FRAGMENTATION INTO SPIN-1 HADRONS

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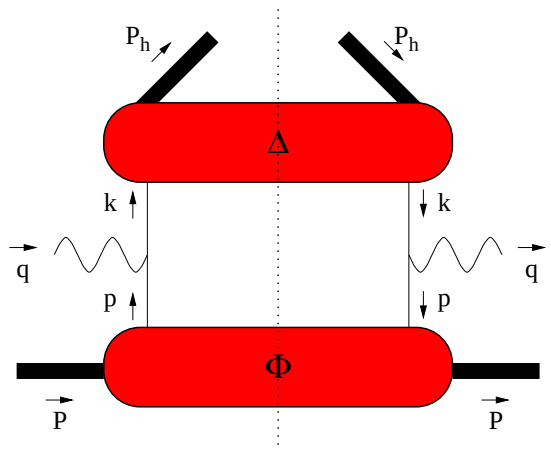
FOM

- Semi-inclusive DIS
- Fragmentation to spin-1 hadrons
- Transversity measurements
- Positivity bounds and modeling attempts

Semi-inclusive DIS

$$\mathrm{d}\sigma(l + H \to l' + h + X) \propto L_{\mu\nu}\, W^{\mu\nu}$$



$$2MW^{\mu\nu} \propto \mathrm{Tr}\left[\Phi(x_{_B})\ \gamma^{\mu}\ \Delta(z_h)\ \gamma^{\nu}\right]$$

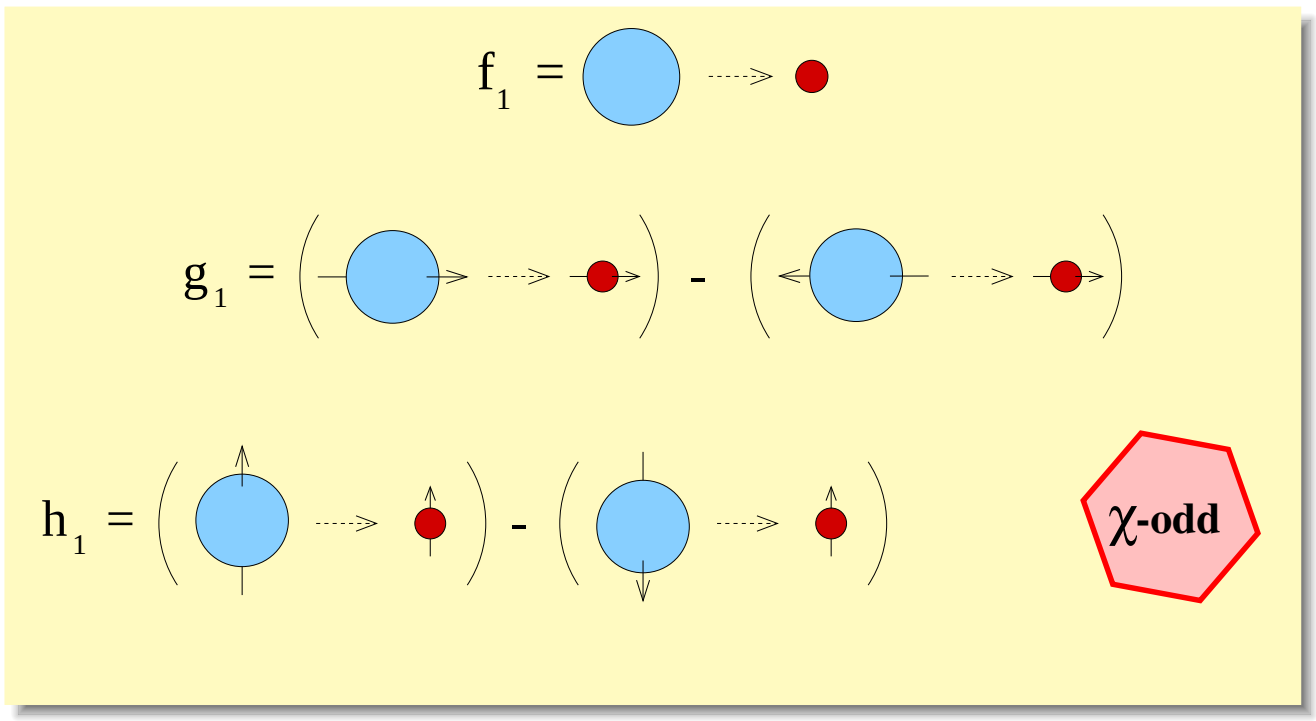
$$x_{_B} = \frac{Q^2}{2\,P\cdot q} \qquad z_h = \frac{2\,P_h\cdot q}{Q^2}$$

$$\Phi(x) = \frac{1}{2}\int \mathrm{d}^2\boldsymbol{p}_T\ \mathrm{d}p^-\ \Phi(p,P)\Big|_{p^+=xP^+}$$

$$\Delta(z) = \frac{z}{4}\int \mathrm{d}^2\boldsymbol{k}_T\ \mathrm{d}k^+\ \Delta(k,P_h)\Big|_{k^-=\frac{P_h^-}{z}}$$

Distribution functions for spin-1/2

From decomposition of Φ correlation function.

