

Bhabha scattering at the ILC

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What is the ILC?

The International Linear Collider will be the next e^+e^- collider based on superconducting technology:

- first phase: $\sqrt{s} \leq 500$ GeV
- upgrade: $\sqrt{s} \approx 1$ TeV
- luminosity $\mathcal{L} \approx 3 - 5 \cdot 10^{34} \text{cm}^{-2}\text{s}^{-1} \Rightarrow \sim 300 - 500 \text{fb}^{-1}/\text{year}$
- polarised electron beams ($P = 80\text{-}90\%$) and, as an option, polarised positron beams ($P = 40\text{-}60\%$).
- GigaZ option: 10^9 events at the Z pole with polarised beams
- Time scale:
 - Conceptual design by end 2006
 - International Liner Collider Technical Design Report by 2007
 - Site selection and approval in 2008
 - Begin data taking: 2015

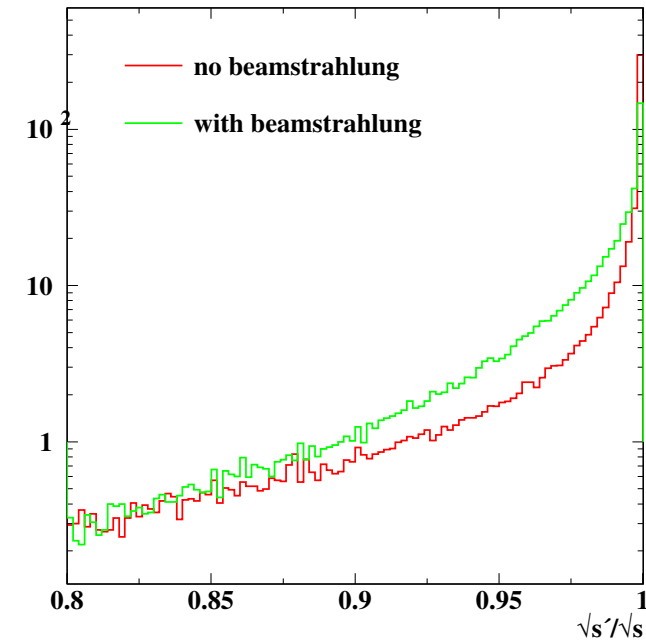
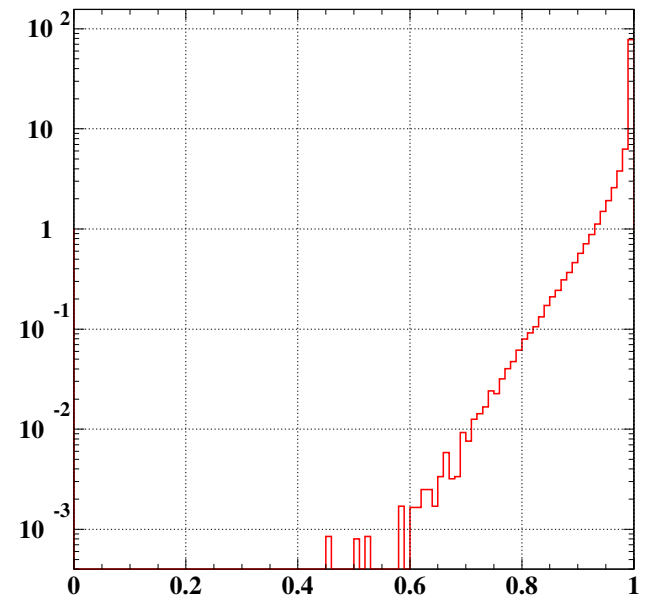
Expected statistics at the International Linear Collider (ILC)

- $\text{few} \cdot 10^4 e^+e^- \rightarrow HZ/\text{year}$ at $\sqrt{s} \approx 350 \text{ GeV}$ ($m_H \approx 120 \text{ GeV}$)
- $10^5 e^+e^- \rightarrow t\bar{t}/\text{year}$ at $\sqrt{s} \approx 350 \text{ GeV}$
- $5 \cdot 10^5 e^+e^- \rightarrow q\bar{q}/\text{year}$ at $\sqrt{s} \approx 500 \text{ GeV}$ (no rad. ret)
- $10^5 e^+e^- \rightarrow \mu^+\mu^-/\text{year}$ at $\sqrt{s} \approx 500 \text{ GeV}$ (no rad. ret)
- $10^6 e^+e^- \rightarrow W^+W^-/\text{year}$ at $\sqrt{s} = 500 - 1000 \text{ GeV}$
- $10^9 e^+e^- \rightarrow Z/\text{year}$ at $\sqrt{s} \approx 91 \text{ GeV}$

New problem at the ILC: beamstrahlung

- Beams at IP are extremely collimated with many electrons/bunch
→ very high charge density
⇒ Electrons of one bunch radiate against the coherent field of the other bunch (Beamstrahlung)
- Average energy loss for colliding e^+e^- pairs at 500 GeV: $\sim 1.5\%$
- For continuum processes beamstrahlung comparable to ISR, however with shorter tails
- Beamstrahlung has to be included in every generator!

