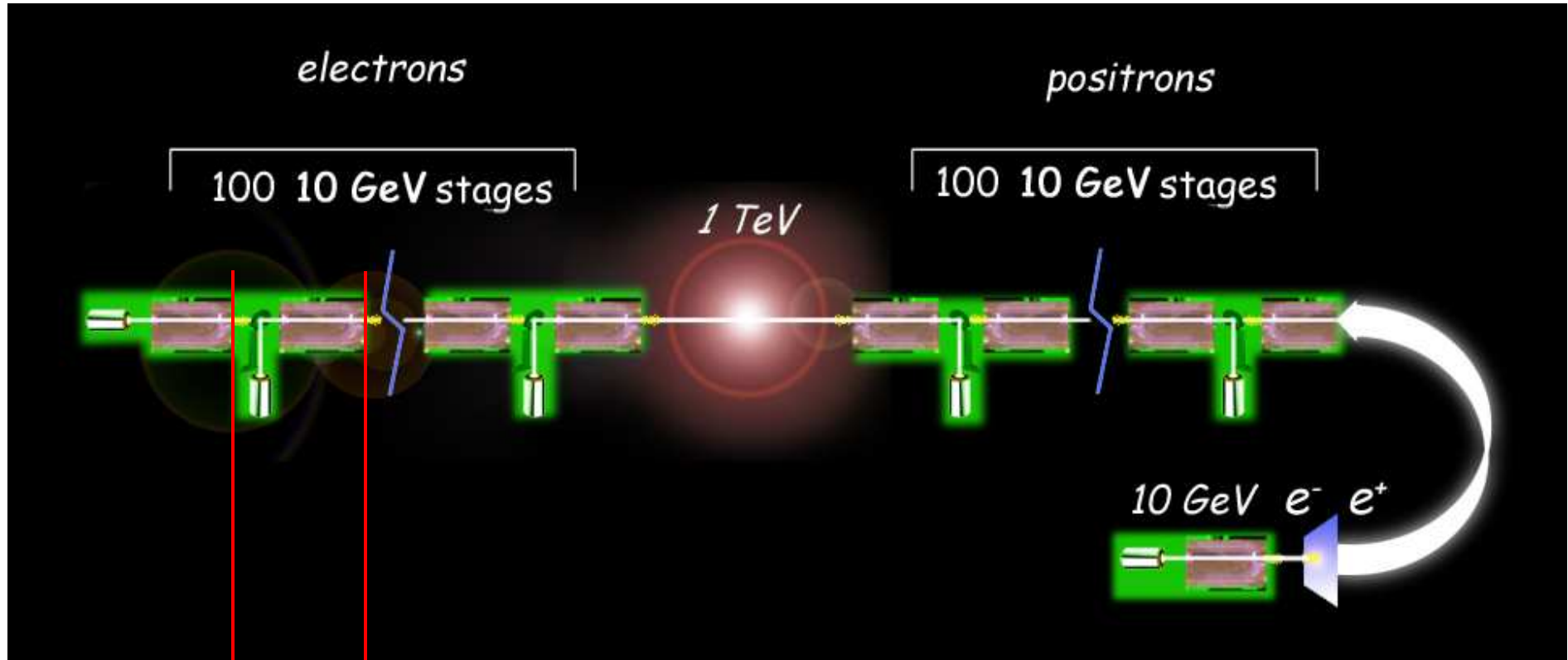


Challenges of laser-driven plasma wakefield acceleration



not yet proven !

Linear collider based on plasma wakefield acceleration



cell length ?

Circular colliders (synchrotrons):

particle type	limitation
<ul style="list-style-type: none">proton synchrotrons $938 \text{ MeV}/c^2$	dipole magnet
<ul style="list-style-type: none">electron synchrotrons $0.511 \text{ MeV}/c^2$	synchrotron radiation
<ul style="list-style-type: none">muon synchrotrons $105.7 \text{ MeV}/c^2$	mean lifetime: $\tau = 2.2 \mu\text{s} ?$ $\tau^* = \gamma \tau = 21 \text{ ms}$ at $E = 1 \text{ TeV}$ travel 6300 km 1000 turns in a synchrotron with $B = 7 \text{ T}$



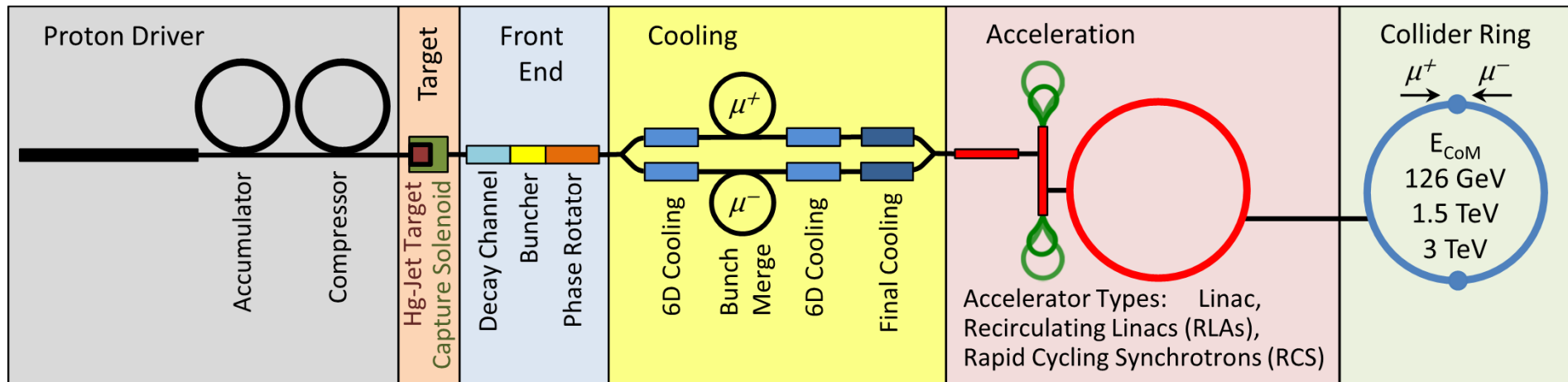
Circular colliders (synchrotrons):

