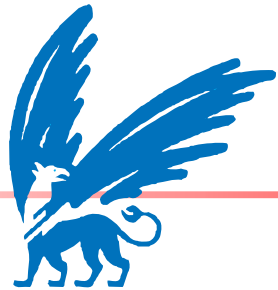
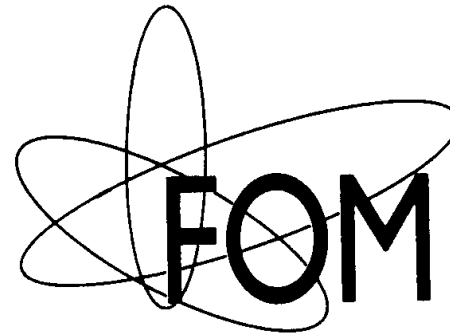


FRAGMENTATION INTO SPIN-1 HADRONS

A. Bacchetta



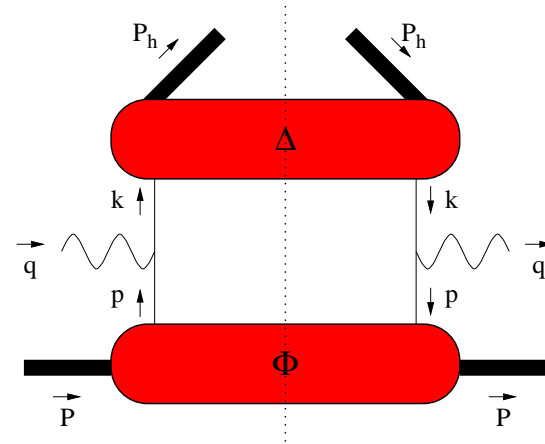
vrije Universiteit *amsterdam*



- Semi-inclusive DIS
- Fragmentation to spin-1 hadrons
- Transversity measurements
- Positivity bounds and modeling attempts

Semi-inclusive DIS

$$d\sigma(l + H \rightarrow l' + h + X) \propto L_{\mu\nu} W^{\mu\nu}$$



$$2MW^{\mu\nu} \propto \text{Tr} [\Phi(x_B) \gamma^\mu \Delta(z_h) \gamma^\nu]$$

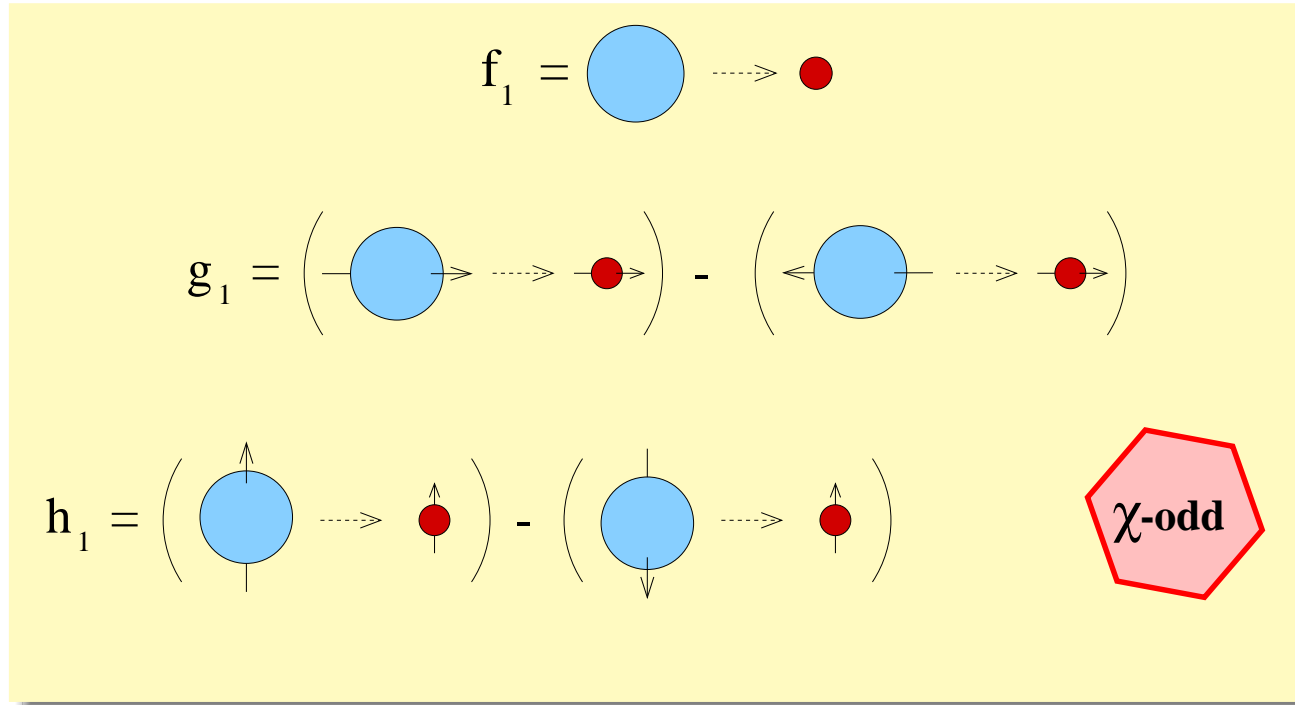
$$x_B = \frac{Q^2}{2P \cdot q} \quad z_h = \frac{2P_h \cdot q}{Q^2}$$

