

WG4 Discussion on Warped-space KK gauge bosons

Shrihari Gopalakrishna

Brookhaven National Lab

LHC2FC Workshop, CERN

17th February 2009

- Focus on heavy EW spin-1 resonances
- Warped (RS) model
 - $SU(3)_{QCD} \times SU(2)_L \times SU(2)_R \times U(1)_X$ bulk gauge group
 - Precision electroweak observables require $M_{Z'}$, $M_{W_1^\pm} \gtrsim 2$ TeV
 - Makes discovery challenging at the LHC
- Equivalently, 4-D strongly coupled theory (AdS/CFT Corresp.)
- What are the general issues at colliders?

5D Warped Space

[Randall, Sundrum, 99]

$$ds^2 = e^{-2k|y|}(\eta_{\mu\nu} dx^\mu dx^\nu) + dy^2$$

Z_2 Orbifold -

- Planck (UV) Brane
- TeV (IR) Brane

R : radius of Ex. Dim.

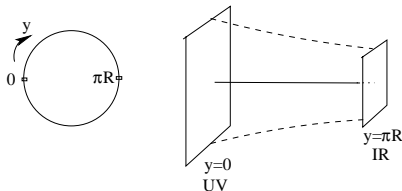
k : curvature

Hierarchy prob soln:

- TeV Brane Higgs : $M_{EW} \sim ke^{-k\pi R}$: Choose $k\pi R \sim 34$

Bulk fields \rightarrow AdS/CFT

Bulk Fermions explain flavor (in FCNC safe way)



EWPT constraint

[Csaki, Erlich, Terning - 02] [Agashe, Delgado, May, Sundrum - 03]

Bulk $SU(2)_L \times SU(2)_R \times U(1)_X$ and 3rd gen quarks (2, 2)

- T and $Zb\bar{b}$ protected. S constraint [Agashe, Contino, DaRold, Pomarol - 06]
- EWPT $\Rightarrow M_{Z'} \gtrsim 2 - 3 \text{ TeV}$ [Carena, Ponton, Santiago, Wagner - 06,07]

Note that: $M_{Z'} = 2 \text{ TeV} \Rightarrow M_{KK \text{ Grav}} = 3 \text{ TeV}$.

FCNC constraint

[Agashe, Perez, Soni, 04] [Casagrande, Goertz, Haish, Neubert, Pfoh, 08]

Non-trivial flavor structure needed

[Fitzpatrick, Perez, Randall - 07]

[Csaki, Falkowski, Weiler, 08]

Our study on Electroweak KK

Agashe, Davoudiasl, SG, Han, Huang, Perez, Si, Soni.

arXiv:0709.0007 & 0810.1497

For KK Glue and Graviton see:

- KK Gluon at the LHC

$L = 100 \text{ fb}^{-1}$ LHC reach is 4 TeV

[Agashe, Belyaev, Krupovnickas, Perez, Virzi, 06]
[Lillie, Randall, Wang, 07] [Lillie, Shu, Tait, 07]

- KK Graviton at LHC

$L = 300 \text{ fb}^{-1}$ LHC reach is about 2 TeV

[Agashe, Davoudiasl, Perez, Soni 07]
[Fitzpatrick, Kaplan, Randall, Wang, 07]

- Associated Prod.

[Guchait, Mahmoudi, Sridhar, 08]

[Other Refs. later]

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: $(W_L^3, W_R^3, X) \rightarrow Z'$
- Two charged gauge bosons: $(W_L^\pm, W_R^\pm) \rightarrow W'$

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: $(W_L^3, W_R^3, X) \rightarrow Z'$
- Two charged gauge bosons: $(W_L^\pm, W_R^\pm) \rightarrow W'$

Kaluza-Klein (KK) expansion: $A(x, y) = \sum_0^\infty f_n(y) A^{(n)}(x)$

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: $(W_L^3, W_R^3, X) \rightarrow Z'$
- Two charged gauge bosons: $(W_L^\pm, W_R^\pm) \rightarrow W'$

