QCD Analysis of Polarized Structure Functions

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Outline:

- **Formalism**
- **World Data**
- **Parton Distributions**
- **New Parton Parametrizations with Errors**
- $\Lambda_{\rm QCD}$ **and** $\alpha_s(M_Z)$
- **Factorization Scheme Invariant Evolution**
- **Conclusions**

Transversity Workshop, Zeuthen, July 9th 2001

Polarized Structure Function $g_1(x)$

- Quark Parton Model (LO):

$$g_1(x) = \frac{1}{2}\sum_{i=1}^{n_f} e_i^2 \left[\Delta q_i(x) + \Delta\overline{q}_i(x)\right]$$

$$\Delta q_i(x) \quad = \quad (q_i^+(x) - q_i^-(x)) \ ,$$

$q_i^{+(-)}$ is the quark density with helicity aligned(anti-aligned) to the helicity of the parent nucleon (for $\Delta\overline{q}_i(x)$ accordingly)

- QCD Improved Quark Parton Model (NLO):

$$g_1(x,Q^2) \quad = \quad \frac{1}{2}\left(\frac{1}{n_f}\sum\nolimits_{i=1}^{n_f} e_i^2\right)\left[\delta C_S \otimes \Delta\Sigma + \delta C_G \otimes \Delta G + \delta C_{NS} \otimes \Delta q^{NS}\right]$$

The symbol $\otimes$ denotes convolution w.r.t. x of the polarized parton densities $\Delta q_i(x,Q^2)$ with the Wilson coeffient functions $\delta C_i(x,\alpha_s(Q^2))$:

$$f(x)\otimes g(x) = \int_0^1 dx_1 \int_0^1 dx_2 \delta(x - x_1x_2) f(x_1) g(x_2)$$

Polarized Structure Function $g_1(x)$

– Singlet Polarized Quark Distribution :

$$\Delta\Sigma(x,Q^2) = \sum_{i=1}^{n_f} \Delta^+ q_i(x,Q^2)$$

with

$$\Delta^+ q_i(x,Q^2) = \left[\Delta q_i(x,Q^2) + \Delta\overline{q}_i(x,Q(2)\right]$$

– Polarized Gluon Distribution : $\Delta G(x,Q^2)$

– Non-Singlet Polarized Quark Distribution :

$$\Delta q^{NS}(x,Q^2) = \frac{\sum_{i=1}^{n_f}\left(e_i^2 - \frac{1}{n_f}\sum_{k=1}^{n_f} e_k^2\right)\Delta^+ q_i(x,Q^2)}{\frac{1}{n_f}\sum_{k=1}^{n_f} e_k^2}$$

- For $n_f = 3\,(u,d,s,\bar{u},\bar{d},\bar{s})$:

$$\begin{aligned}
\Delta\Sigma &= (\Delta u + \Delta\bar{u}) + (\Delta d + \Delta\bar{d}) + (\Delta s + \Delta\bar{s}) \\
\Delta q^{NS}_{p(n)} &= +(-)\frac{3}{4}\Delta q_3 + \frac{1}{4}\Delta q_8
\end{aligned}$$

with

$$\begin{aligned}
\Delta q_3 &= (\Delta u + \Delta\bar{u}) - (\Delta d + \Delta\bar{d}) \\
\Delta q_8 &= (\Delta u + \Delta\bar{u}) + (\Delta d + \Delta\bar{d}) - 2(\Delta s + \Delta\bar{s})
\end{aligned}$$

Evolution Equations

- Evolution Equations to go from Q_0^2 to Q^2 ($t \equiv lnQ^2$)
  - for $\Delta\Sigma$ and ΔG :

$$\frac{d}{dt}\begin{pmatrix} \Delta\Sigma(x,t) \\ \Delta G(x,t) \end{pmatrix} = \frac{\alpha_s(t)}{2\pi}\begin{pmatrix} \Delta P_{qq}(x) & \Delta P_{qG}(x) \\ \Delta P_{Gq}(x) & \Delta P_{GG}(x) \end{pmatrix} \otimes \begin{pmatrix} \Delta\Sigma(x,t) \\ \Delta G(x,t) \end{pmatrix}$$

  - for Δq^{NS} :

$$\frac{d}{dt}\Delta q^{NS}(x,t) = \frac{\alpha_s(t)}{2\pi}\Delta P_{qq}^{NS}(x) \otimes \Delta q^{NS}(x,t)$$

with $\Delta P_{ij}(x)$ the polarized splitting functions.

- The polarized Wilson coefficient functions $\delta C_i(x, \alpha_s(Q^2))$ and the polarized splitting functions $\Delta P_{ij}(x, \alpha_s(Q^2))$ are known in the $\overline{MS}$ scheme up to $\mathcal{O}(\alpha_s^2)$. [E.B. Zijlstra and W.L. van Neerven, Nucl. Phys. B417 (1994) 61, R. Mertig and W.L. van Neerven, Z. Phys. C70 (1996) 637, W. Vogelsang, Phys. Rev. D54 (1996) 2023]

$\Longrightarrow$ NLO QCD Analysis: Determination of $\Lambda_{\rm QCD}$ and the Parton densities with Errors.

Parametrization

- General choice for the parametrization of the polarized parton distributions at Q_0^2:

$$x\Delta q_i(x, Q_0^2) = \eta_i A_i x^{a_i}(1-x)^{b_i}(1+\gamma_i x + \rho_i x^{\frac{1}{2}})$$

- Normalization:

$$\begin{aligned} A_i^{-1} &= \left(1+\gamma_i \frac{a_i}{a_i+b_i+1}\right)\frac{\Gamma(a_i)\Gamma(b_i+1)}{\Gamma(a_i+b_i+1)} \\ &\quad +\rho_i \frac{\Gamma(a_i+0.5)\Gamma(b_i+1)}{\Gamma(a_i+b_i+1.5)} \end{aligned}$$

such that

$$\int_0^1 dx \Delta q_i(x, Q_0^2) = \eta_i$$

are the first moment of $\Delta q_i(x, Q_0^2)$

- The polarized parton distributions to be fitted are:

$$\Delta u_v \ , \ \Delta d_v \ , \ \Delta \bar{q} \ , \ \Delta G$$

where the index v denotes the $valence$ quark.

Note that: $\Delta q + \Delta \bar{q} = \Delta q_v + 2\Delta \bar{q}$.

Choice of Parameters

- $Q_0^2 = 4.0\,GeV^2$
- SU(3) flavour symmetry assumed

  η_{u_v} and η_{d_v} determined from F and D the SU(3) parameters involved in the matrix elements describing the neutron and hyperon β-decays:

  $\eta_{u_v} = 2F = 0.926$; $\eta_{d_v} = F - D = -0.341$
- Flavor symmetric sea assumed

  $\Delta\bar{u}(x, Q_0^2) = \Delta\bar{d}(x, Q_0^2) = \Delta\bar{s}(x, Q_0^2) = \Delta\bar{q}(x, Q_0^2)$
- No assumption made concerning positivity and helicity retention
- For u_v and d_v : $\rho_{u_v} = \rho_{d_v} = 0$
- For the gluon : $\gamma_G = \rho_G = 0$ (Gluon A)
- For the sea : $\gamma_{\bar{q}} = \rho_{\bar{q}} = 0$ (Sea A)
- The normalizations of the different data sets against each other were fitted and fixed afterwards
- The remaining 12 parameters to be determined are:

  Δu_v: a_u , b_u , γ_u Δd_v: a_d , b_d , γ_d

  $\Delta\bar{q}$: $\eta_{\bar{q}}$, $a_{\bar{q}}$, $b_{\bar{q}}$ ΔG: η_G , a_G , b_G

Note:

$$x\Delta q_i(x, Q_0^2) = \eta_i A_i x^{a_i}(1-x)^{b_i}(1 + \gamma_i x + \rho_i x^{\frac{1}{2}})$$