Normalised e^- spectrum



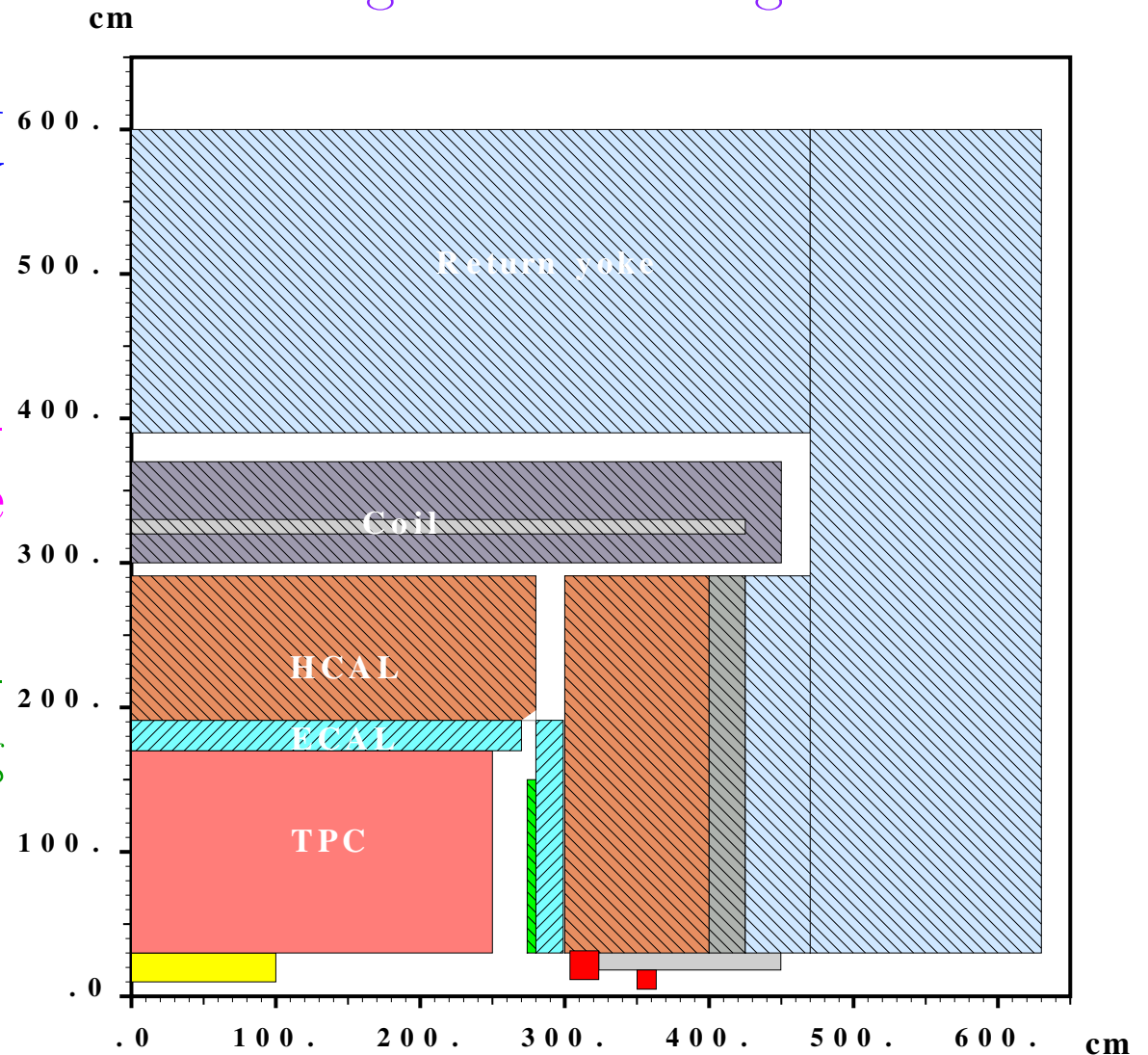
The Detector at the ILC

The concept for a large detector with gaseous tracking

Hermetic detector with full tracking and calorimetry above $\theta = 7^\circ$

Efficient electron identification in the full tracking range

Unambiguous charge identification in the full tracking range