Kaluza-Klein (KK) expansion: $A(x, y) = \sum_0^\infty f_n(y) A^{(n)}(x)$

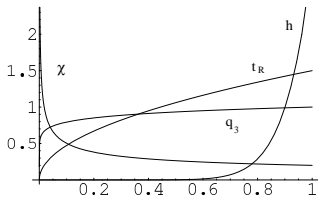
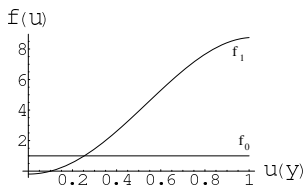
Symmetry breaking

- $SU(2)_R \times U(1)_X \rightarrow U(1)_Y$ & $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$
 - Mixes zero-mode and KK-states
 - Mass eigenstates Z' and W'

Wave functions

Bulk field EOM gives profiles in extra-dimension

Fermion bulk mass (c parameter) controls localization



Overlap integral gives $\psi\psi V_\mu$: $\mathcal{I} = \int [dy] g_\psi(y)^2 f_V(y)$

Compared to SM

- Z' couplings to h enhanced (also V_L by Equivalence Theorem!)
- Z' couplings to t_R enhanced
- Z' couplings to χ suppressed

Define: $\xi \equiv \sqrt{k\pi R} \approx 5$

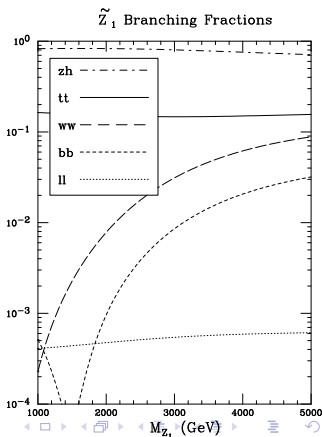
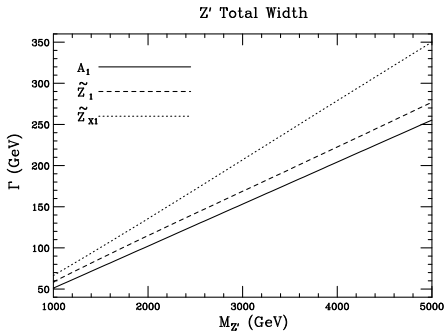
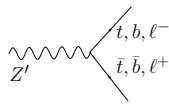
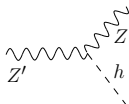
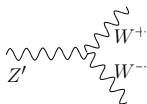
Z' overlap with Higgs $\rightarrow \xi$

Z' overlap with fermions:

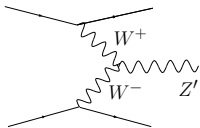
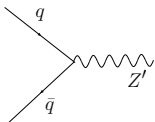
	Q_L^3	t_R	light fermions
\mathcal{I}^+	1	ξ	$-\frac{1}{\xi}$
\mathcal{I}^-	1	ξ	0

$$\bar{\psi}_{L,R} \gamma^\mu \left[eQ\mathcal{I}A_{1\mu} + g_Z (T_L^3 - s_W^2 T_Q) \mathcal{I}Z_{1\mu} + g_{Z'} (T_R^3 - s'^2 T_Y) \mathcal{I}Z_{X1\mu} \right] \psi_{L,R}$$

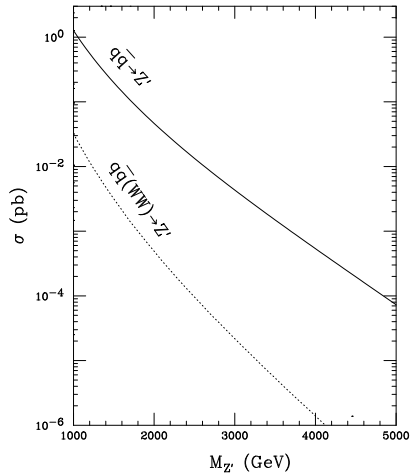
Z' decays and BR



Z' production c.s. (LHC)