$$\Phi(x) = \frac{1}{2} \int d^2p_T dp^- \Phi(p, P) \Big|_{p^+ = xP^+}$$

$$\Delta(z) = \frac{z}{4} \int d^2k_T dk^+ \Delta(k, P_h) \Big|_{k^- = \frac{P_h^-}{z}}$$

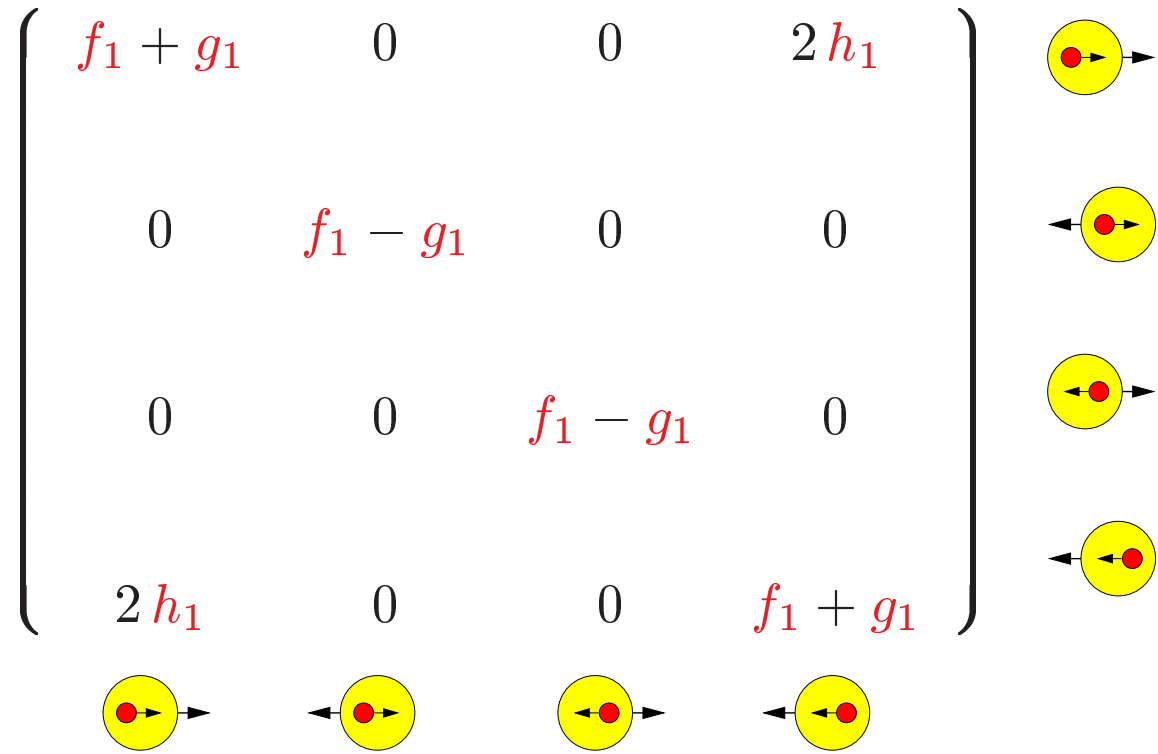
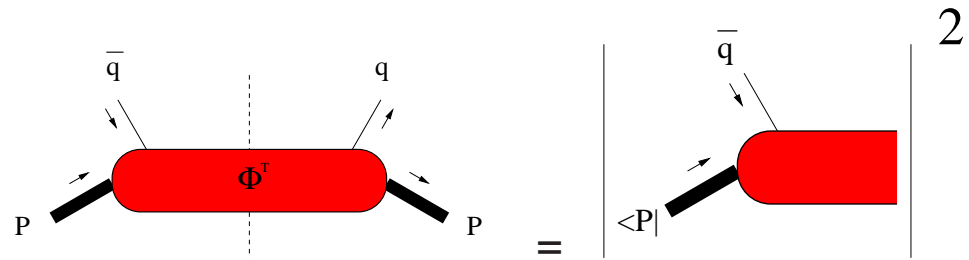
Distribution functions for spin-1/2

From decomposition of Φ correlation function.



