

# First results with two light flavours of maximally twisted mass quarks

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LHP 2006, JLAB

# European Twisted Mass Collaboration

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# Outline

## Introduction

Twisted Mass Fermions

## Preparing the Ground

Quenched Experiments

Understanding the Phasestructure

Algorithmic Improvements

## Results

Setup

First Results

Scaling

## Conclusion

# Twisted Mass fermions

- ▶ the twisted mass Dirac operator:

$$D_{\text{tm}} = D_{\text{W}} + m_0 + i\mu\gamma_5\tau_3.$$

[Frezzotti, Grassi, Sint, Weisz, 1999]

- ▶ Wilson Dirac operator  $D_{\text{W}}$  with bare mass  $m_0$ .
- ▶ twisted mass parameter  $\mu$ .
- ▶  $\tau_3$  third Pauli matrix acting in flavour space
- ▶  $D_{\text{tm}}$  is protected against unphysically small eigenvalues

## Automatic $\mathcal{O}(a)$ improvement

If  $m_0$  is tuned to its critical value  $m_{\text{crit}}$  (maximal twist) then ...

- ▶ observables are automatically  $\mathcal{O}(a)$  improved.

[Frezzotti, Rossi, 2003]

Shown to work in practice for various observables in the quenched approximation [Jansen et al., 2004, 2005; Abdel-Rehim et al., 2004, 2005]

- ▶ Simplifies mixing during renormalisation
- ▶ Only one parameter ( $m_0$ ) must be tuned

but...

- ▶ parity and flavour symmetry explicitly broken

## Idea of the Proof

$$\langle O(\mathbf{x}) \rangle^{\text{lat}} = \langle O(\mathbf{x}) \rangle^c - a \int d\mathbf{y} \langle O(\mathbf{x}) \mathcal{L}_1(\mathbf{y}) \rangle^c + a \sum_k \langle O_k(\mathbf{x}) \rangle^c + \mathcal{O}(a^2)$$

- ▶ r.h.s.: all expt. values with cont. action: must obey symmetries of cont. action
- ▶ all operators in expansion must share lattice symmetries of  $O$
- ▶ example: cont. symmetry modified Parity

$$\tilde{\mathcal{P}} : \begin{cases} \psi(\vec{\mathbf{x}}, t) & \rightarrow \gamma_0 \exp(i\omega\gamma_5\tau_3)\psi(-\vec{\mathbf{x}}, t) \\ \bar{\psi}(\vec{\mathbf{x}}, t) & \rightarrow \bar{\psi}(-\vec{\mathbf{x}}, t) \exp(i\omega\gamma_5\tau_3)\gamma_0, \end{cases}$$

- ▶  $O$  must be even under  $\tilde{\mathcal{P}}$ , clover term is odd: term cancels in the expansion

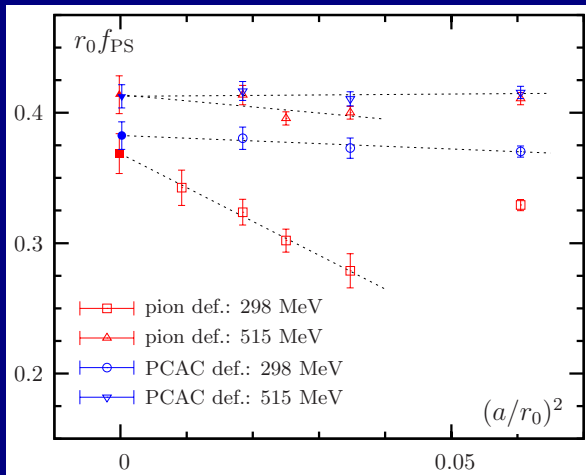
## A choice? Tuning to full twist

- ▶  $\mathcal{O}(a)$  ambiguity in  $m_{\text{crit}}$  does not harm automatic  $\mathcal{O}(a)$  improvement