Total Z' Cross Section at LHC

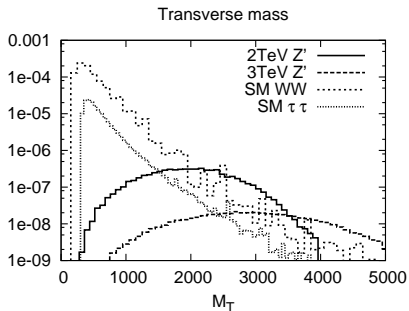
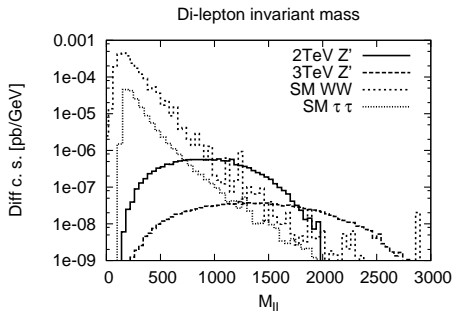


\mathcal{L} needed : (... , ...) fb^{-1} for (2 TeV, 3 TeV) is shown below

- $pp \rightarrow Z' \rightarrow W^+ W^-$ (100, 1000) fb^{-1}
 - Fully leptonic : $W \rightarrow \ell\nu$; $W \rightarrow \ell\nu$
 - Semi leptonic : $W \rightarrow \ell\nu$; $W \rightarrow (jj)$
- $pp \rightarrow Z' \rightarrow Z h$
 - $m_h = 120\text{GeV}$: $Z \rightarrow \ell^+ \ell^-$; $h \rightarrow b\bar{b}$ (200, 1000) fb^{-1}
 - $m_h = 150\text{GeV}$:
 $Z \rightarrow (jj)$; $h \rightarrow W^+ W^- \rightarrow (jj) \ell\nu$ ($< 100, 300$) fb^{-1}
- $pp \rightarrow Z' \rightarrow \ell^+ \ell^-$
 - Clean but needs high luminosity (1000, -) fb^{-1}
- $pp \rightarrow Z' \rightarrow t\bar{t}, b\bar{b}$ [Djouadi, Moreau, Singh - 07]
 - KK gluon degenerate

$$pp \rightarrow Z' \rightarrow W^+W^- \rightarrow \ell\nu\ell\nu$$

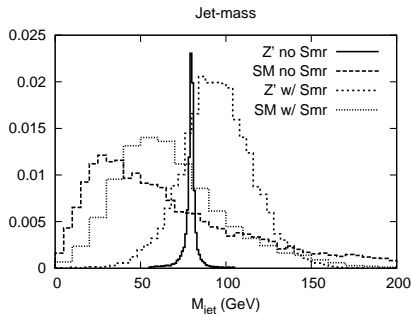
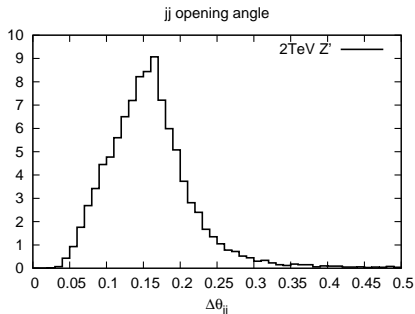
2 ν 's \Rightarrow cannot reconstruct event



$$M_{eff} \equiv p_{T_{\ell_1}} + p_{T_{\ell_2}} + \cancel{p}_T \quad M_{T_{WW}} \equiv 2\sqrt{p_{T_{\ell\ell}}^2 + M_{\ell\ell}^2}$$

\mathcal{L} needed: 100 fb^{-1} (2 TeV) ; 1000 fb^{-1} (3 TeV)

$$pp \rightarrow Z' \rightarrow W^+W^- \rightarrow \ell \nu jj$$



jj Collimation implies forming m_W nontrivial : use jet-mass

In our study: Jet-mass after Parton shower in Pythia

[Discussions w/ Frank Paige] [Almeida, Lee, Perez, Sterman, Sung, Virzi, 08]

To account for (HCal) expt. uncert.