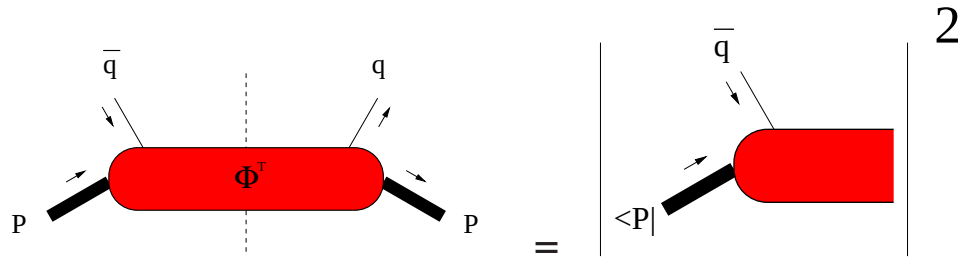


Quark+target spin structure

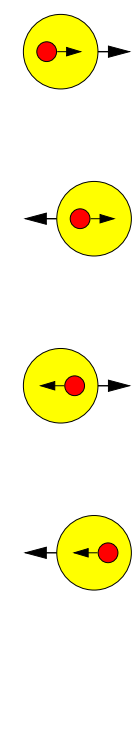
Full scattering matrix $M = (\Phi\gamma^+)^T$

for $\overline{q} + P \to X$

in parton $\otimes$ hadron helicity spaces



$$\begin{pmatrix} f_1+g_1 & 0 & 0 & 2\,h_1 \\ 0 & f_1-g_1 & 0 & 0 \\ 0 & 0 & f_1-g_1 & 0 \\ 2\,h_1 & 0 & 0 & f_1+g_1 \end{pmatrix}$$



Soffer bound

Positive definiteness of this matrix requires:

Positivity of diagonal elements $\Rightarrow$ trivial bounds

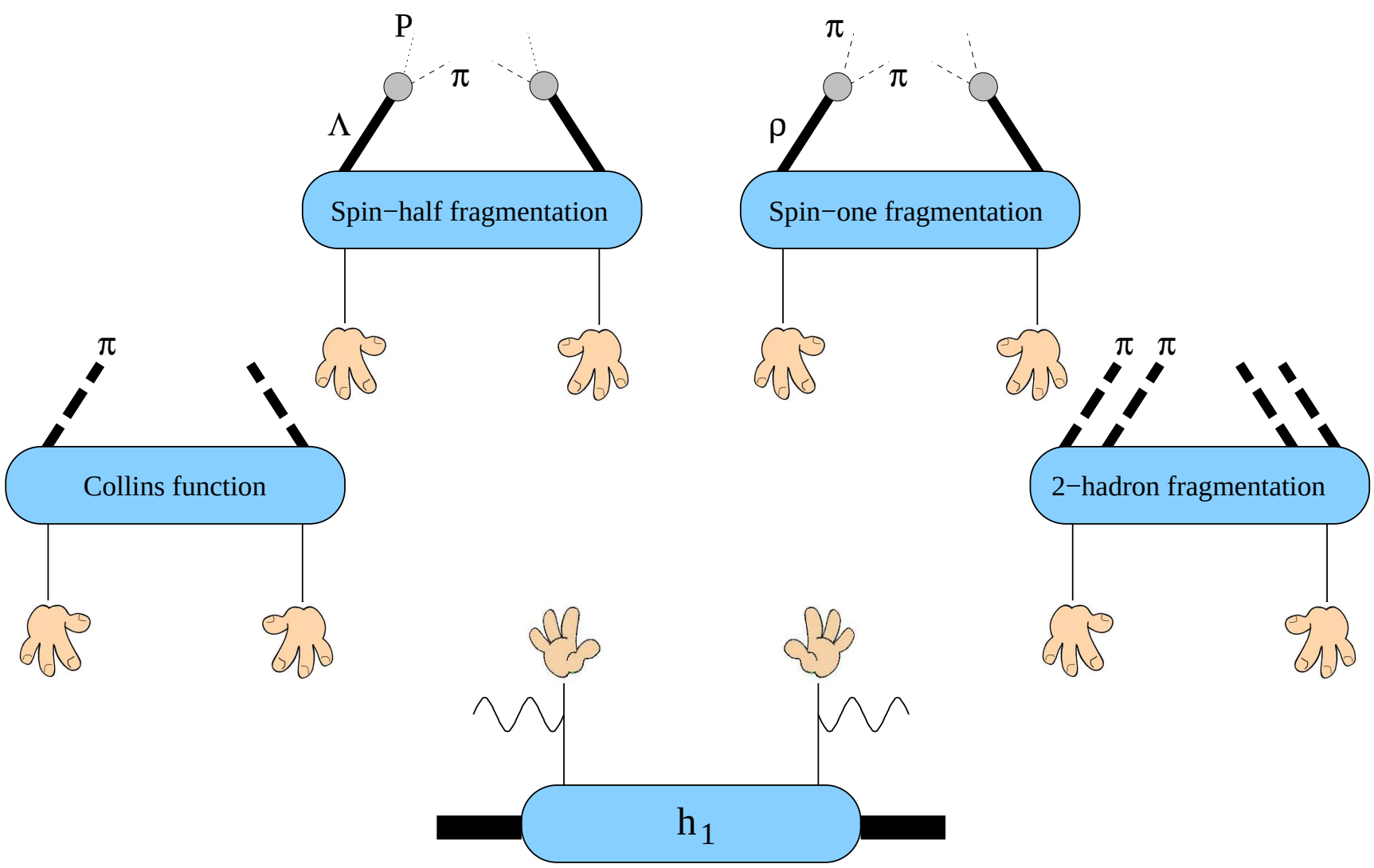
$$f_1(x) \geq 0$$

$$|g_1(x)| \leq f_1(x)$$

Positivity of eigenvalues $\Rightarrow$ less trivial (Soffer) bound

$$|h_1(x)| \leq \frac{1}{2}\left(f_1(x) + g_1(x)\right)$$

Chiral-odd partners for the transversity distribution



Fragmentation to spin-1 hadrons

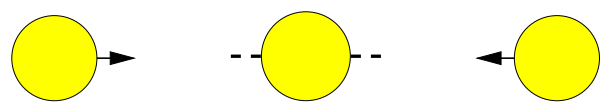
We need to measure the polarization of the outgoing hadron by means of VECTOR AND TENSOR ANALYZING POWERS of the decay

$$\Delta_{ij}(k, P_h, A_h^i, A_h^{ij}) = \int \frac{\mathrm{d}^4\xi}{(2\pi)^4}\, \mathrm{e}^{-\mathrm{i}k\cdot\xi} \langle 0|\psi_i(0)|P_h, A_h^i, A_h^{ij}\rangle\langle P_h, A_h^i, A_h^{ij}|\overline{\psi}_j(\xi)|0\rangle$$

or as a matrix in the hadron's helicity space

$$\Delta_{ij;\Lambda\Lambda'}(k, P_h) = \int \frac{\mathrm{d}^4\xi}{(2\pi)^4}\, \mathrm{e}^{-\mathrm{i}k\cdot\xi} \langle 0|\psi_i(0)|P_h, \Lambda\rangle\langle P_h, \Lambda'|\overline{\psi}_j(\xi)|0\rangle$$

where Λ and Λ' are hadron's helicities (1,0,-1)



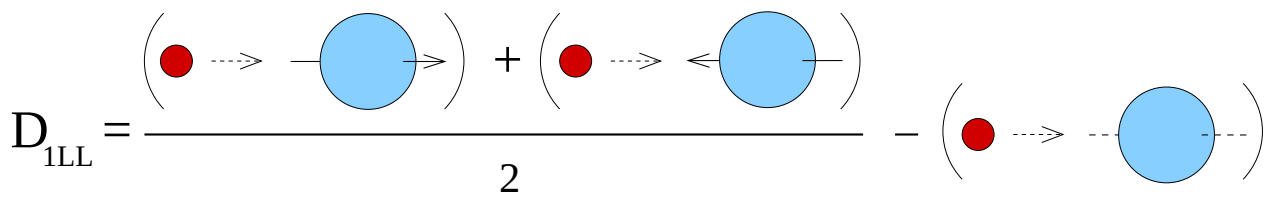
$$\Delta(k, P_h, A_h^i, A_h^{ij}) = \mathrm{Tr}\left[R_h^{(\mathrm{decay})}(A_h^i, A_h^{ij})\ \Delta(k, P_h)\right]$$

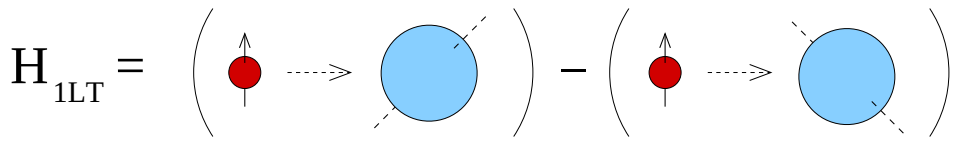
General decomposition of correlation function

Leading order in $\frac{1}{Q}$

no transverse momentum

$$\Delta(z) = \frac{1}{2}\Big\{ D_1(z) + G_1(z)\,\gamma_5\, A_L + H_1(z)\,\gamma_5\, \not{A}_T + D_{1LL}(z)\, A_{LL} + \mathrm{i}\, H_{1LT}(z)\, \not{A}_{LT} \Big\}\gamma^+$$