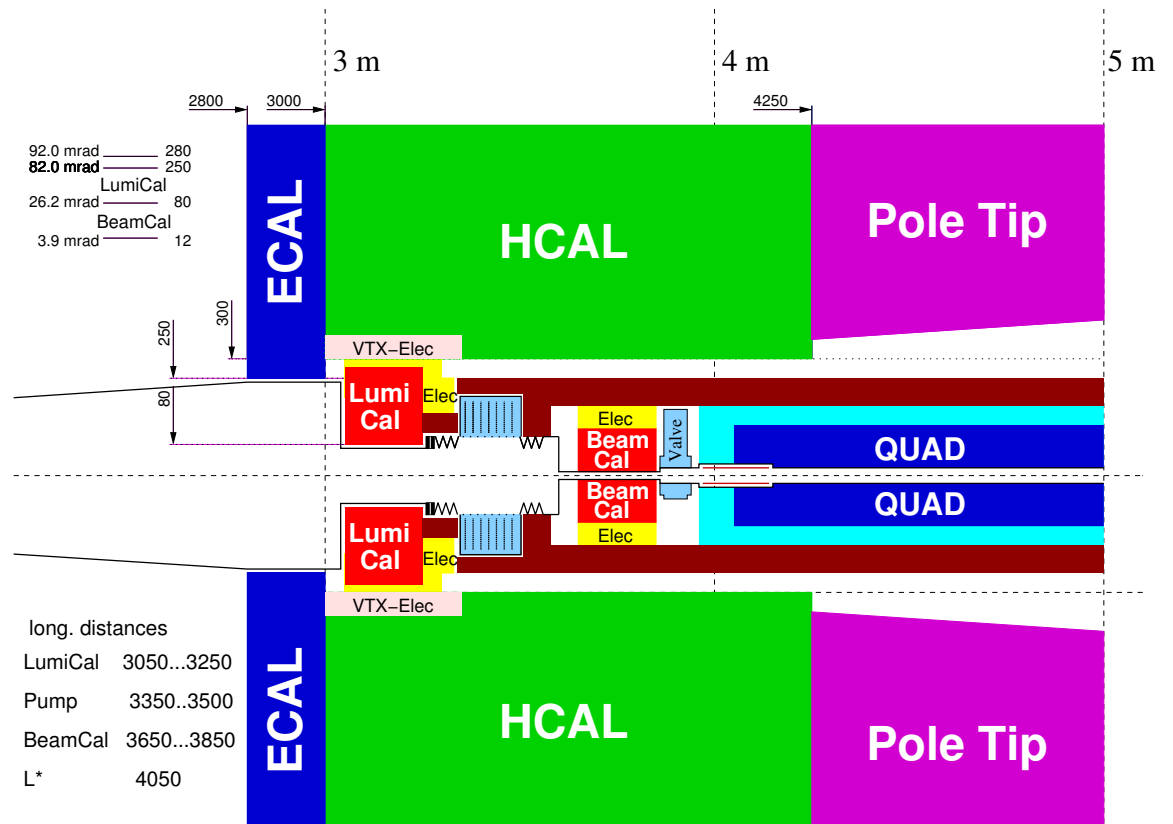
The forward region of the detector

BeamCal

- $4 \text{ mrad} < \theta < 25 \text{ mrad}$
- huge background from beam-beam interactions
- can only be used for machine tuning and $\gamma\gamma$ veto

LumiCal

- $25 \text{ mrad} < \theta < 80 \text{ mrad}$
- almost no background
- will be used for a precision luminosity measurement