Smearing by $\delta E = 80\%/\sqrt{E}$; $\delta\eta, \delta\phi = 0.05$

Tracker + ECal (2 cores?) have better resolutions

[F. Paige] [M. Strassler]

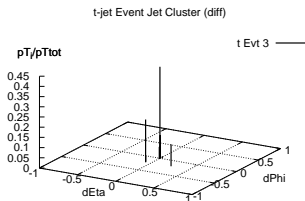
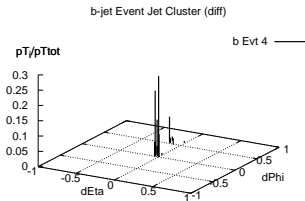
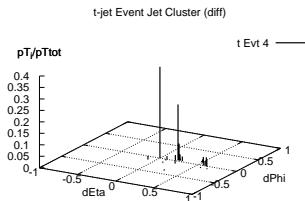
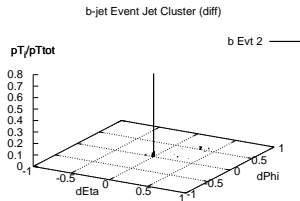
\mathcal{L} needed: 100 fb^{-1} (2 TeV) ; 1000 fb^{-1} (3 TeV)

$$W'^{\pm} \rightarrow t b \rightarrow l \nu b b$$

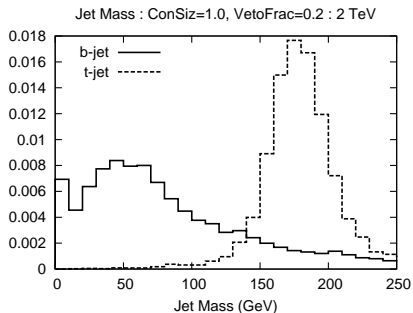
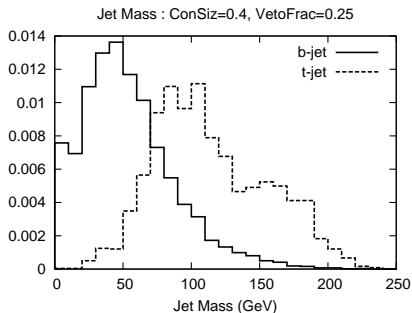
Signal c.s. $\sim 1fb$

Bkgnd is single top + QCD W b b AND ...

$t\bar{t}$: hadronically decaying top can fake a b



$$W'^{\pm} \rightarrow t b \rightarrow l \nu b b$$



Jet-mass cut: cone size 1.0 and $0 < j_M < 75 \Rightarrow 0.4\%$ of *top* fakes *b*
 \mathcal{L} needed: 100 fb^{-1} (2 TeV)

Conclusions

- Probing Z' and W' at the LHC challenging but possible
 - KK Grav needs large luminosity : SLHC?
- Improve reach by better “merged” jet reconstruction (jet mass)

Future Colliders (SLHC, VLHC, ILC, CLIC, muon collider)

- Can better detector resolutions give more suppression of QCD?
- Increased luminosity can give reasonable event samples
- What happens to signal vs. bkgnd with increasing \sqrt{s} ?
- e^+e^- needs to have adequate \sqrt{s} (EWPT constraints)

BACKUP SLIDES

EWPT

Bulk $SU(2)_L \times U(1)_Y$

- $\log \frac{M_{Pl}}{M_{EW}}$ enhanced contributions to S, T [Csaki, Erlich, Terning - 02]
- $Zb\bar{b}$ coupling shifted

Bulk $SU(2)_L \times SU(2)_R \times U(1)_X$

- T-parameter protected [Agashe, Delgado, May, Sundrum - 03]

Bulk $SU(2)_L \times SU(2)_R \times U(1)_X$ and 3rd gen quarks (2, 2)

- T and $Zb\bar{b}$ protected [Agashe, Contino, DaRold, Pomarol - 06]

EWPT $\Rightarrow M_{Z'} \gtrsim 2 - 3 \text{ TeV}$

[Carena, Ponton, Santiago, Wagner - 06,07]

FCNC constraint

Non-trivial flavor structure needed For MFV example see [Fitzpatrick, Perez, Randall - 07]

EW Gauge Sector

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: (W_L^3, W_R^3, X)
- Two charged gauge bosons: (W_L^\pm, W_R^\pm)

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: (W_L^3, W_R^3, X)
- Two charged gauge bosons: (W_L^\pm, W_R^\pm)