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<ul style="list-style-type: none"> muon synchrotrons 105.7 MeV/c² 	mean lifetime: $\tau = 2.2 \mu\text{s}$ $p \rightarrow \pi \rightarrow \mu$ (muon beams produced as tertiary beams)



Muon collider principle

Muon collider block diagram



Collider:

$$E_{cm} = 3 \text{ TeV}$$

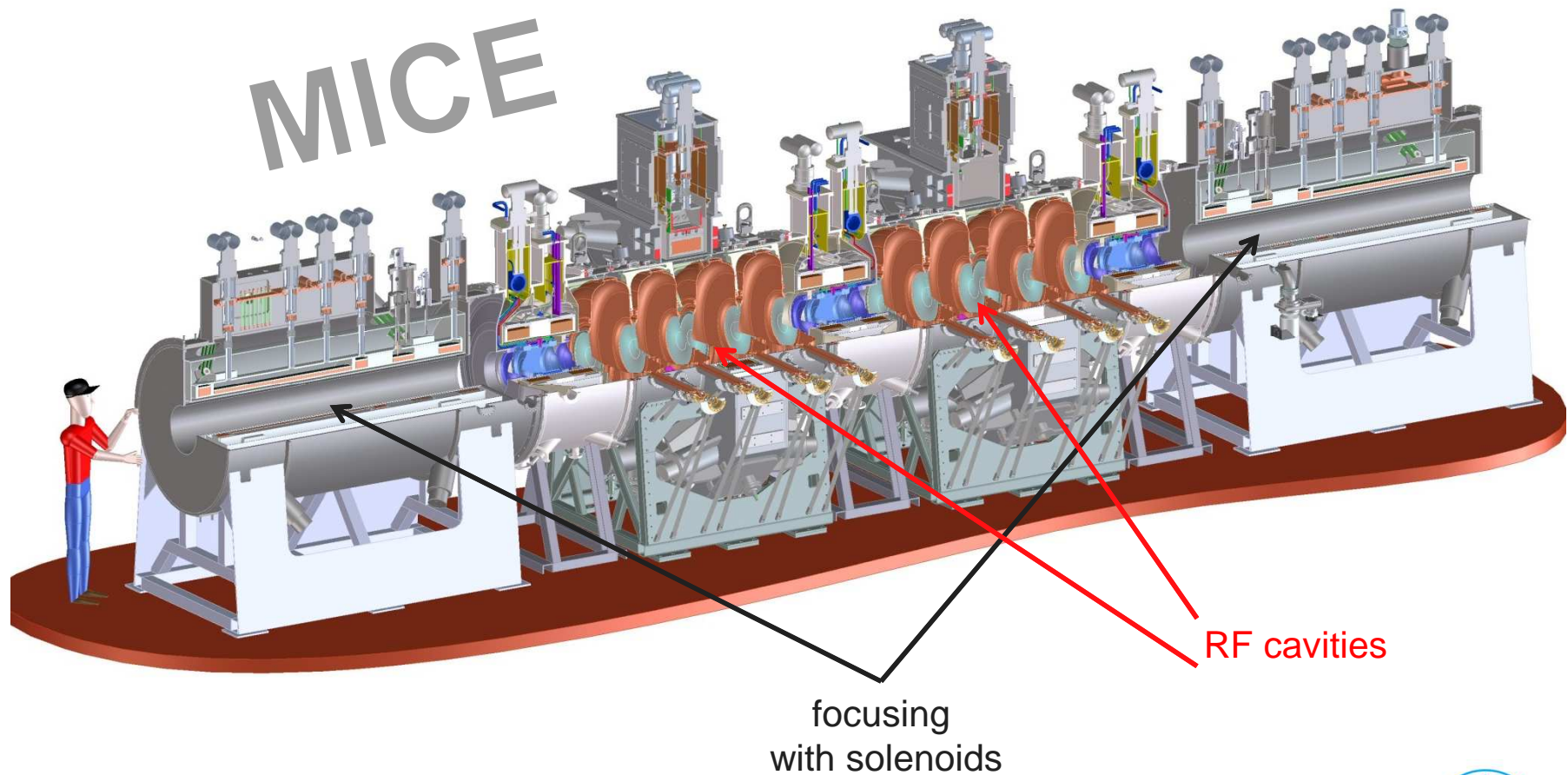
circumference = 4.5 km

$$L = 3 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$



MICE at FERMILAB (Chicago)

The Muon Ionization Cooling Experiment (MICE)
→ demonstrate the method and validate simulations



Summing-up

Basics of synchrotron radiation

particle type	limitation
• proton synchrotrons	dipole magnet
• electron synchrotrons	synchrotron radiation

International Linear Collider (ILC):

- luminosity eq.
- damping rings
- positron source
- power efficiency in superconducting cavities

Two very promising (and challenging) research areas in accel. physics:

- laser-driven plasma wakefield acceleration
- muon collider



Thank you for your attention

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