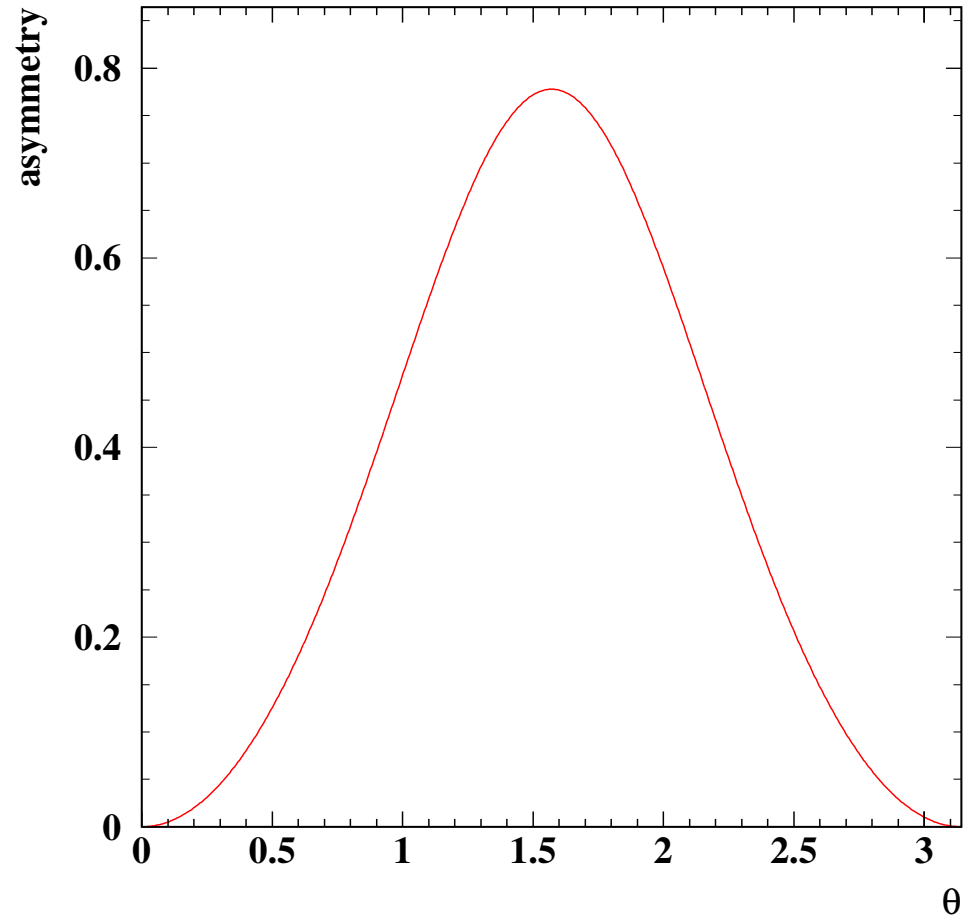


Bhabha scattering and polarisation

- At ILC both beams can be polarised
- Even for pure QED there is a cross section asymmetry between parallel and antiparallel beams proportional to $\mathcal{P}_{e^+}\mathcal{P}_{e^-}$ (ILC: $\mathcal{P}_{e^+}\mathcal{P}_{e^-} \sim 0.5$)
- This asymmetry is huge in the centre and small but sizable in the forward region
- Polarisation of both beams has to be included in the calculations and the generators

$$A = \frac{\sigma(\Rightarrow) - \sigma(\Leftarrow)}{\sigma(\Rightarrow) + \sigma(\Leftarrow)} \text{ as a function of } \theta$$

for QED only



The Luminosity Measurement at ILC

What precision do we need?

Luminosity precision is determined by statistics of interesting processes

- $e^+e^- \rightarrow W^+W^-$: ~ 10 pb at $\sqrt{s} = 340$ GeV scaling with $1/s$
⇒ $\mathcal{O}(10^6)$ events ⇒ need 10^{-3} precision
- $e^+e^- \rightarrow f\bar{f}$: ~ 5 pb at $\sqrt{s} = 340$ GeV scaling with $1/s$
⇒ $\mathcal{O}(10^6)$ events ⇒ need 10^{-3} precision
- **GigaZ**: aim for 10^9 hadronic Z decays. Relevant physics quantities (except N_ν) need also leptonic decays (10% of hadronic decays)
⇒ need 10^{-4} precision

$$e^+e^- \rightarrow f\bar{f}$$

$e^+e^- \rightarrow f\bar{f}$ is sensitive to physics at very high scales (compositeness, Z' , extra space dimensions)

Sensitivity is mainly via interference with Standard Model amplitude
 $\Rightarrow \propto 1/M^2$