The World Data: $g_1/F_1(x,Q^2)$ or $A_1(x,Q^2)$

Published Experimental Data above $Q^2 = 1.0\ GeV^2$

Experiment	type	x–range	Q^2–range $[GeV^2]$	ΔN [%]	# data points
E143	g_1/F_1	0.027 – 0.749	1.17 – 9.52	3.7	82
HERMES	g_1/F_1	0.028 – 0.660	1.13 – 7.46	3.0	39
E155	g_1/F_1	0.015 – 0.750	1.22 – 34.72	7.6	24
SMC	A_1	0.005 – 0.480	1.30 – 58.0	4.0	59
EMC	A_1	0.015 – 0.466	3.50 – 29.5	14.0	10
proton					214
E143	g_1/F_1	0.027 – 0.749	1.17 – 9.52	4.9	82
E155	g_1/F_1	0.015 – 0.750	1.22 – 34.79	4.0	24
SMC	A_1	0.005 – 0.479	1.30 – 54.8	4.0	65
deuteron					171
E142	A_1	0.035 – 0.466	1.10 – 5.50	3.0	28
HERMES	A_1	0.033 – 0.464	1.22 – 5.25	3.0	9
E154	A_1	0.017 – 0.564	1.20 – 15.0	3.0	11
neutron					48
total					433

$$g_1/F_1 \approx \frac{1}{(1+\gamma^2)} A_1 \ , \quad where \quad \gamma^2 = Q^2/\nu^2$$

$$F_1 = \frac{(1+\gamma^2)}{2x(1+R)} F_2$$

F_2–Parametrization: NMC, M. Arneodo et al., Phys. Lett. **B364** (1995) 107.

R–Parametrization: SLAC, L. Withlow et al., Phys. Lett. **B250** (1990) 193.

The World Data: $g_1(x, Q^2)$

Published Experimental Data above $Q^2 = 1.0\ GeV^2$

Experiment	x-range	Q^2-range $[GeV^2]$	ΔN [%]	# data points
E143	0.031 – 0.749	1.27 – 9.52	3.7	28
HERMES	0.028 – 0.660	1.13 – 7.46	3.0	39
E155	0.015 – 0.750	1.22 – 34.73	7.6	24
SMC	0.005 – 0.480	1.30 – 58.0	4.0	12
proton				103
E143	0.031 – 0.749	1.27 – 9.52	4.9	28
E155	0.015 – 0.750	1.22 – 34.79	4.0	24
SMC	0.005 – 0.479	1.30 – 54.8	4.0	12
deuteron				64
E142	0.035 – 0.466	1.10 – 5.50	3.0	8
HERMES	0.033 – 0.464	1.22 – 5.25	3.0	9
E154	0.017 – 0.564	1.20 – 15.0	3.0	17
neutron				34
total				201

Parameter Values at $Q_0^2 = 4.0\,GeV^2$

8+1 Parameter Fit based on the Asymmetry Data:

Parameter	LO		NLO	
	value	error	value	error
η_{u_v}	0.926	fixed	0.926	fixed
a_{u_v}	0.209	0.034	0.331	0.078
b_{u_v}	2.417	0.205	3.257	0.344
γ_{u_v}	21.34	fixed (*)	27.22	fixed (*)
η_{d_v}	-0.341	fixed	-0.341	fixed
a_{d_v}	0.125	0.077	0.175	0.119
b_{d_v}	3.091	1.441	3.852	1.835
γ_{d_v}	38.50	fixed (*)	22.95	fixed (*)
$\eta_{\bar{q}}$	-0.431	0.142	-0.461	0.133
$a_{\bar{q}}$	0.338	0.265	0.493	0.364
$b_{\bar{q}}$	4.98	fixed (*)	5.82	fixed (*)
η_G	1.182	0.704	0.887	0.642
a_G	2.219	1.024	1.674	0.887
b_G	4.65	fixed (*)	4.91	fixed (*)
χ^2 / NDF	0.99		0.92	

$\Rightarrow$ The parameters marked by (*) have been fitted first and then fixed since the present data do not constrain its values well enough.

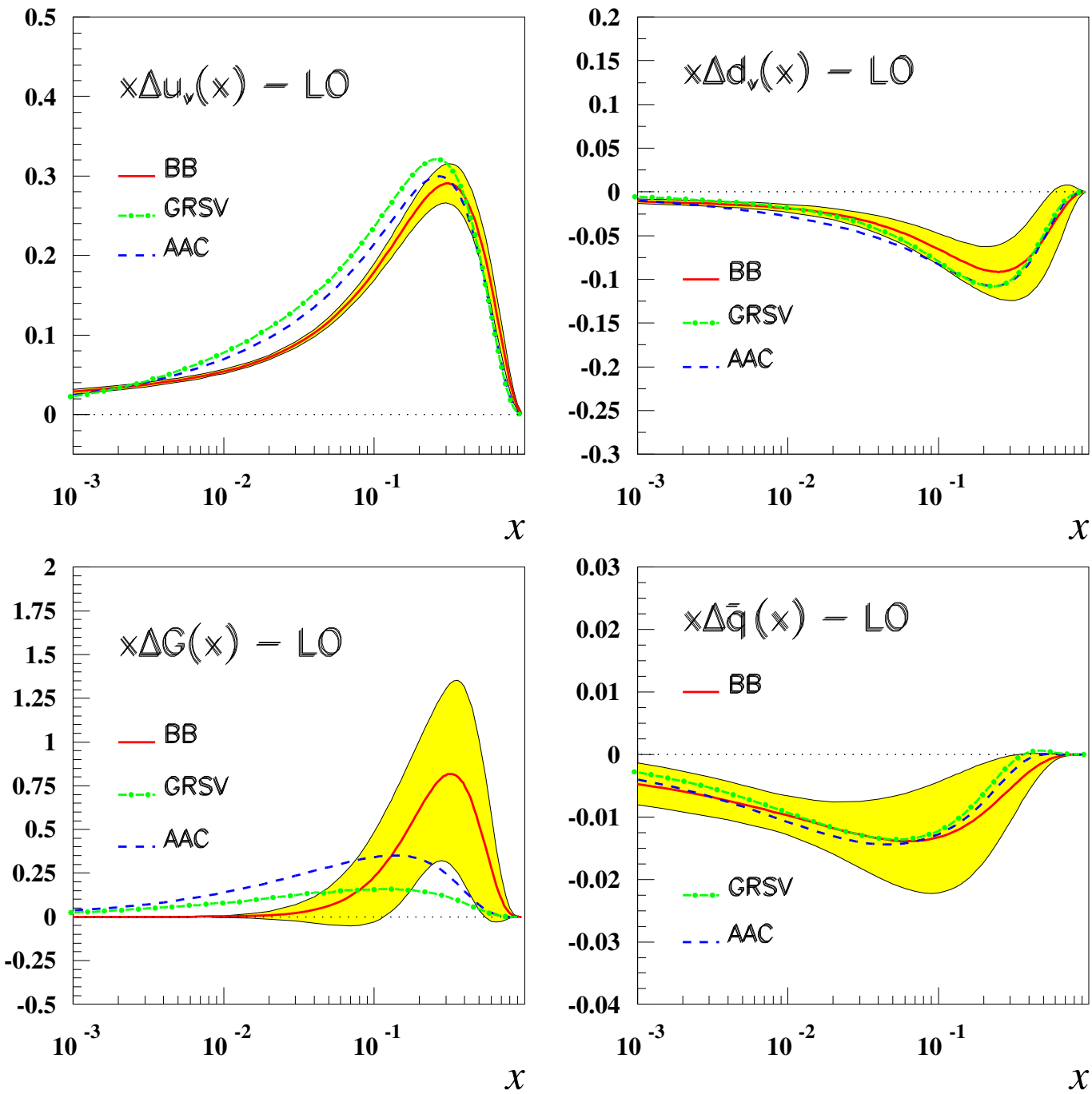
Covariance Matrices – 8 + 1 Parameter Fit

LO	$\Lambda^{(4)}_{QCD}, GeV$	a_{u_v}	b_{u_v}	a_{d_v}	b_{d_v}	$\eta_{\bar{q}}$	$a_{\bar{q}}$	η_G	a_G
$\Lambda^{(4)}_{QCD}$	4.72e-3								
a_{u_v}	2.94e-5	1.19e-3							
b_{u_v}	-1.10e-3	5.49e-3	4.21e-2						
a_{d_v}	-6.10e-4	-1.48e-3	-7.33e-3	5.87e-3					
b_{d_v}	-9.74e-3	-2.37e-2	-3.36e-2	7.09e-2	2.08e-0				
$\eta_{\bar{q}}$	3.75e-5	2.88e-3	1.25e-2	-4.74e-3	-3.30e-2	2.01e-2			
$a_{\bar{q}}$	5.42e-4	7.81e-3	3.75e-2	-1.51e-2	-1.82e-1	3.13e-2	7.00e-2		
η_G	2.98e-3	-8.58e-3	-3.99e-2	1.38e-2	-4.47e-3	-6.51e-2	-8.82e-2	4.96e-1	
a_G	-2.57e-3	6.08e-3	4.03e-2	-2.29e-2	-3.93e-1	4.13e-2	8.29e-2	-3.50e-1	1.05e-0