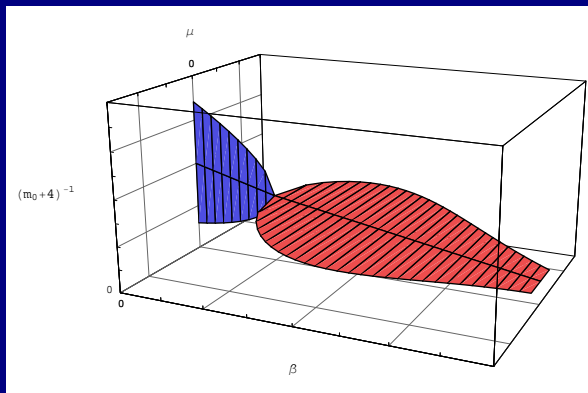
Several definitions available:

1.  $m_{\text{crit}}$  where  $m_{\text{PS}} = 0$  at  $\mu = 0$  (Pion def.)
  2.  $m_{\text{crit}}$  where  $m_{\text{PCAC}} = 0$  at  $\mu = 0$  (PCAC def.)
  3.  $m_{\text{crit}}$  where  $m_{\text{PCAC}} = 0$  for each value of  $a\mu$  separately
  4. ...
- ▶ Theoretical discussion in [Frezzotti, Martinelli, Papinutto, Rossi; Sharpe; Aoki, Bär]: do not use 1.

# Quenched Continuum Scaling of $f_{PS}$

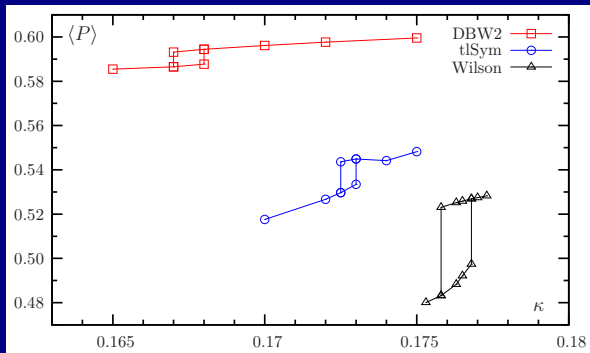


# Generic Phasestructure of Wilson Fermions



This is generic for Wilson type fermions!

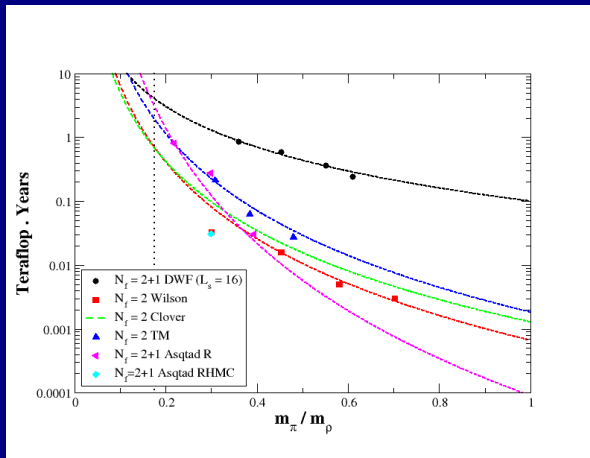
# Choosing the Gauge Action



- ▶ tISym is a theoretically sound compromise

# Speeding up the HMC

Resurrection of  
 Wilson fermions due  
 to algorithmic  
 improvements



# Setup

- ▶ We are using the tree-level Symanzik improved gauge action [Weisz, 1983]
- ▶  $N_f = 2$  mass-degenerate flavours of maximally twisted mass quarks
- ▶ Algorithm: HMC with multiple time scales and mass preconditioning [Urbach et al., 2005]
- ▶ Plan:
  - ▶ 3 lattice spacings: 0.075 – 0.125 fm
  - ▶ pseudo scalar masses in the range 250 – 550 MeV
  - ▶ volumes  $\geq 2$  fm
- ▶  $m_0$  tuned to  $m_{\text{crit}}$  at the lowest mass at each lattice spacing