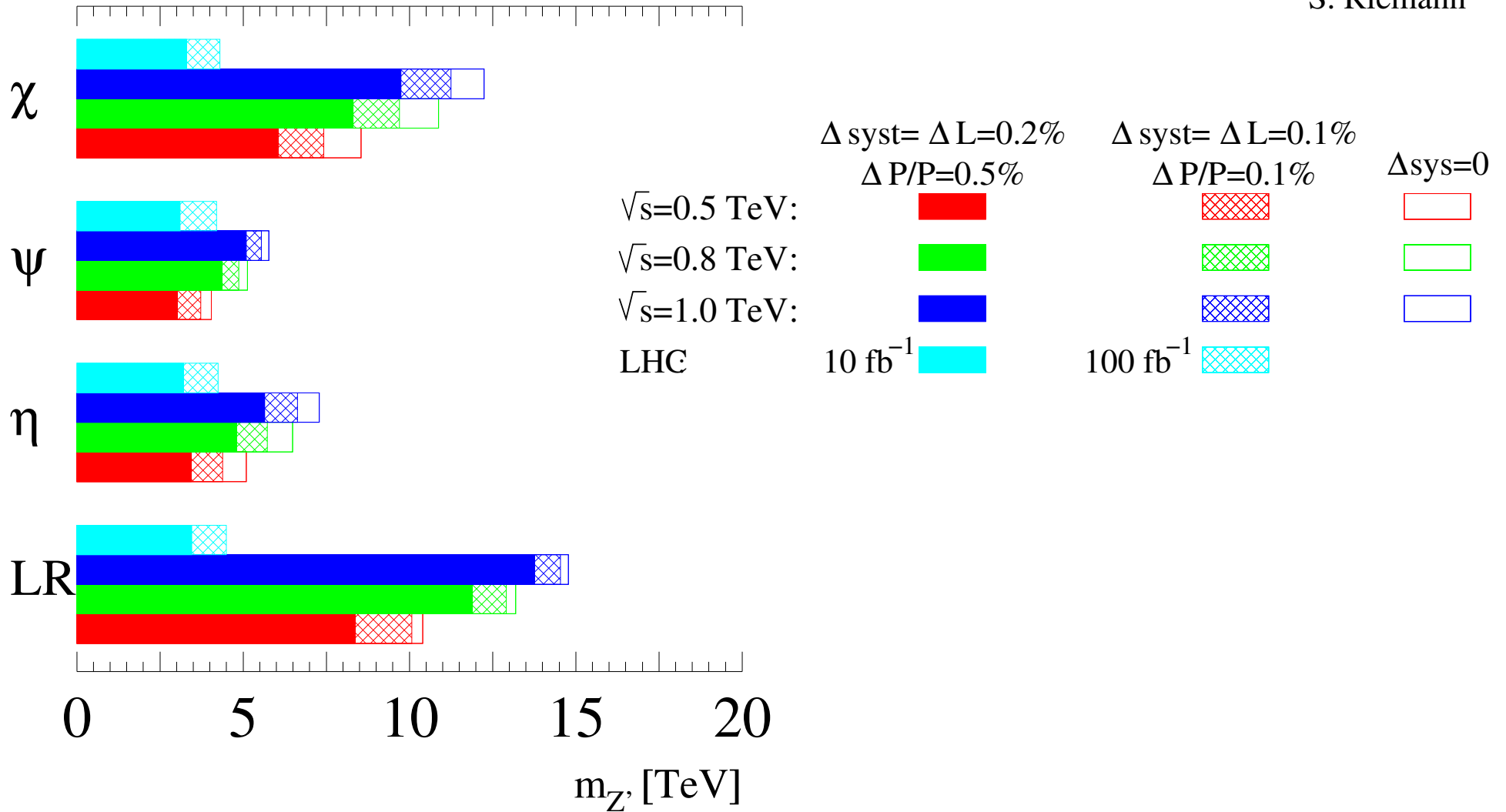
All observables (cross section, left-right asymmetry, forward-backward asymmetry) are important

Systematic errors (e.g. luminosity) effect results significantly

Z' limits in different models

$L=1\text{ab}^{-1}$ $P_{-}=0.8$ $P_{+}=0.6$

S. Riemann



GigaZ

GigaZ = 10^9 Z at $\sqrt{s} \approx m_Z$

Main aim: $\sin^2 \theta_{eff}^l$ determination \Rightarrow no \mathcal{L} dependence

Important additional information from “lineshape” parameters
 $\Gamma_Z, \sigma_0^{\text{had}}, R_l$