- $SU(2)_R \times U(1)_X \rightarrow U(1)_Y : (W_L^3, W_R^3, X) \rightarrow (W_L^3, B, Z_X)$
 - $Z_X \equiv \frac{1}{\sqrt{g_X^2 + g_R^2}}(g_R W_R^3 - g_X X) \rightarrow (-, +) ; W_R^\pm \rightarrow (-, +)$
 - $B \equiv \frac{1}{\sqrt{g_X^2 + g_R^2}}(g_X W_R^3 + g_R X) \rightarrow (+, +) ; W_L^\pm \rightarrow (+, +)$

Symm breaking by BC: $Z_X(-, +)$ means $Z_X|_{y=0} = 0 ; \partial_y Z_X|_{y=\pi R} = 0$

Bulk Gauge group : $SU(2)_L \times SU(2)_R \times U(1)_X$

- Three neutral gauge bosons: (W_L^3, W_R^3, X)
- Two charged gauge bosons: (W_L^\pm, W_R^\pm)
- $SU(2)_R \times U(1)_X \rightarrow U(1)_Y$: $(W_L^3, W_R^3, X) \rightarrow (W_L^3, B, Z_X)$
 - $Z_X \equiv \frac{1}{\sqrt{g_X^2 + g_R^2}}(g_R W_R^3 - g_X X) \rightarrow (-, +)$; $W_R^\pm \rightarrow (-, +)$
 - $B \equiv \frac{1}{\sqrt{g_X^2 + g_R^2}}(g_X W_R^3 + g_R X) \rightarrow (+, +)$; $W_L^\pm \rightarrow (+, +)$

Symm breaking by BC: $Z_X(-, +)$ means $Z_X|_{y=0} = 0$; $\partial_y Z_X|_{y=\pi R} = 0$

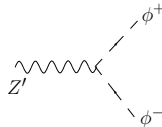
- $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$: $(W_L^3, B, Z_X) \rightarrow (A, Z, Z_X)$
 - By TeV brane Higgs

EWSB induced $Z'W^+W^-$ coupling

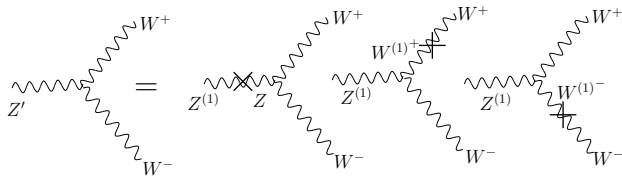
$Z^{(1)}V^{(0)}V^{(0)}$ is zero by orthogonality ...

... but induced after EWSB

Using Goldstone equivalence:



In Unitary Gauge:



Even though $\xi \cdot (\frac{v}{M_{KK}})^2$ suppressed ...

... can be overcome by $(\frac{M_{KK}}{m_Z})^2$ (from long. pol. vectors)

Case (i): Channel	$M_{W'_1}$ (TeV)	\mathcal{L} (fb^{-1})	S/B	S/\sqrt{B}
$t b \rightarrow \ell \nu b \bar{b}$	3	300	5.8 (0.9)	0.995 (0.95) CL
$Z W \rightarrow \ell \ell \ell \nu$	3	1000	6	0.99 CL
$m_h = 120: W h \rightarrow \ell \nu b \bar{b}$	3	300	1	0.99 CL
$m_h = 150: W h \rightarrow (ij) \ell \nu (jj)$	3	300	2	0.9987 CL
Case (ii): Channel	$M_{W'_1}$ (TeV)	\mathcal{L} (fb^{-1})	S/B	S/\sqrt{B}
$t b \rightarrow \ell \nu b \bar{b}$	2	1000	0.4 (0.2)	3.4 (2.5) σ
$Z W \rightarrow \ell \ell \ell \nu$	3	1000	10	> 0.9995 CL
$m_h = 120: W h \rightarrow \ell \nu b \bar{b}$	3	300	2.4	> 0.9995 CL
$m_h = 150: W h \rightarrow (ij) \ell \nu (jj)$	3	300	4	$\gg 0.9995$ CL