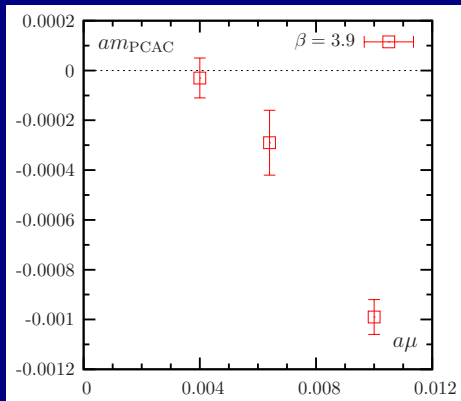
# Setup

$\beta$	$L^3 \times T$	$a$ [fm]	$m_{PS}$ [MeV]	$N_{\text{therm}}$	$N_{\text{traj}}$
3.9	$24^3 \times 24$	$\approx 0.095$	$\approx 280$	1500	5000
			$\approx 350$	1500	5000
			$\approx 430$	1500	5000
			$\approx 510$	1500	5000
4.05	$32^3 \times 64$	$\approx 0.075$	$\approx 280$	1500	1200
			$\approx 350$	1500	500

$\beta = 4.05$  very preliminary!

## Tuning to full twist

- ▶  $m_{\text{PCAC}} = 0$  at lowest  $a\mu$
- ▶  $m_{\text{PCAC}} = 0$  dependence visible
- ▶ no news: known from quenched



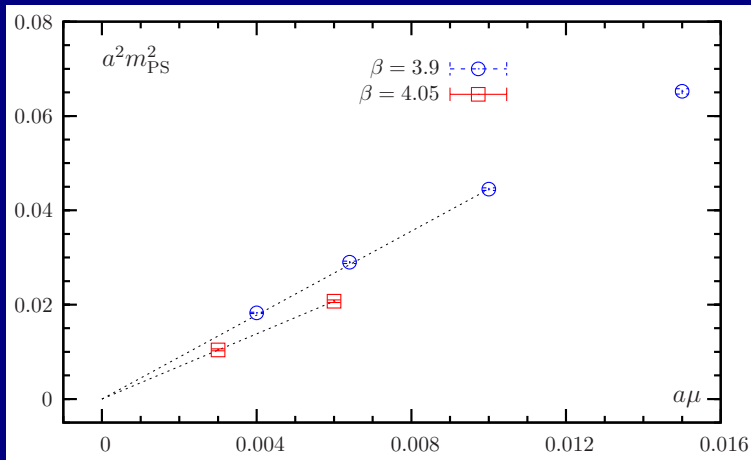
## Some simple observables

- ▶ First observables to look at:  $m_{\text{PS}}$  and  $f_{\text{PS}}$
- ▶  $f_{\text{PS}}$  at maximal twist can be obtained from

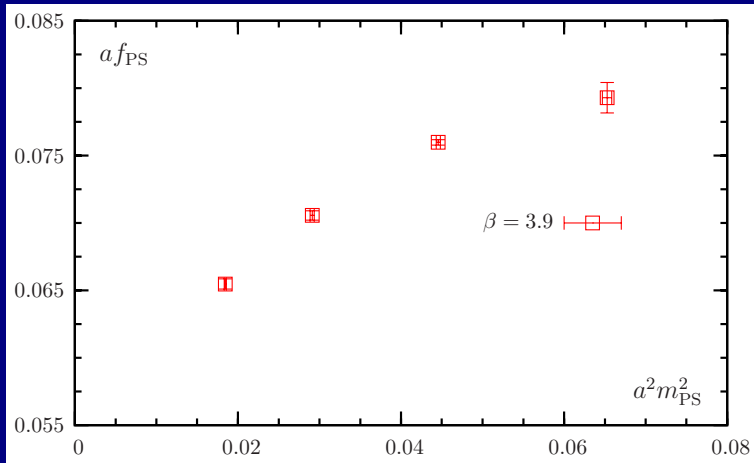
$$f_{\text{PS}} = \frac{2\mu}{m_{\text{PS}}^2} |\langle 0 | P^1(0) | \pi \rangle|$$

- ▶ Note that at maximal twist  $f_{\text{PS}}$  does not need to be renormalised
- ▶ We estimate finite size (FS) effects with NLO ChPT formula from Gasser and Leutwyler [Gasser, Leutwyler, 1987]
- ▶ Checked against resummed Lüscher formula [Lüscher, 1986; Colangelo, Dürr, Haefeli, 2005]
- ▶ We use spin diluted random time slice sources, fuzzing [Michael] and variational methods [Michael, 1985 ;Lüscher, Wolff, 1990]