NLO	$\Lambda^{(4)}_{QCD}, GeV$	a_{u_v}	b_{u_v}	a_{d_v}	b_{d_v}	$\eta_{\bar{q}}$	$a_{\bar{q}}$	η_G	a_G
$\Lambda^{(4)}_{QCD}$	3.41e-3								
a_{u_v}	4.17e-5	6.06e-3							
b_{u_v}	-2.24e-4	2.38e-2	1.18e-1						
a_{d_v}	-3.03e-4	-6.12e-3	-2.44e-2	1.43e-2					
b_{d_v}	1.31e-3	-3.67e-2	-1.69e-2	1.03e-1	3.37e-0				
$\eta_{\bar{q}}$	-1.65e-4	-5.32e-4	-5.07e-3	3.29e-3	6.13e-2	1.77e-2			
$a_{\bar{q}}$	-3.58e-4	2.29e-2	8.96e-2	-2.80e-2	-5.59e-2	2.05e-2	1.33e-1		
η_G	5.88e-4	1.04e-2	4.96e-2	-1.39e-2	-1.28e-1	-7.20e-2	-6.17e-2	4.12e-1	
a_G	-2.37e-4	4.25e-2	1.53e-1	-5.01e-2	3.32e-2	4.09e-2	2.72e-1	-1.73e-1	7.87e-1

Pol. Parton Densities at $Q_0^2 = 4.0\,GeV^2$

- 8+1 Parameter Fit based on the Asymmetry Data:



$\Rightarrow$ Yellow error band: Fully correlated 1σ Gaussian error propagation at the input scale Q_0^2.

Pol. Parton Densities at $Q_0^2 = 4.0\,GeV^2$

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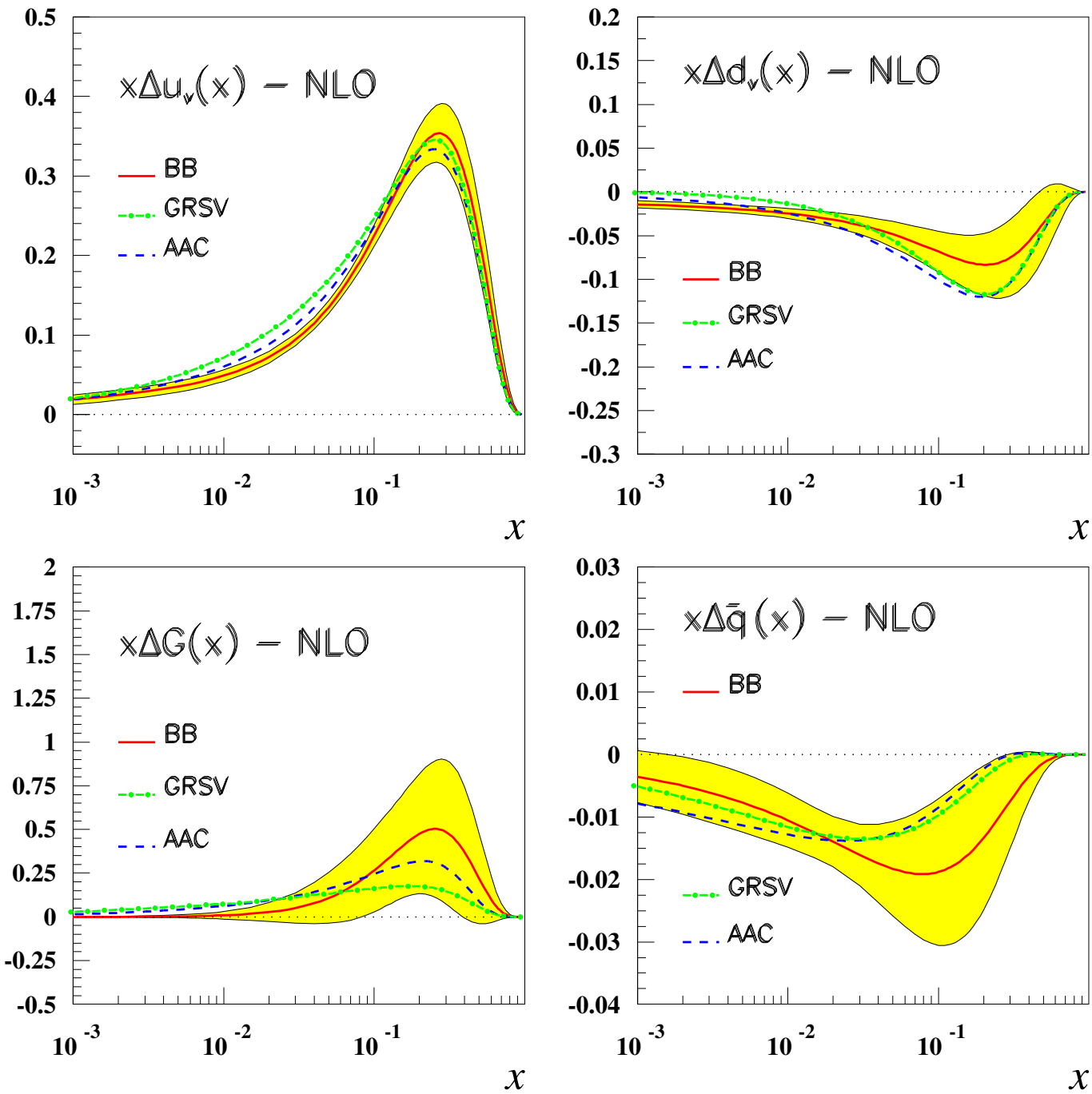