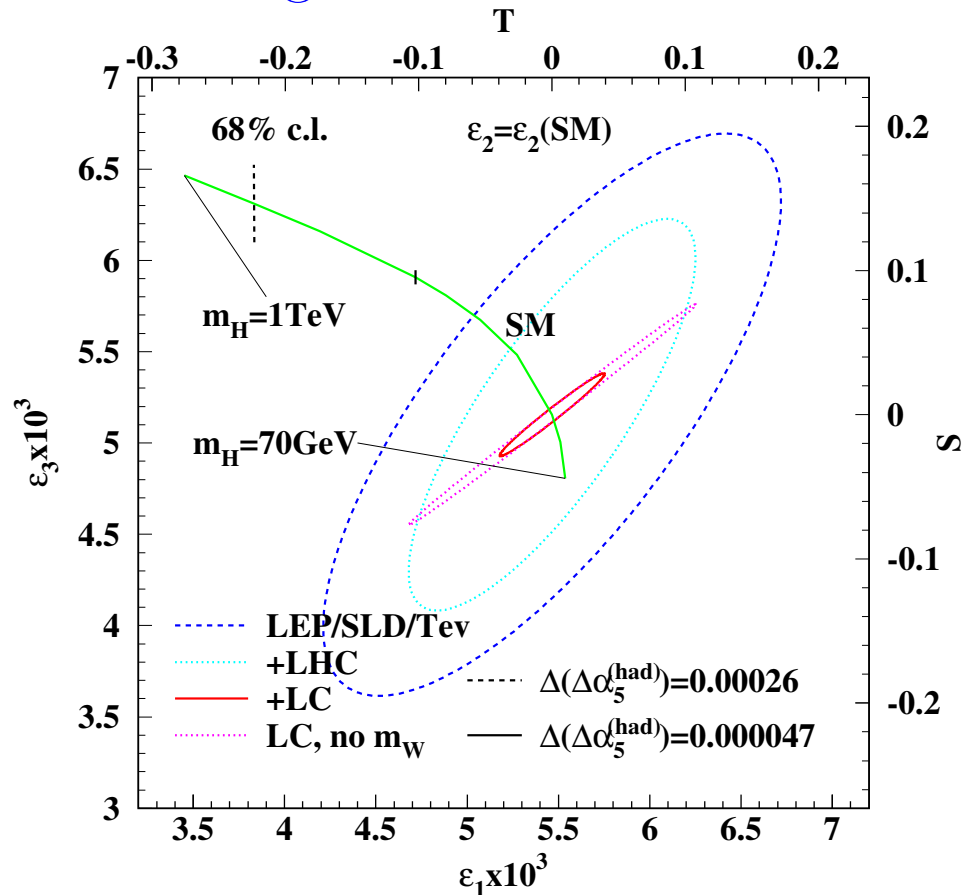
Interesting information is obtained from combination of these parameters:

$$\sigma_0^{\text{had}} = \frac{12\pi \Gamma_e \Gamma_{\text{had}}}{m_Z \Gamma_Z^2}$$
$$R_l = \frac{\Gamma_{\text{had}}}{\Gamma_l}$$

\Rightarrow need all parameters with about the same accuracy

- Γ_Z : difficult to estimate (beam energy, beamstrahlung, beamspread) but $\Delta\Gamma_Z = 1 \text{ MeV}$ ($\Delta\Gamma_Z/\Gamma_Z = 4 \cdot 10^{-4}$) seems realistic
 - R_l : $\Delta R_l/R_l = 10^{-4}$ from lepton statistics
- ⇒ need lumi error (exp+theo) $\Delta\mathcal{L}/\mathcal{L} \sim 2 \cdot 10^{-4}$

Gain of GigaZ:



- small axis: $\sin^2 \theta_{eff}^l$
⇒ no luminosity dependence
- large axis: m_W if $\varepsilon_2 = U = SM$, otherwise Γ_l
⇒ luminosity precision essential
Important in interpretations outside SM!

The Luminometer (current planning)

- Calorimeter with high granularity
- No tracking in front
- ⇒ Will do “calorimetric measurement”, i.e. no separation of nearby electron and photon
- $25 \text{ mrad} < \theta < 80 \text{ mrad}$
- All similar to LEP

Theoretical uncertainties

- Uncertainty was $\Delta\mathcal{L} = 0.05\%$ at LEP
- Sufficient for high energy, if constant with \sqrt{s}
- Definitely too large at GigaZ

Polarisation effects

- Asymmetry $3 \cdot 10^{-4}$ at 25 mrad and $3 \cdot 10^{-3}$ at 80 mrad ($\mathcal{P}_{e^+}\mathcal{P}_{e^-} = 1$)
- Marginal effect at for high energy precision
- Relevant for GigaZ
- Polarisation affects also asymmetries where luminosity normally cancels:
 - 0.1% asymmetry affects A_{LR} at GigaZ by $0.36 \cdot 10^{-4} = 1.2\sigma$ using the Blondel scheme
 - also at high energy the beam polarisation can be measured with the Blondel scheme
 - 0.1% asymmetry affects the measured asymmetry by $0.5 \cdot 10^{-3} = 0.5\sigma$