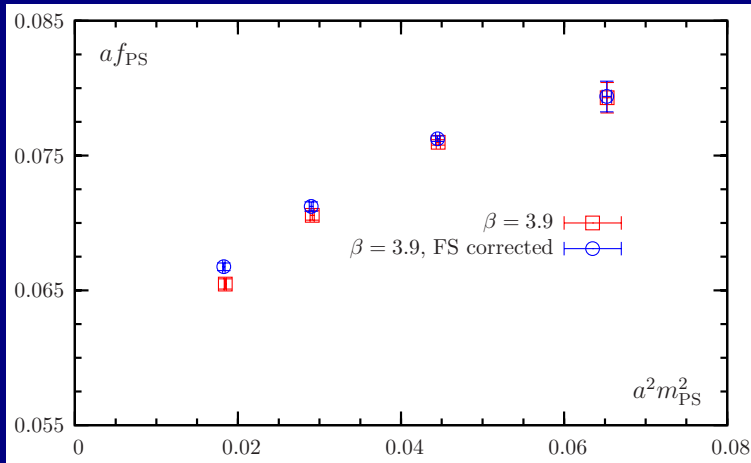
# $am_{\text{PS}}$ versus $a\mu$ at $\beta = 3.9$



# $f_{\text{PS}}$ at $\beta = 3.9$



# $f_{\text{PS}}$ at $\beta = 3.9$



## $f_{\text{PS}}$ at $\beta = 3.9$

- ▶ Fits to the data with ChPT formulae in progress
- ▶ when the data is extrapolated linearly to the physical point we obtain:

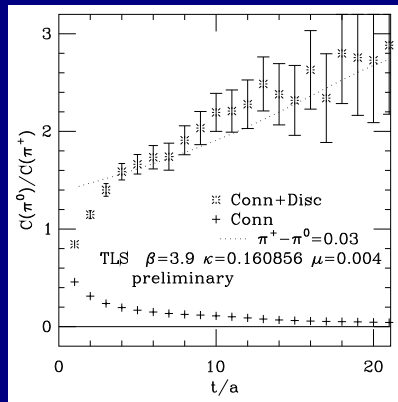
$$f_{\pi} = 126.3 \pm 0.8 \pm 0.7 \text{ MeV}$$

- ▶ First error comes from  $m_{\text{PS}}$ ,  $f_{\text{PS}}$  and extrapolation, the second from  $r_0/a$
- ▶  $r_0 = 0.5 \text{ fm}$  was used
- ▶ In our normalisation  $f_{\text{PS}} = 131 \text{ MeV}$

## Effects of Isospin breaking

Flavour symmetry explicitly broken by twisted mass term at finite  $a$

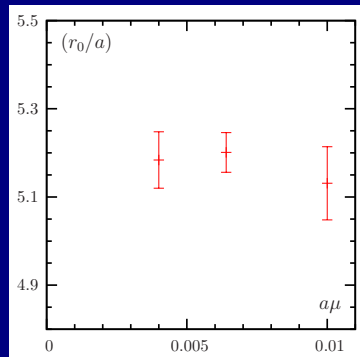
- ▶ Expected to be largest in  $m_{\text{PS}}^{\pm} - m_{\text{PS}}^0$
- ▶  $\beta = 3.9$ ,  $a\mu = 0.004$ ,  $am_{\text{PS}} = 0.1358(5)$ :  
 $m_{\text{PS}}^{\pm} - m_{\text{PS}}^0 = 0.03(1)$
- ▶ A factor of 2 smaller than quenched
- ▶ Splitting in the vector mass compatible with 0

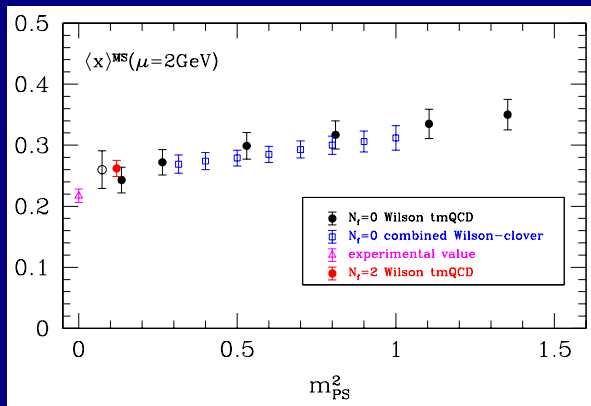


# Setting the scale

- ▶ Using the Sommer scale  $r_0$
- ▶ Value of  $r_0$  not very well known
- ▶ Scale setting with f.i.  $f_K$ ,  $m_{K^*}$  etc. in progress
- ▶ Here I use  $r_0 = 0.5$  fm
- ▶  $\mu$  dependence seems to be weak