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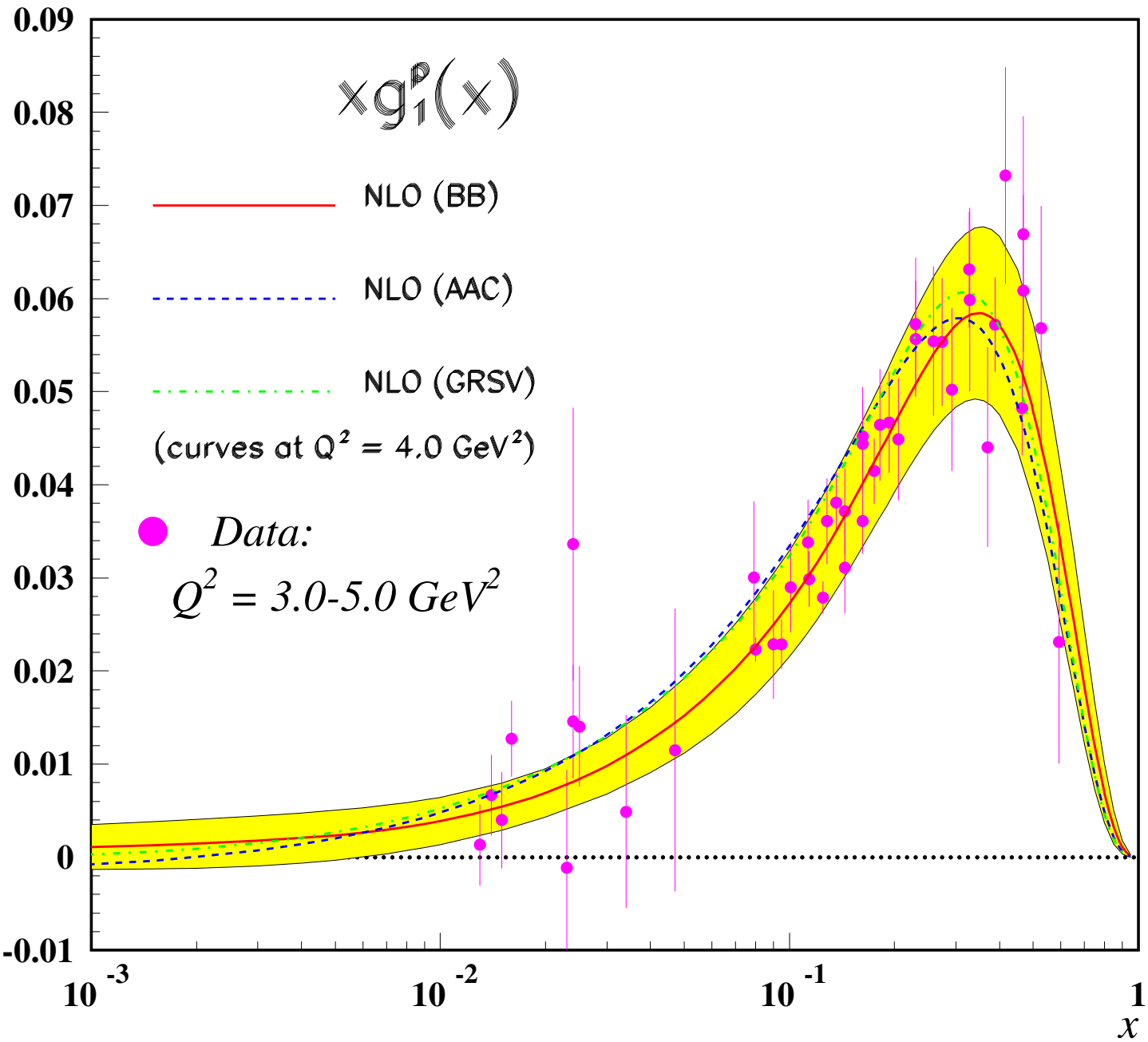
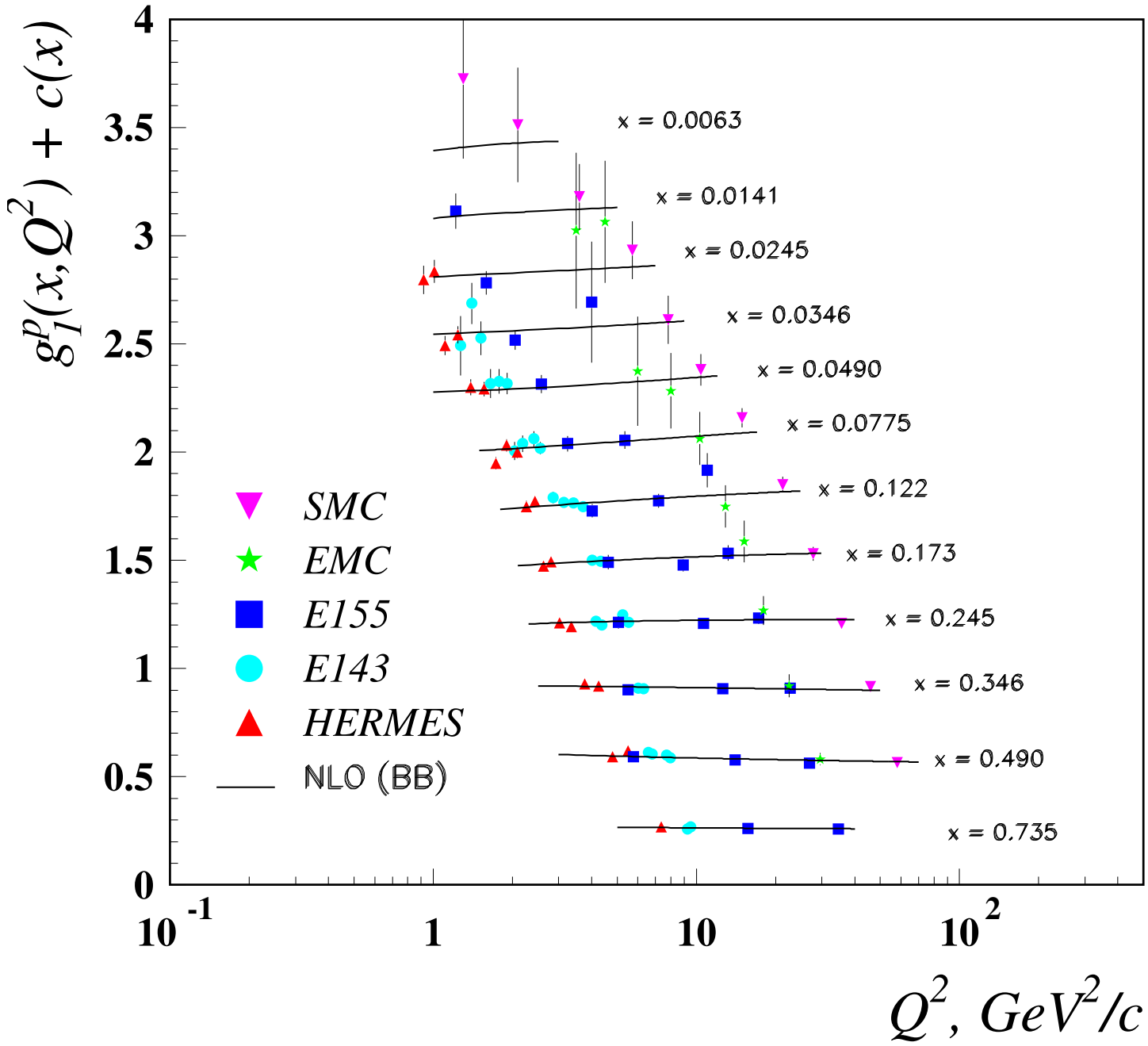
Comparison with measured $xg_1^p(x)$



$\Rightarrow$ Yellow error band: Fully correlated 1σ Gaussian error propagation at $Q^2 = 4.0\ GeV^2$.