Reconstruction of the Luminosity Spectrum

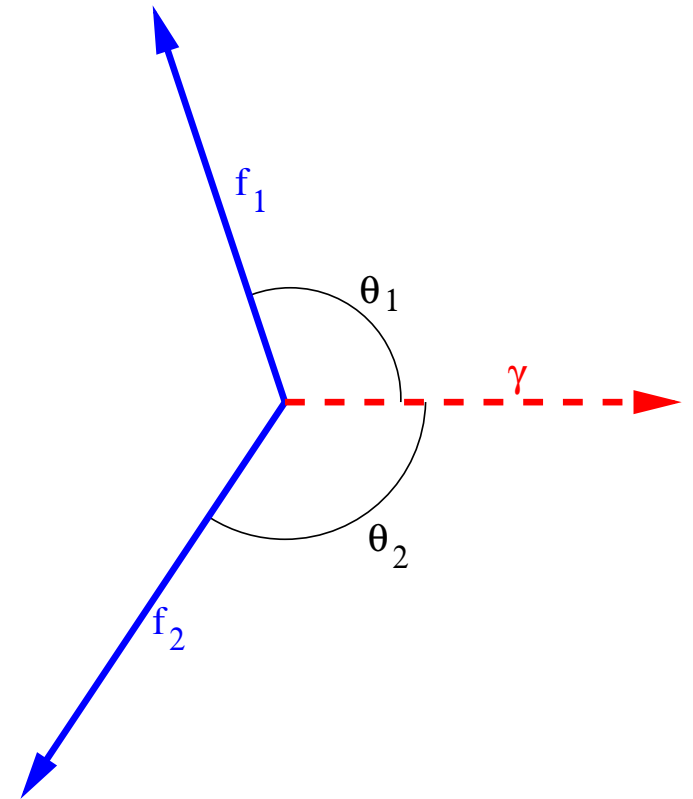
- Since the beam parameters will not be known to with high precision the spectrum of beamstrahlung has to be measured from data
- The energy loss of the outgoing beam is much larger than for the colliding particles
- For this reason the luminosity spectrum has to be measured from annihilation data
- Since one is interested in a $< 10^{-4}$ precision this cannot be done with calorimeters
- Method of choice: Bhabha acolinearity in the forward region:
 - very simple final state
 - very high cross section

The acolinearity method

- Assume only one photon is radiated
 - ⇒ $\sqrt{s'}$ can be calculated from fermion angles only

$$\frac{\sqrt{s'}}{\sqrt{s}} = \sqrt{\frac{\sin \theta_1 + \sin \theta_2 + \sin(\theta_1 + \theta_2)}{\sin \theta_1 + \sin \theta_2 - \sin(\theta_1 + \theta_2)}}$$

- The radiation in both directions can be unfolded in the fit
- This requires the knowledge of correlations
- **ISR/FSR has to be known from theory**



Requirements from theory

- Experimental θ resolution: $\delta\theta = 2 \cdot 10^{-5}$ resulting in an $\sqrt{s'}$ resolution of about 10^{-4}

- A possible photon energy cutoff in the generator therefore has to be below $k/E_b = 10^{-4}$

- This resolution is needed to disentangle beamstrahlung and beam energy spread

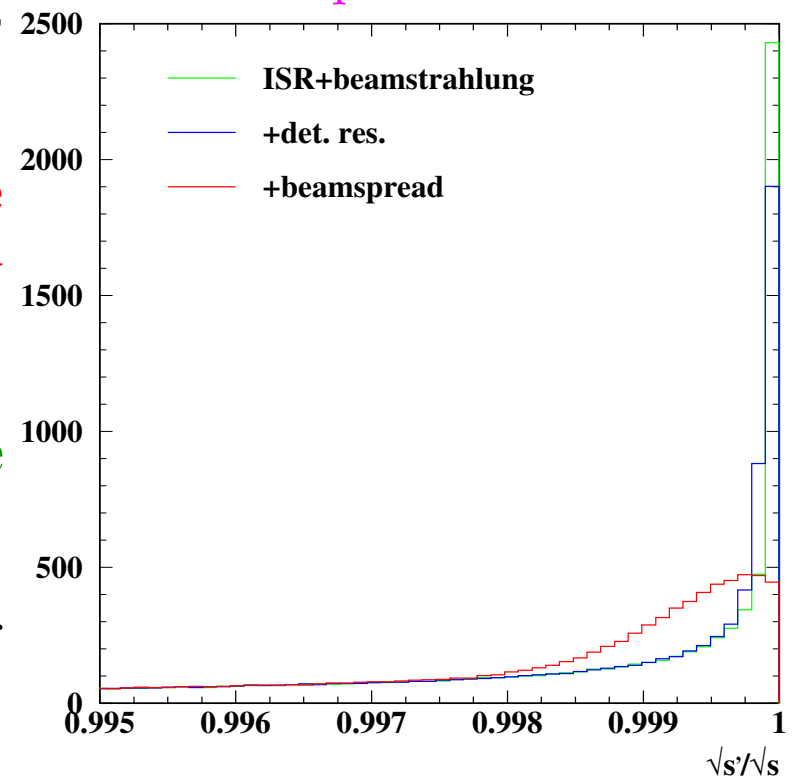
- Statistical error is $\Delta(\sqrt{s'}/\sqrt{s}) = 10^{-5}$ for 100 fb^{-1}