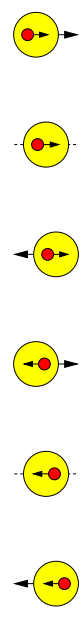




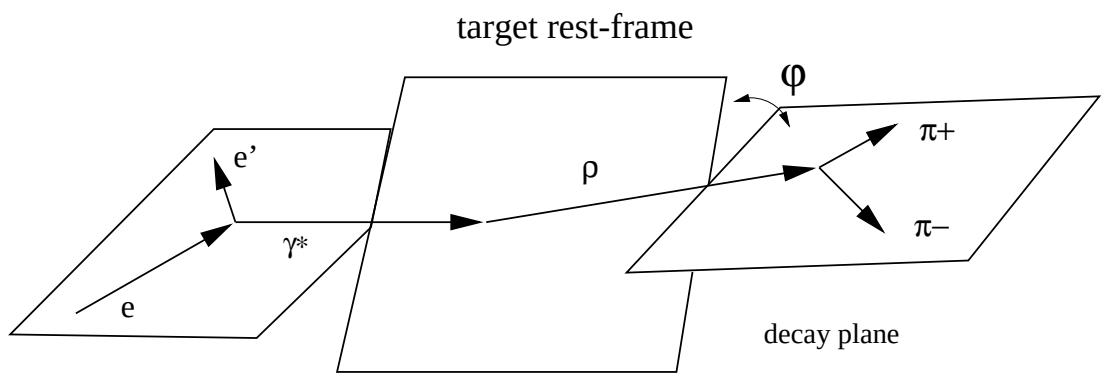
As a matrix in quark+hadron spin space

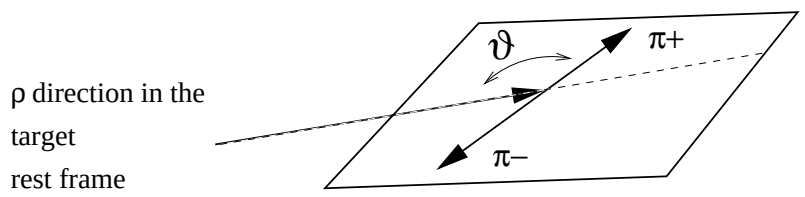
Full scattering matrix M for $q \to \rho + X$

$$
\begin{pmatrix}
D_1 + G_1 - \frac{D_{1LL}}{3} & 0 & 0 & 0 & \sqrt{2}\,(H_1 + \mathrm{i}H_{1LT}) & 0 \\
0 & D_1 + \frac{2\,D_{1LL}}{3} & 0 & 0 & 0 & \sqrt{2}\,(H_1 - \mathrm{i}H_{1LT}) \\
0 & 0 & D_1 - G_1 - \frac{D_{1LL}}{3} & 0 & 0 & 0 \\
0 & 0 & 0 & D_1 - G_1 - \frac{D_{1LL}}{3} & 0 & 0 \\
\sqrt{2}\,(H_1 - \mathrm{i}H_{1LT}) & 0 & 0 & 0 & D_1 + \frac{2\,D_{1LL}}{3} & 0 \\
0 & \sqrt{2}\,(H_1 + \mathrm{i}H_{1LT}) & 0 & 0 & 0 & D_1 + G_1 - \frac{D_{1LL}}{3}
\end{pmatrix}
$$



Analyzing powers





$$
\begin{aligned}
A_{LL} &= \tfrac{1}{3}\left(\cos^2\theta + \cos 2\theta\right) \\
A^x_{LT} &= -\sin 2\theta\ \cos\varphi & A^y_{LT} &= -\sin 2\theta\ \sin\varphi \\
A^{xx}_{TT} &= -\sin^2\theta\ \cos 2\varphi & A^{xy}_{TT} &= -\sin^2\theta\ \sin 2\varphi
\end{aligned}
$$

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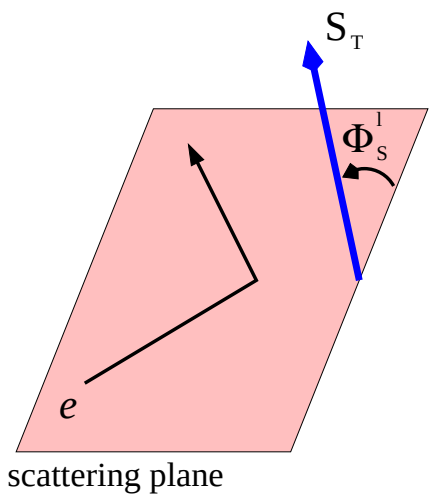
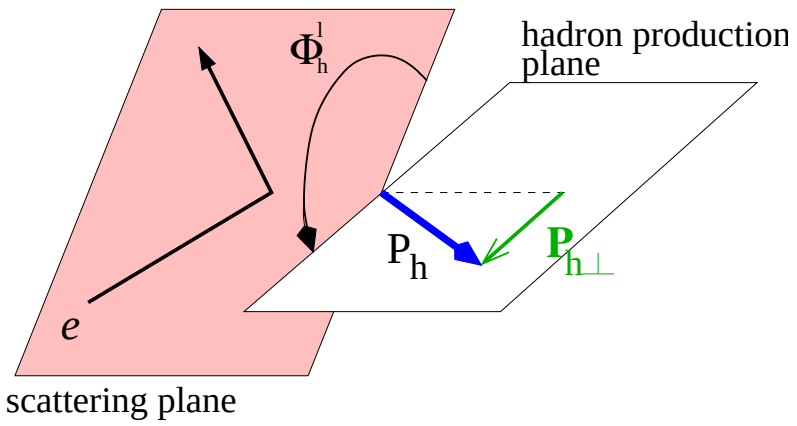
Transversity and asymmetries