$$\beta = 3.9$$

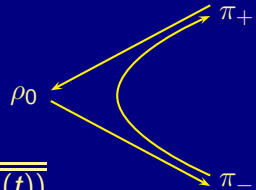


$\langle x \rangle$  for the Pion

Lowest moment of Non-Singlet Pion Parton Distribution  
Function  $\langle x \rangle$

## Exploring $\rho$ -decay [McNeile, Michael, 2003]

- ▶  $\rho$  does not (yet) decay at  $\beta = 3.9$ :  
 $a\Delta m \approx 0.15$  at  $a\mu = 0.004$
- ▶ uncharged vector meson at rest:



$$R(t) = \frac{\rho_0(0) \rightarrow \pi_+\pi_-(t)}{\sqrt{(\rho_0(0) \rightarrow \rho_0(0))(\pi_+\pi_-(0) \rightarrow \pi_+\pi_-(t))}}$$

- ▶ Assuming  $m_\rho \approx m_{\pi\pi}$ ,  $\mathbf{x} = \langle \rho | \pi\pi \rangle$ :

$$R \propto \sum_{t=0}^T \mathbf{x} e^{-m_\rho t} e^{-m_{\pi\pi}(T-t)} = e^{-mT} \cdot \mathbf{x} \cdot t$$

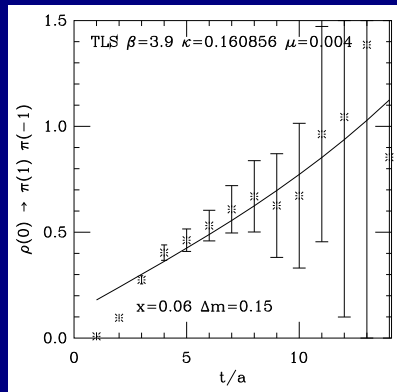
for  $xt \ll 1$

## Exploring $\rho$ -decay

- ▶ We extract  $\bar{g} = 1.2(2)$ , with

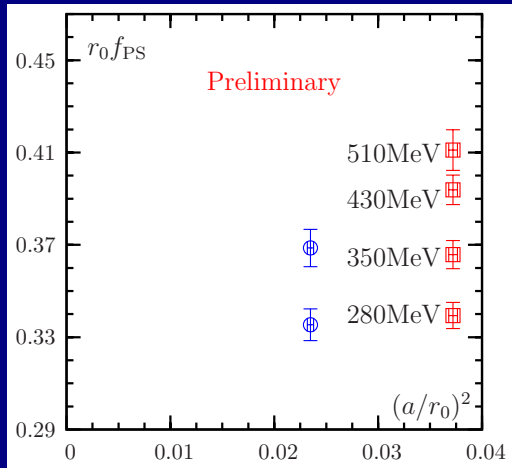
$$\bar{g} = \Gamma m_\rho E_{\pi\pi} / k^3$$

- ▶ experimental number:  $\bar{g} = 1.39$
- ▶ Already well compatible
- ▶ Larger lattices needed for a decaying  $\rho$ .



# Preliminary continuum scaling

- ▶  $r_0/a$  value at the lowest  $\mu$  value
- ▶ masses not yet exactly matched



# Conclusion

- ▶ mtmQCD stands on a sound basis
- ▶ First encouraging results with  $N_f = 2$  flavours of maximally twisted mass quarks
- ▶ We can reach values for  $m_{PS}$  as low as 280 MeV
- ▶ Flavour symmetry breaking effects are visible, but significantly smaller than quenched
- ▶ Lattice artifacts in  $f_{PS}$  seem to be small
- ▶ First physics applications in progress
- ▶  $N_f = 2 + 1 + 1$  in progress: talk of Enno