$g_1^p(x)$ versus Q^2



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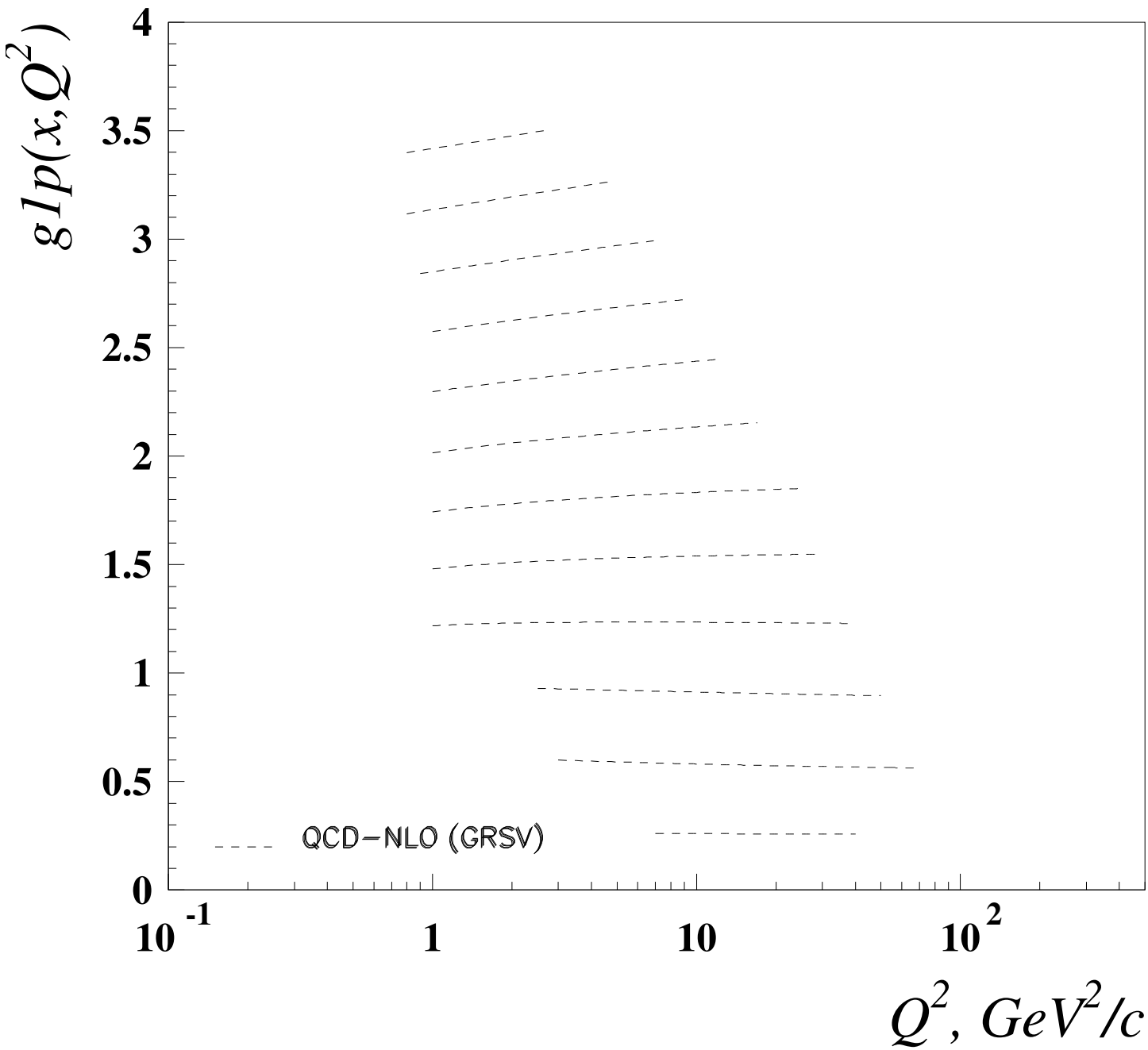
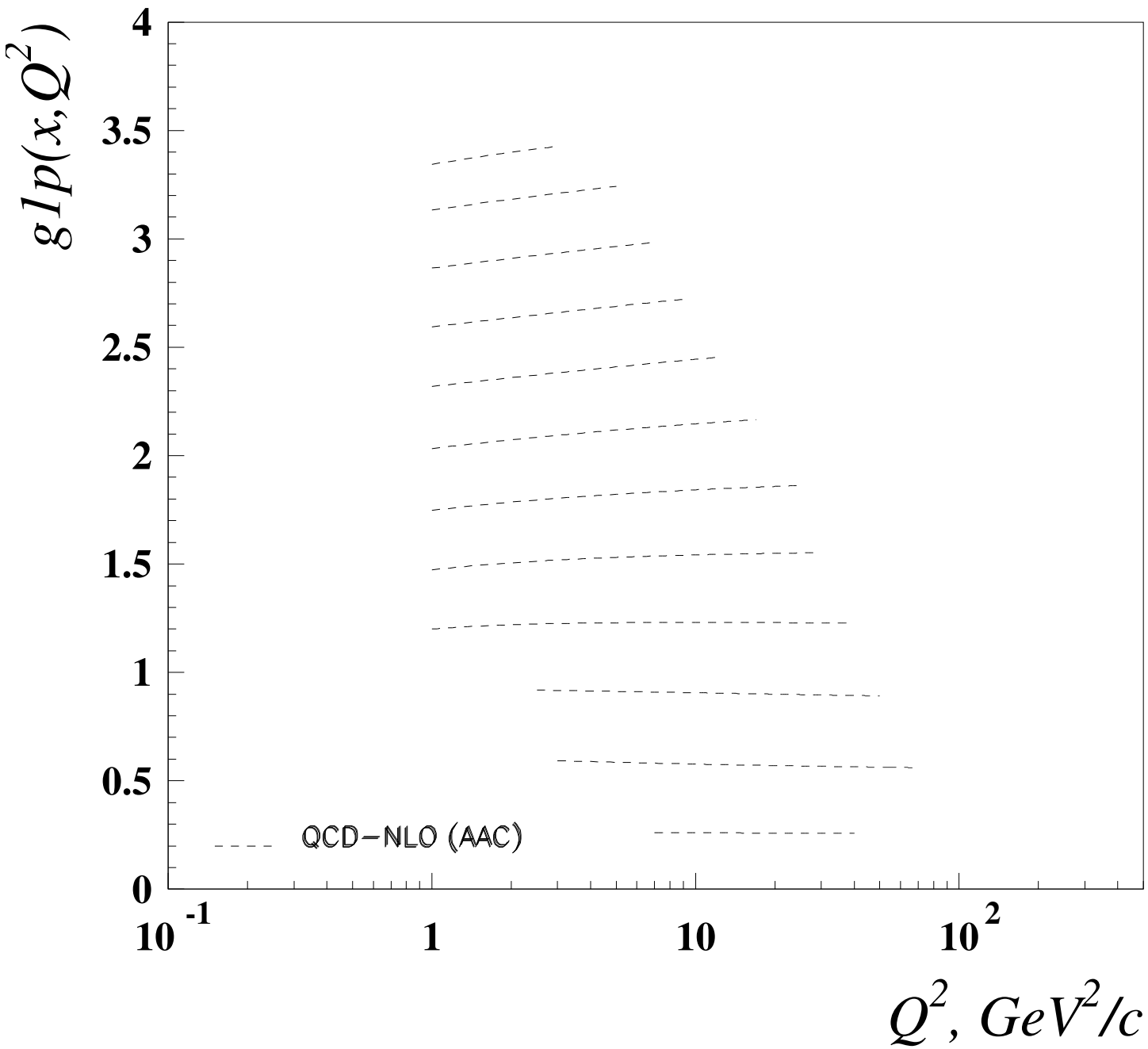
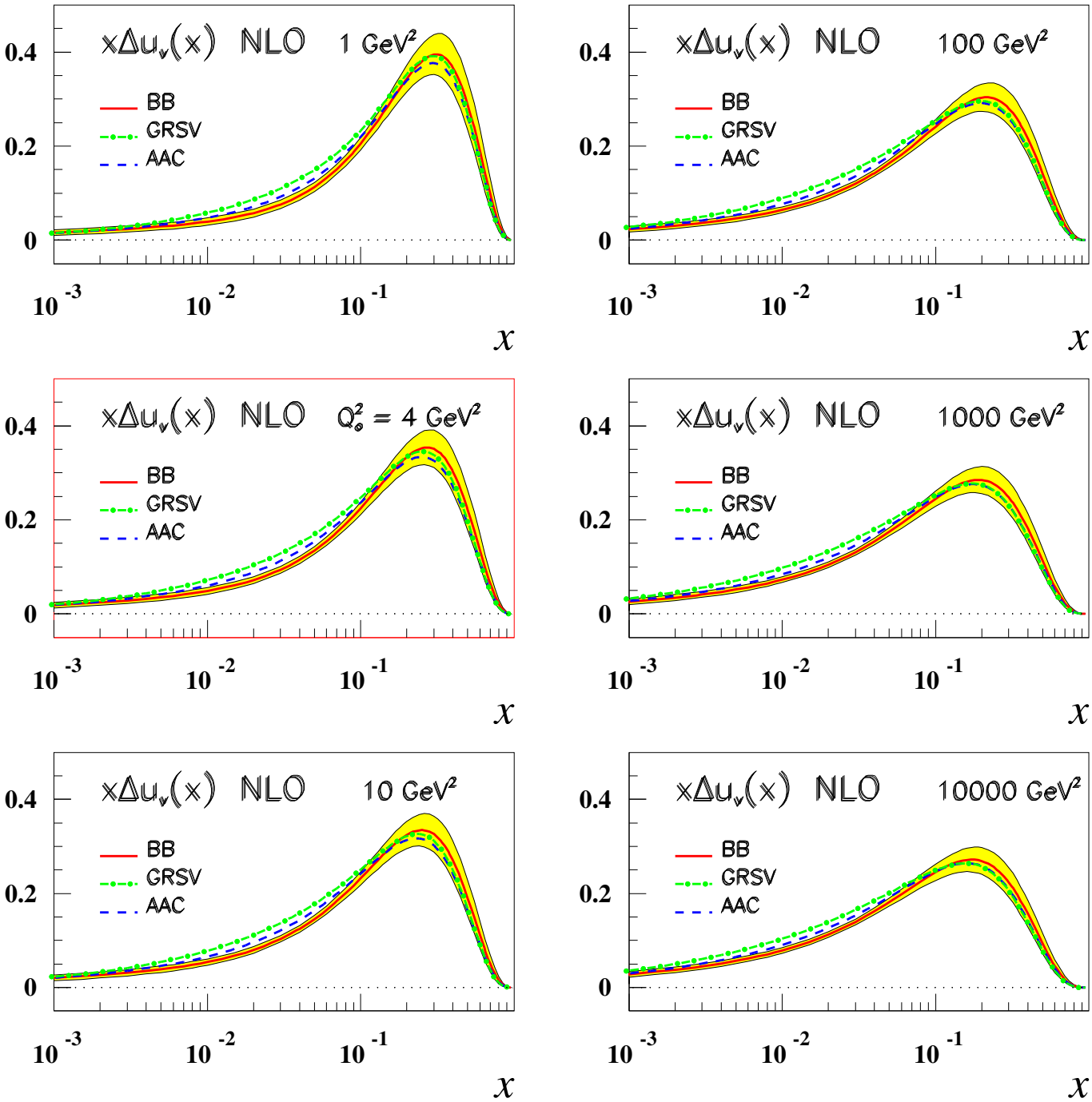