$$\mathrm{d}\sigma^{\uparrow} - \mathrm{d}\sigma^{\downarrow} = f\left(\phi_S^{\ell}, |P_{h\perp}|, \phi_h^{\ell}, x_{_B}, z_h\right)$$

$$\left\langle \begin{matrix}\text{Weight}\\ \text{Function}\end{matrix} \right\rangle_T (x_{_B}, z_h) = \int \mathrm{d}\phi_S^{\ell}\ \mathrm{d}|P_{h\perp}|\ \mathrm{d}\phi_h^{\ell}\ \begin{matrix}\text{Weight}\\ \text{Function}\end{matrix}\ \left(\mathrm{d}\sigma^{\uparrow} - \mathrm{d}\sigma^{\downarrow}\right)$$

$$\propto\ h_1(x_{_B}) \times \text{?}\,(z_h)$$



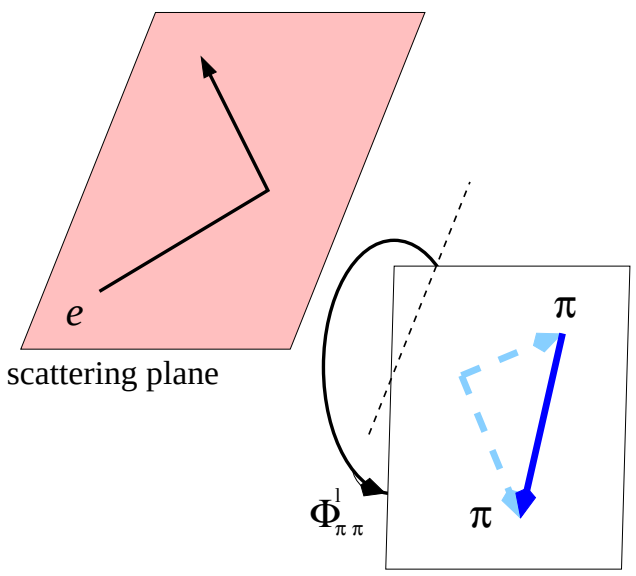


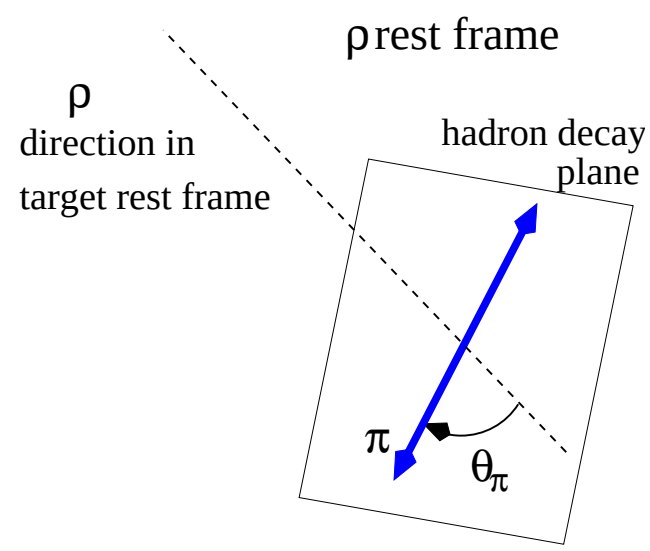
Asymmetries in spin-1 production

$$\mathrm{d}\sigma^{\uparrow} - \mathrm{d}\sigma^{\downarrow} = f\left(\theta_\pi, \phi^\ell_{\pi\pi}, \phi^\ell_S, |P_{\rho\perp}|, \phi^\ell_\rho, x_{_B}, z_\rho\right)$$

$$\left\langle \begin{matrix}\text{Weight}\\ \text{Function}\end{matrix} \right\rangle_T (x_{_B}, z_\rho) = \int \mathrm{d}\theta_\pi\ \mathrm{d}\phi^\ell_{\pi\pi}\ \mathrm{d}\phi^\ell_S\ \mathrm{d}|P_{\rho\perp}|\ \mathrm{d}\phi^\ell_\rho \begin{matrix}\text{Weight}\\ \text{Function}\end{matrix} \left(\mathrm{d}\sigma^{\uparrow} - \mathrm{d}\sigma^{\downarrow}\right)$$

$$\propto\ h_1(x_{_B}) \times \text{?}\,(z_\rho)$$





Asymmetries containing the transversity distribution

$$\left\langle \sin\left(\phi^{\ell}_{\pi\pi} + \phi^{\ell}_{S}\right) \right\rangle_{T} \propto h_1(x_B)\, H_{1LT}(z_\rho)$$