- The radiation has to be precise on this level

- The acolinearity method uses the charged tracks only

⇒ need to have an exact description of FSR and ISR/FSR interference

Reconstructed $\sqrt{s'}/\sqrt{s}$
spectra



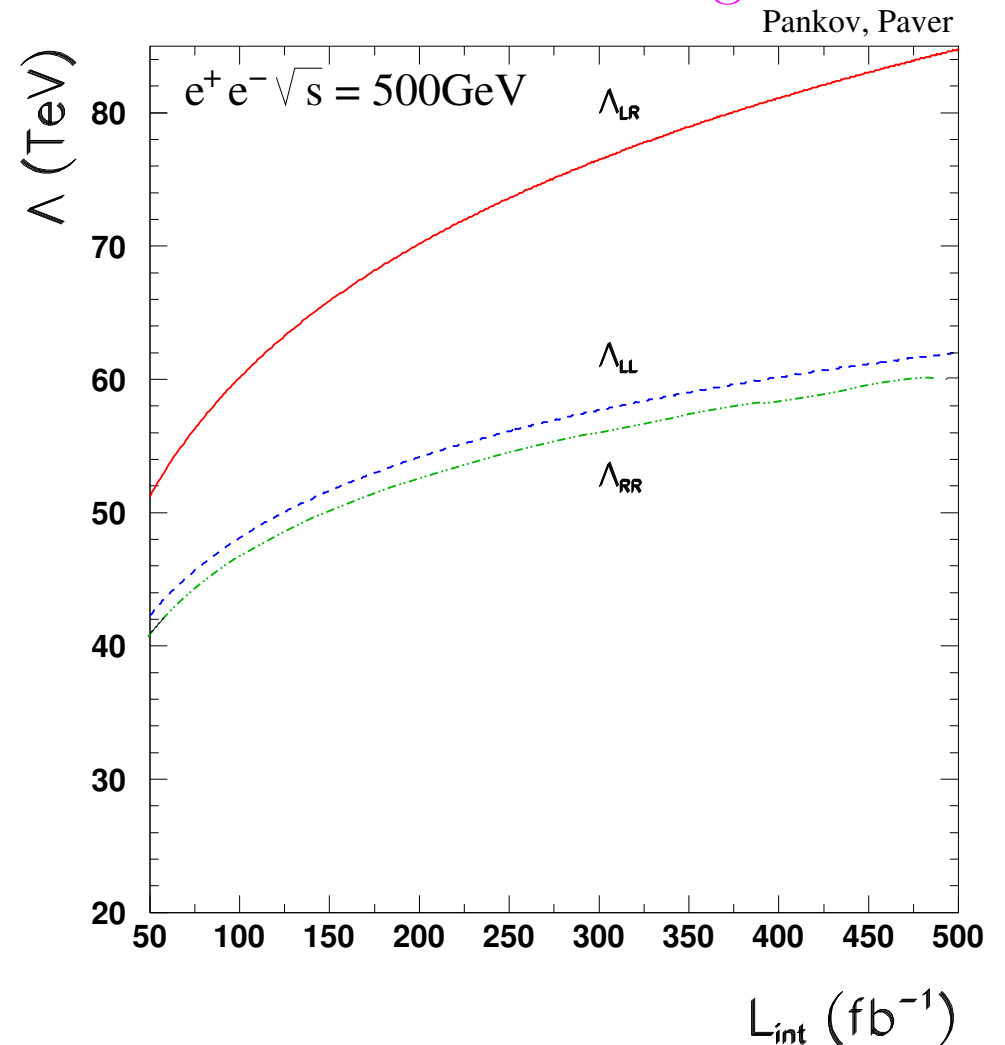
Physics with large angle Bhabhas at the ILC

- Large angle Bhabha scattering can be used as a general probe for new physics
- The most general description of this are contact interactions
- $\mathcal{O}(10^5)$ events per year are expected requiring the corresponding theoretical uncertainty
- Polarisation is very important to distinguish between the different helicity structures
- As a unique feature of Bhabha scattering the $J=0$ state can isolate the t-channel for vector currents and scalar s-channel exchange

Results

- Contact interaction limits of $\Lambda > 60 - 80 \text{ TeV}$ can be reached, depending on the helicity structure
($\sqrt{s} = 500 \text{ GeV}$, $\mathcal{L} = 500 \text{ fb}^{-1}$)
- These limits are better than and complementary to muons
- In principle similar limits can be reached for e^-e^- , do we have Moller scattering under control theoretically?

Contact interaction limits from Bhabha scattering



Conclusions

- Bhabha scattering is needed at the ILC for technical measurements (luminosity, luminosity spectrum) and for physics
- The required precision is up to 10^{-4}
- A few new technical requirements are needed to use the theory predictions at the ILC:
 - beamstrahlung
 - polarisation of both beams
 - non calorimetric measurements
- The 2-loop calculations are definitely a huge step towards these goals