$g_1^p(x)$ versus Q^2



Evolution of Polarized Parton Densities

- 8+1 Parameter Fit based on the Asymmetry Data:



$\Rightarrow$ Yellow error band: Fully correlated 1σ Gaussian error propagation through the evolution equation.

Evolution of Polarized Parton Densities

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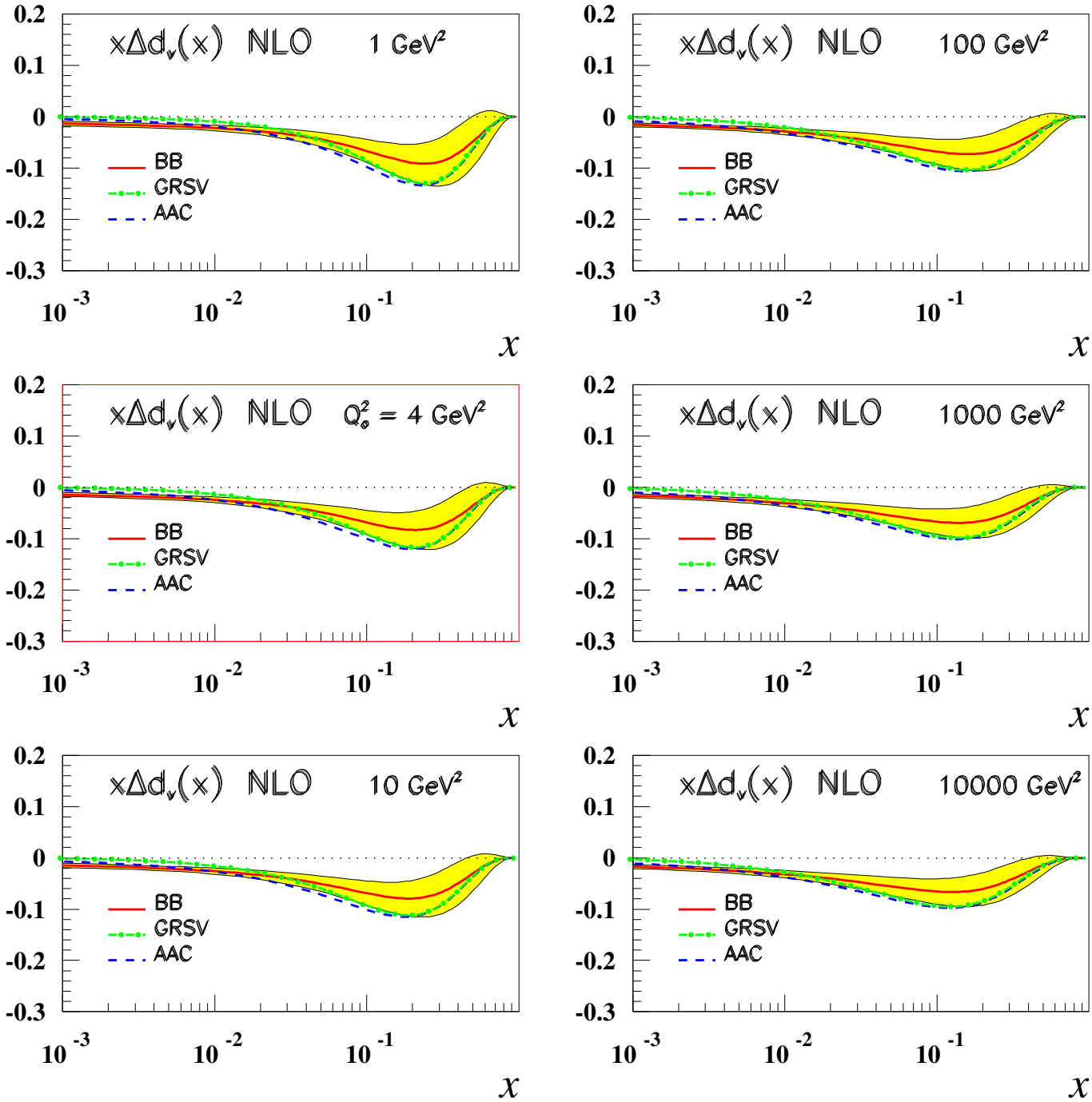

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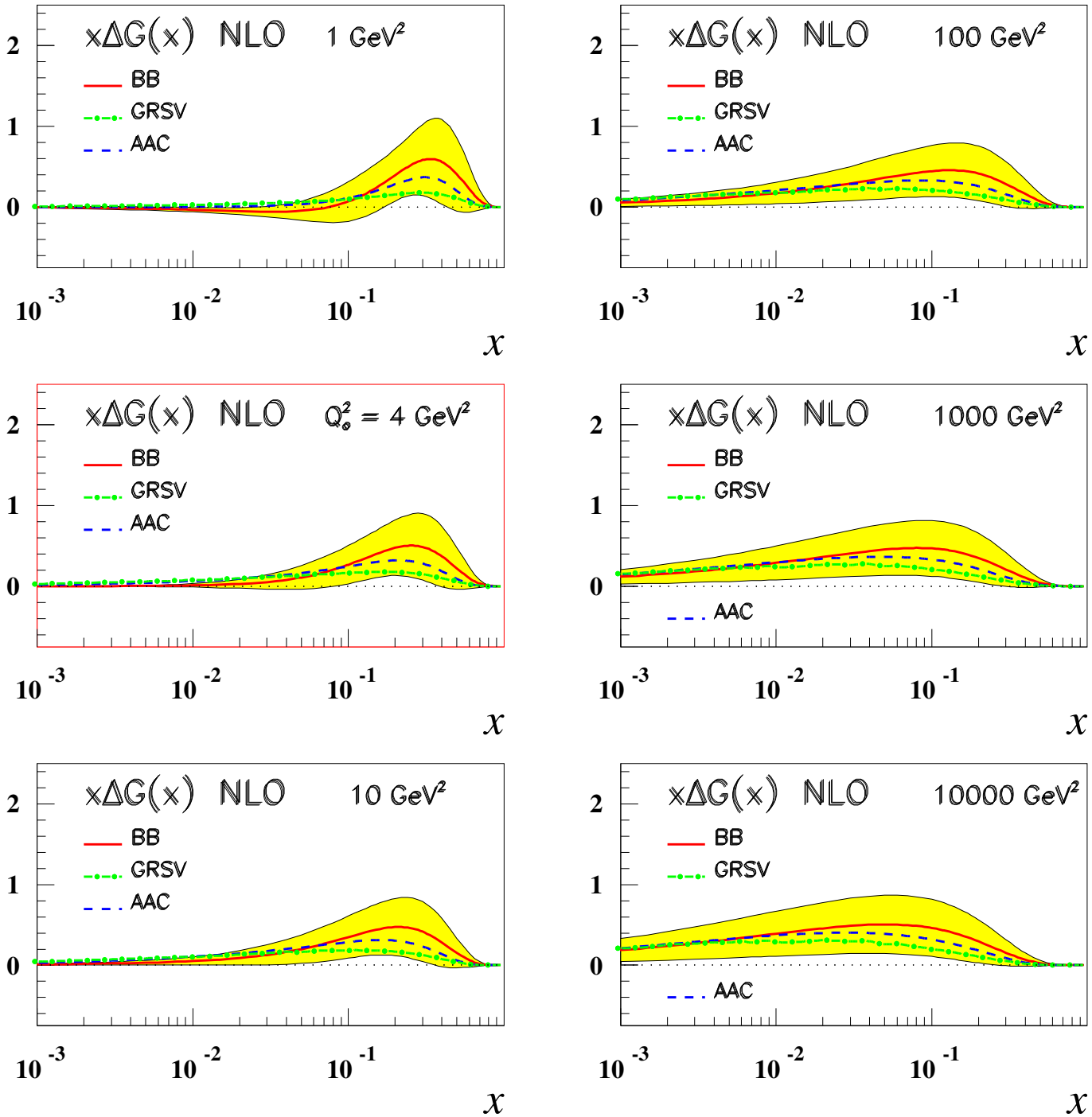