χ-odd T-odd

- Requires measurement of $\phi^{\ell}_{\pi\pi}, \phi^{\ell}_{S}, x_B, z_\rho$
- No complications due to presence of hadron transverse momentum

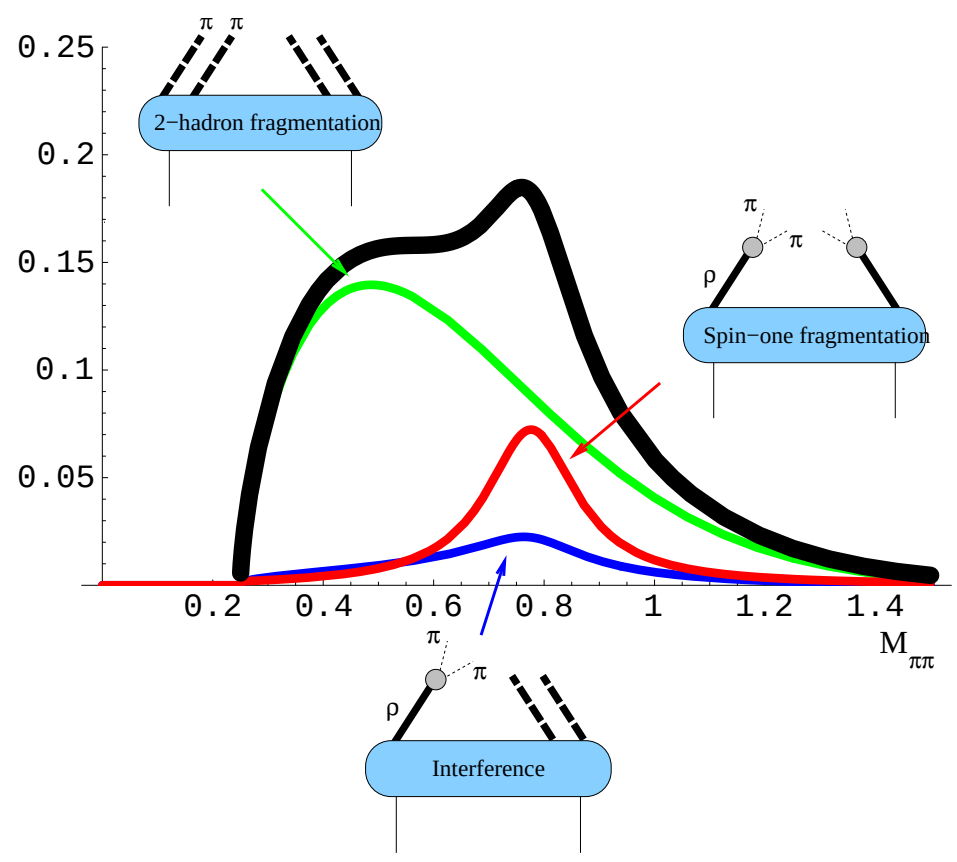
$$\left\langle \frac{|P_{\rho\perp}|}{2\, z\, M_h} \sin\left(2\phi^{\ell}_{\pi\pi} + \phi^{\ell}_{S} - \phi^{\ell}_{\rho}\right) \right\rangle_{T} \propto h_1(x_B)\, H^{(1)}_{1TT}(z_\rho)$$

$$\left\langle \frac{|P_{\rho\perp}|}{2\, z\, M_h} \sin\left(\phi^{\ell}_{\rho} + \phi^{\ell}_{S}\right) \right\rangle_{T} \propto h_1(x_B)\, H^{\perp(1)}_{1LL}(z_\rho)$$

$$\left\langle \frac{|P_{\rho\perp}|^2}{2\, z^2\, M_h^2} \sin\left(\phi^{\ell}_{\pi\pi} - \phi^{\ell}_{S} - 2\phi^{\ell}_{\rho}\right) \right\rangle_{T} \propto h_1(x_B)\, H^{\perp(2)}_{1LT}(z_\rho)$$