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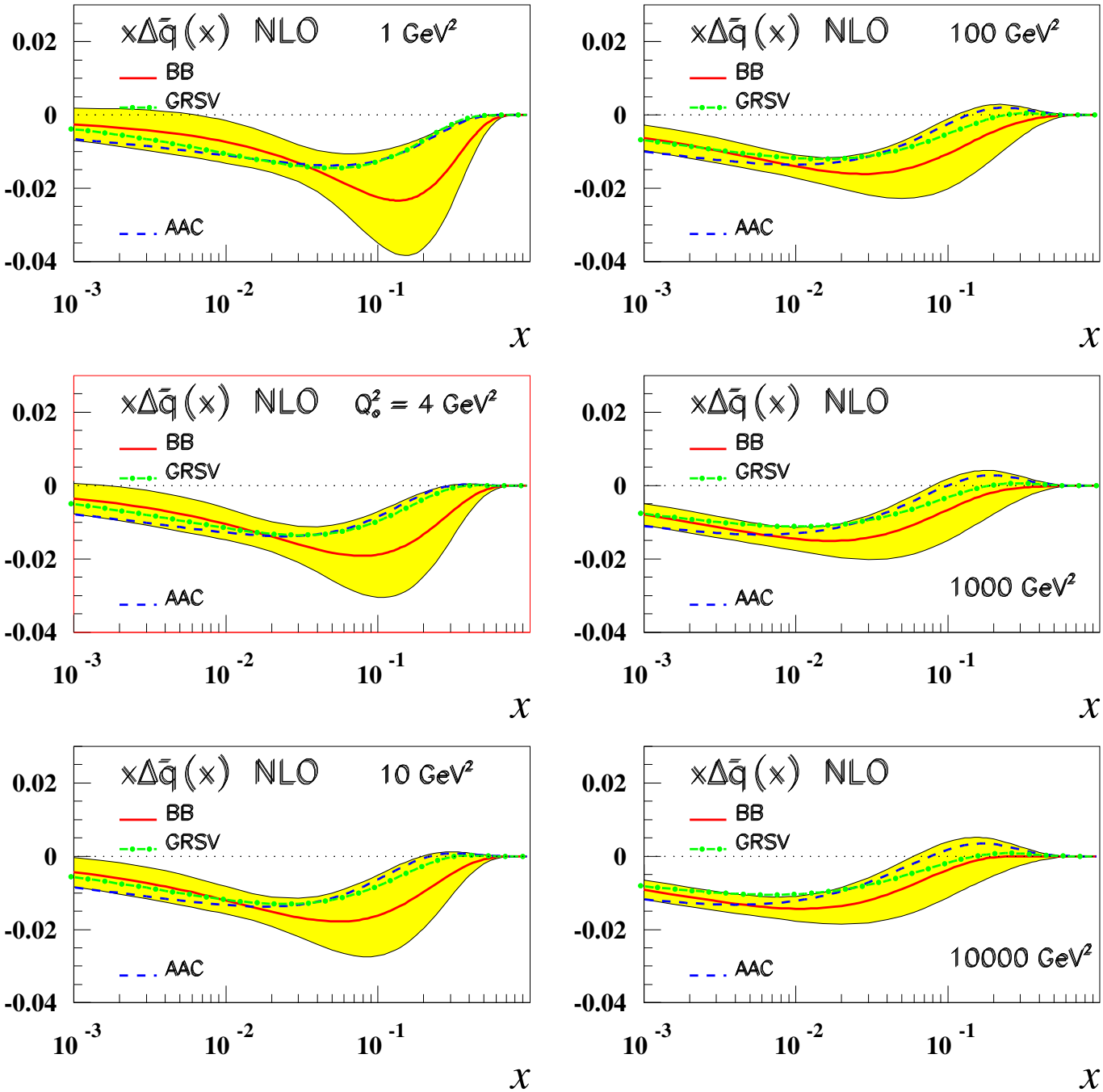

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8+1 parameter NLO fit: $\Lambda^{(4)}_{QCD} \Rightarrow \alpha_s(M_Z^2)$

$\Lambda^{(4)}_{QCD}$ [Gev]	$A1$		$g1$	
	value	error	value	error
FS/RS=1.0/1.0	0.241	$\pm$ 0.058	0.241	$\pm$ 0.062
FS/RS=0.5/1.0	0.197	– 0.044	0.193	– 0.048
FS/RS=2.0/1.0	0.306	+ 0.065	0.306	+ 0.065
FS/RS=1.0/0.5	0.358	+ 0.117	0.348	+ 0.107
FS/RS=1.0/2.0	0.183	– 0.058	0.183	– 0.058

- $A1$:

$$\alpha_s(M_Z^2) = 0.114 \begin{matrix} +0.004 \\ -0.005 \end{matrix} \quad \begin{matrix} +0.005 \\ -0.004 \end{matrix} \quad \begin{matrix} +0.008 \\ -0.005 \end{matrix}$$

(fit) (fac) (ren)

- $g1$:

$$\alpha_s(M_Z^2) = 0.114 \begin{matrix} +0.004 \\ -0.005 \end{matrix} \quad \begin{matrix} +0.005 \\ -0.004 \end{matrix} \quad \begin{matrix} +0.007 \\ -0.005 \end{matrix}$$

- SMC: $0.121 \pm 0.002(stat) \pm 0.006(syst + theor)$

E154: $0.108 - 0.116 \quad (bad\, for \geq 0.120)$

ABFR: $0.120 \begin{matrix} +0.004 \\ -0.005 \end{matrix} (exp) \begin{matrix} +0.009 \\ -0.006 \end{matrix} (theor)$

world average (PDG): 0.118 ± 0.002

Fac. Scheme Invariant Combinations

- Instead of PROCESS-INDEPENDENT SCHEME-DEPENDENT Evolution Equations for PARTONS one may think of PROCESS-DEPENDENT SCHEME-INDEPENDENT Evolution Equations for OBSERVABLES, F_A, F_B.

  $\Rightarrow$ The input densities are measured ! Control over the input directly.
  $\Rightarrow$ No ΔG–Ansatz necessary.
  $\Rightarrow$ A one parameter fit only – Λ_{QCD}.

Evolution Equations : [J. Blümlein, V. Ravindran, and W. L. van Neerven, Nucl. Phys **B586** (2000) 349.]

$$\frac{\partial}{\partial t}\begin{pmatrix} F_A^N \\ F_B^N \end{pmatrix} = -\frac{1}{4}\begin{pmatrix} K_{AA}^N & K_{AB}^N \\ K_{BA}^N & K_{BB}^N \end{pmatrix}\begin{pmatrix} F_A^N \\ F_B^N \end{pmatrix}$$

evolution variable :

$$t = -\frac{2}{\beta_0}\log\left(\frac{a_s(Q^2)}{a_s(Q_0^2)}\right)$$

$\Longrightarrow$ The evolution kernels K_{IJ}^N are also Physical Quantities! The Factorization Scheme Independence holds order by order.

The Renormalization Scale Dependence disappears only with more higher orders.

$\Longrightarrow$ A possible choice: $F_A = g_1$ and $F_B = \partial g_1/\partial t$.