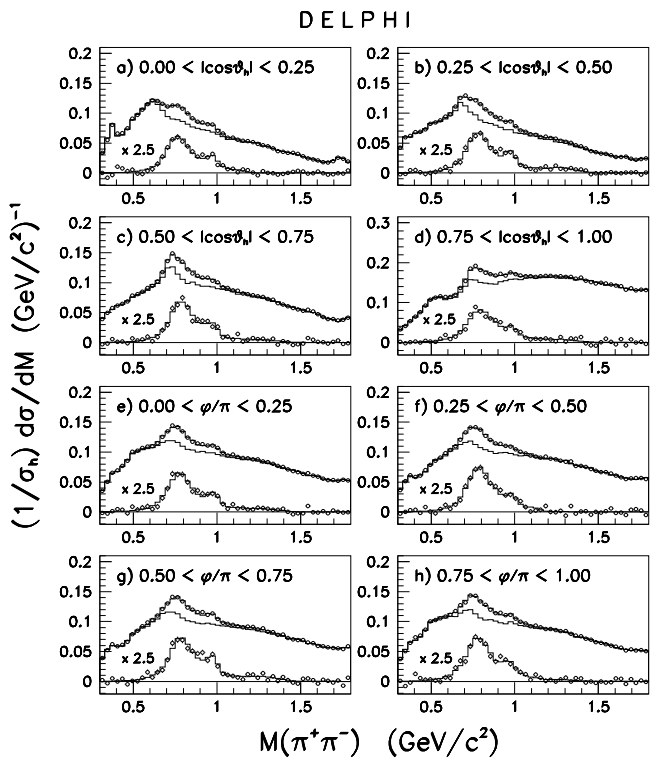
Where to observe these functions?

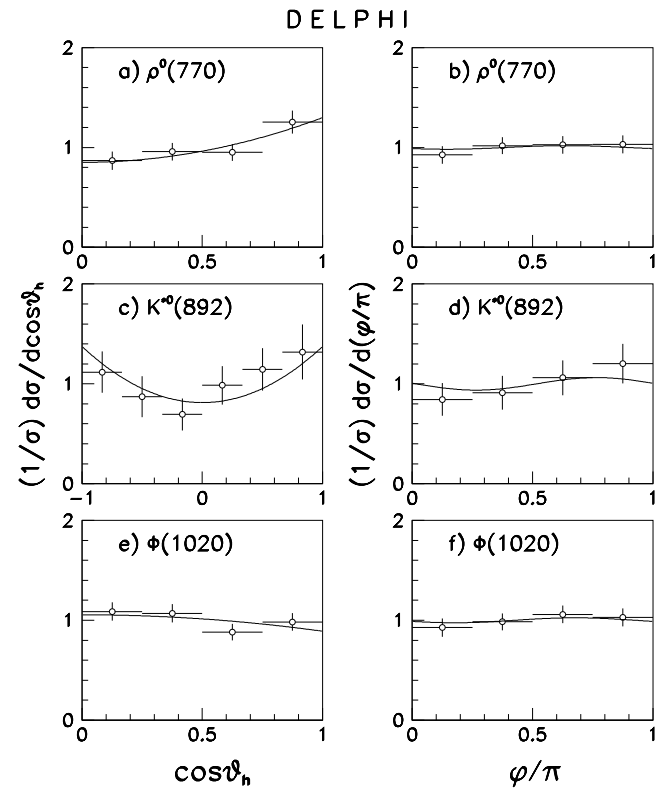
Mass spectrum of pion couples



Measurements in $e^+ e^-$ annihilation

Delphi collaboration Phys. Lett. B 406 (1997) 271



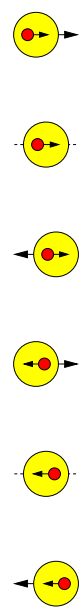


Measurement of decay distribution seems to be feasable.

Quark+hadron spin structure

Full scattering matrix M for $q \to \rho + X$

$$
\begin{pmatrix}
D_1 + G_1 - \frac{D_{1LL}}{3} & 0 & 0 & 0 & \sqrt{2}\,(H_1 + \mathrm{i}H_{1LT}) & 0 \\
0 & D_1 + \frac{2\,D_{1LL}}{3} & 0 & 0 & 0 & \sqrt{2}\,(H_1 - \mathrm{i}H_{1LT}) \\
0 & 0 & D_1 - G_1 - \frac{D_{1LL}}{3} & 0 & 0 & 0 \\
0 & 0 & 0 & D_1 - G_1 - \frac{D_{1LL}}{3} & 0 & 0 \\
\sqrt{2}\,(H_1 - \mathrm{i}H_{1LT}) & 0 & 0 & 0 & D_1 + \frac{2\,D_{1LL}}{3} & 0 \\
0 & \sqrt{2}\,(H_1 + \mathrm{i}H_{1LT}) & 0 & 0 & 0 & D_1 + G_1 - \frac{D_{1LL}}{3}
\end{pmatrix}
$$



Positivity bounds

Positive definiteness of this matrix requires:

A.B., P.J. Mulders, hep-ph/0104176

Positivity of diagonal elements $\Rightarrow$ trivial bounds

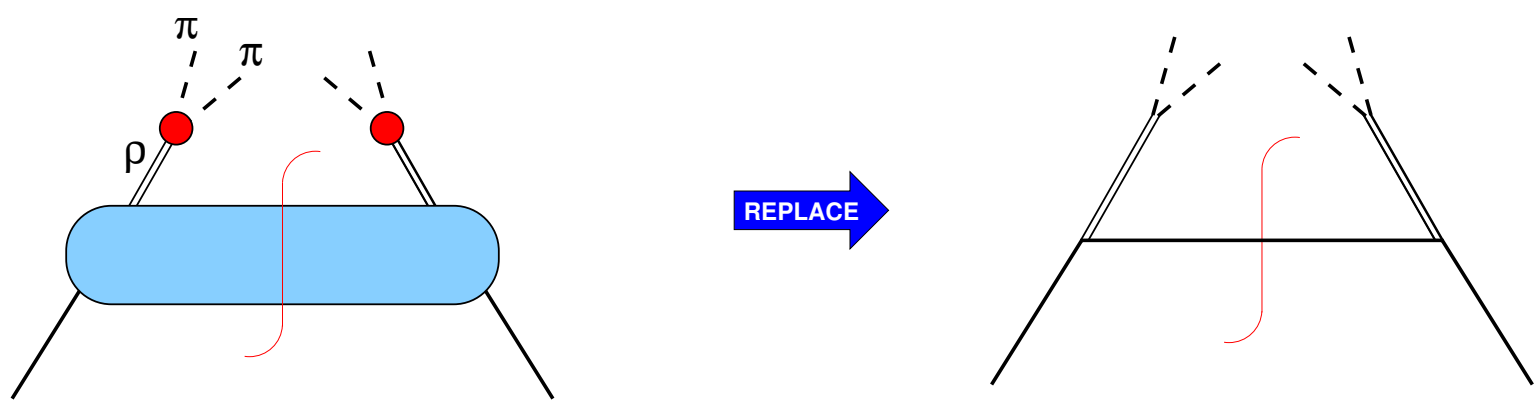
$$
\begin{aligned}
D_1(z) &\geq 0 \\
-\frac{3}{2}\, D_1(z) &\leq D_{1LL}(z) \leq 3\, D_1(z) \\
|G_1(z)| &\leq D_1(z) - \frac{1}{3}\, D_{1LL}(z) \leq \frac{3}{2}\, D_1(z)
\end{aligned}
$$

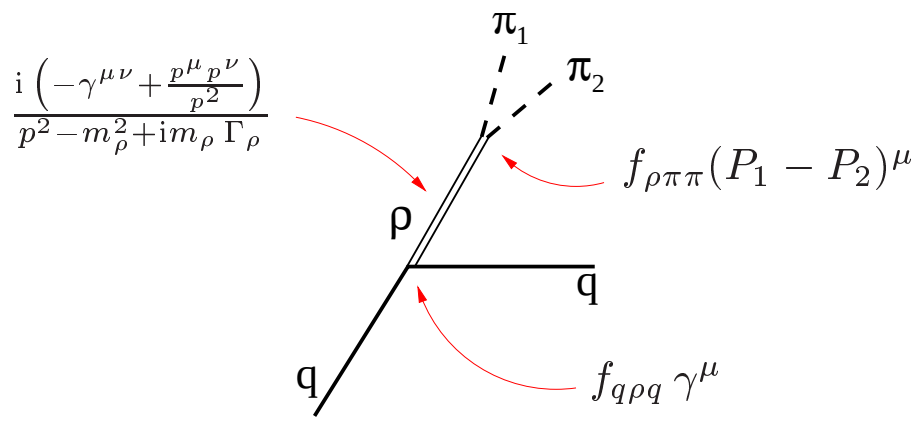
Positivity of 2-dimensional minors $\Rightarrow$ (Soffer-like) bound

$$
[H_{1LT}(z)]^2 \leq \left[D_1(z) + \frac{2}{3}\, D_{1LL}(z)\right] \left[D_1(z) - \frac{1}{3}\, D_{1LL}(z)\right] \leq \frac{9}{2}\, [D_1(z)]^2 .
$$

Generating non-zero T-odd functions

Take a simple model for the fragmentation process



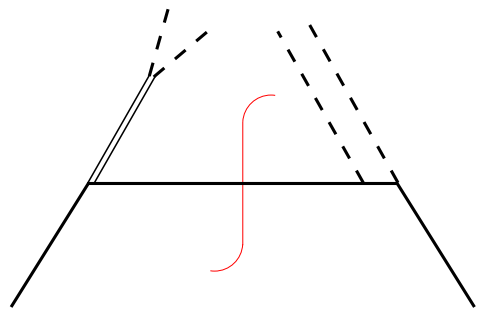


Generating non-zero T-odd functions

Add one-loop corrections (see also talk by Rajen Kundu)

T-odd functions $\neq 0$

"Competing" or "complementary" mechanism: interference fragmentation functions (see talk by Marco Radici)



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Conclusions

- Fragmentation to spin-1 hadron is a viable option for transversity measurements, at the same level as Collins function and interference fragmentation functions.

- Since polarimetry on the final hadron requires the observation of two decay products, it represents a specific contribution to the more general case of two hadron production.

- Information on the unknown fragmentation functions can be extracted from positivity bounds and simple models.