System : $g_1(x,Q^2), \partial g_1/\partial t(x,Q^2)$

Leading Order :

$$K_{22}^{N(0)} = 0$$

$$K_{2d}^{N(0)} = -4$$

$$K_{d2}^{N(0)} = \frac{1}{4}\left(\gamma_{qq}^{N(0)}\gamma_{gg}^{N(0)} - \gamma_{qg}^{N(0)}\gamma_{gq}^{N(0)}\right)$$

$$K_{dd}^{N(0)} = \gamma_{qq}^{N(0)} + \gamma_{gg}^{N(0)}$$

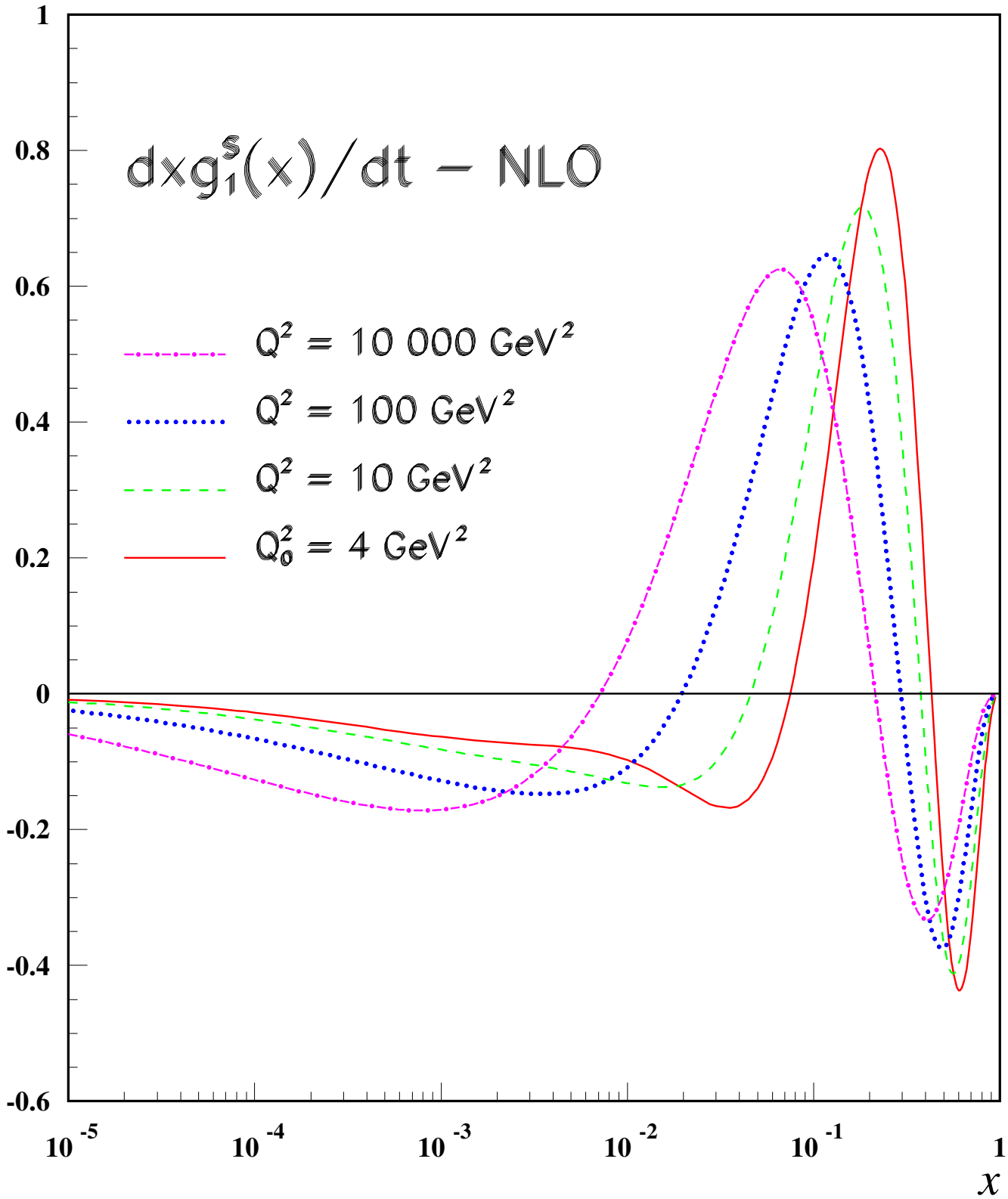
Next-to-Leading Order : [W. Furmanski and R. Petronzio, Z. Phys. **C 11** (1982) 293.]

$$K_{22}^{N(1)} = K_{2d}^{N(1)} = 0$$

$$\begin{aligned}
K_{d2}^{N(1)} = & \frac{1}{4}\left[\gamma_{gg}^{N(0)}\gamma_{qq}^{N(1)} + \gamma_{gg}^{N(1)}\gamma_{qq}^{N(0)} - \gamma_{qg}^{N(1)}\gamma_{gq}^{N(0)} - \gamma_{qg}^{N(0)}\gamma_{gq}^{N(1)}\right] \\
& -\frac{\beta_1}{2\beta_0}\left(\gamma_{qq}^{N(0)}\gamma_{gg}^{N(0)} - \gamma_{qg}^{N(0)}\gamma_{gq}^{N(0)}\right) \\
& +\frac{\beta_0}{2}C_{2,q}^{N(1)}\left(\gamma_{qq}^{N(0)} + \gamma_{gg}^{N(0)} - 2\beta_0\right) \\
& -\frac{\beta_0}{2}\frac{C_{2,g}^{N(1)}}{\gamma_{qg}^{N(0)}}\left[\gamma_{qq}^{N(0)\,2} - \gamma_{qq}^{N(0)}\gamma_{gg}^{N(0)} + 2\gamma_{qg}^{N(0)}\gamma_{gq}^{N(0)} - 2\beta_0\gamma_{qq}^{N(0)}\right] \\
& -\frac{\beta_0}{2}\left(\gamma_{qq}^{N(1)} - \frac{\gamma_{qq}^{N(0)}\gamma_{qg}^{N(1)}}{\gamma_{qg}^{N(0)}}\right)
\end{aligned}$$

$$\begin{aligned}
K_{dd}^{N(1)} = & \gamma_{qq}^{N(1)} + \gamma_{gg}^{N(1)} - \frac{\beta_1}{\beta_0}\left(\gamma_{qq}^{N(0)} + \gamma_{gg}^{N(0)}\right) + 4\beta_0 C_{2,q}^{N(1)} - 2\beta_1 \\
& -\frac{2\beta_0}{\gamma_{qg}^{N(0)}}\left[C_{2,g}^{N(1)}\left(\gamma_{qq}^{N(0)} - \gamma_{gg}^{N(0)} - 2\beta_0\right) - \gamma_{qg}^{N(1)}\right]
\end{aligned}$$

$\partial x g_1^S / \partial t(x, Q^2)$ and shift of $\Lambda_{QCD}^{(4)}$



$\Lambda_{QCD}^{(4)}$: 0.241 $\Longrightarrow$ 0.229 , $\alpha_s(M_Z^2)$: 0.114 $\Longrightarrow$ 0.113

Conclusions

- An LO and NLO QCD Analysis of the current World-Data of Polarized Structure Functions was performed.

- New parametrizations of the parton densities including their Errors were derived. They are available via a fast FORTRAN program for the range:

$$1 < Q^2 < 10^6 \text{ GeV}^2 \text{ and } 10^{-4} < x < 1.$$

- The following results for $\Lambda_{\rm QCD}$ and $\alpha_s(M_Z^2)$ were obtained (8+1 parameter fit):

$$\Lambda^{(4)}_{\rm QCD} = 241 \pm 58 \text{ (fit)} \; {}^{+65}_{-44} \text{ (fac)} \; {}^{+117}_{-58} \text{ (ren)}$$

$$\alpha_s(M_Z^2) = 0.114 \; {}^{+0.004}_{-0.005} \text{ (fit)} \; {}^{+0.005}_{-0.004} \text{ (fac)} \; {}^{+0.008}_{-0.005} \text{ (ren)}$$

- First steps in a fac. scheme invariant QCD evolution based on the structure function $g_1(x, Q^2)$ and $\partial g_1(x, Q^2)/\partial \log Q^2$ were performed yielding similar results for $\alpha_s(M_Z^2)$.

- The latter analysis is a very promising way to proceed in the future, since it allows to extract $\Lambda_{\rm QCD}$ fixing all the input distributions by direct measurement.