Quark+target spin structure

Full scattering matrix $M = (\Phi \gamma^+)^T$
 for $\bar{q} + P \rightarrow X$
 in parton \otimes hadron helicity spaces



Positive definiteness of this matrix requires:

Positivity of diagonal elements \Rightarrow trivial bounds

$$f_1(x) \geq 0$$

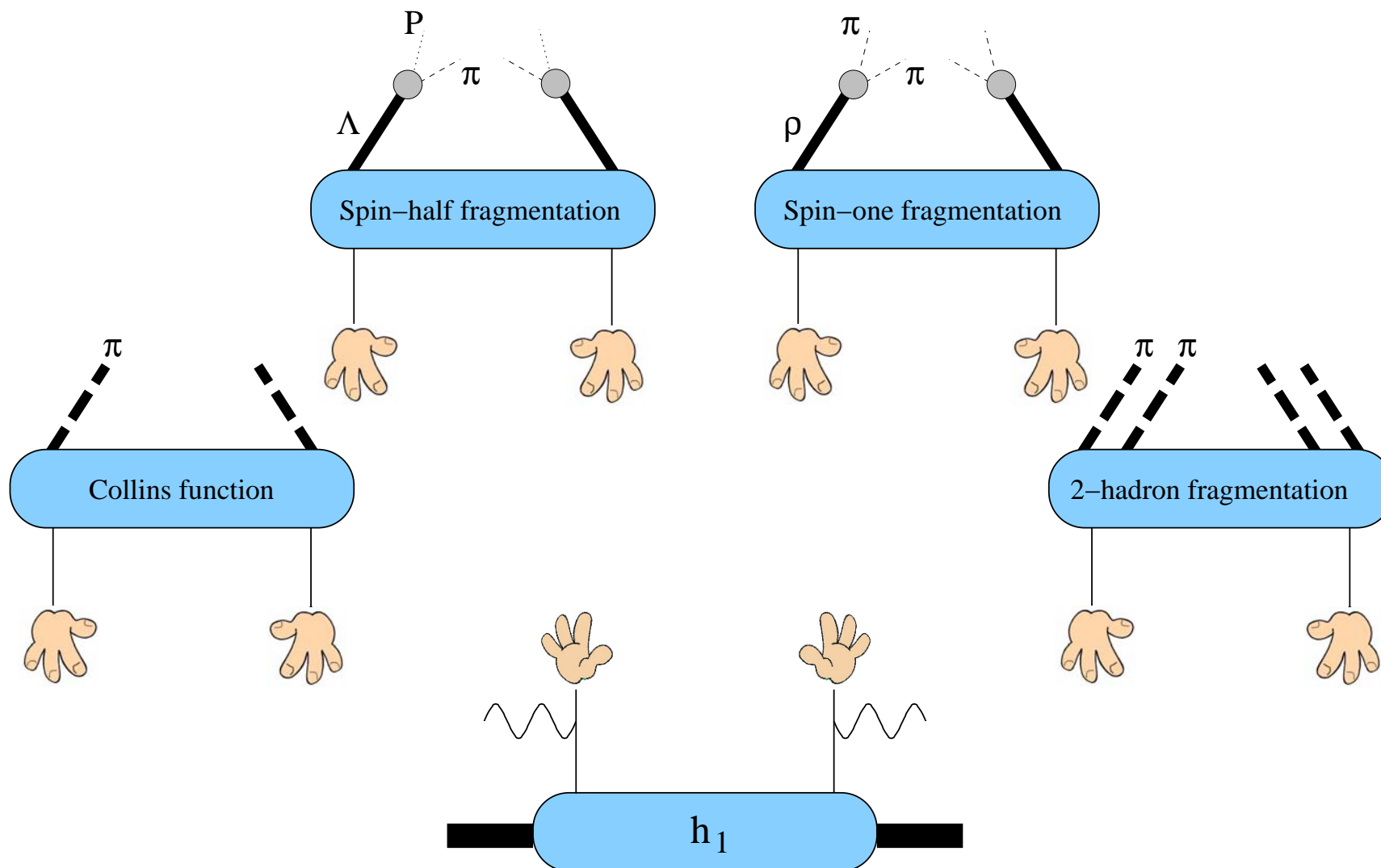
$$|g_1(x)| \leq f_1(x)$$

Positivity of eigenvalues \Rightarrow less trivial (Soffer) bound

$$|h_1(x)| \leq \frac{1}{2} (f_1(x) + g_1(x))$$

Chiral-odd partners for the transversity distribution

ZEUTHEN 2001



Fragmentation to spin-1 hadrons

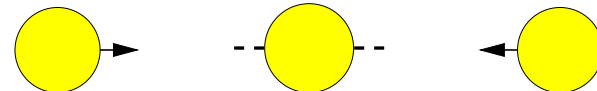
We need to measure the polarization of the outgoing hadron by means of **VECTOR AND TENSOR ANALYZING POWERS** of the decay

$$\Delta_{ij}(k, P_h, A_h^i, A_h^{ij}) = \int \frac{d^4\xi}{(2\pi)^4} e^{-ik \cdot \xi} \langle 0 | \psi_i(0) | P_h, A_h^i, A_h^{ij} \rangle \langle P_h, A_h^i, A_h^{ij} | \bar{\psi}_j(\xi) | 0 \rangle$$

or as a matrix in the hadron's helicity space

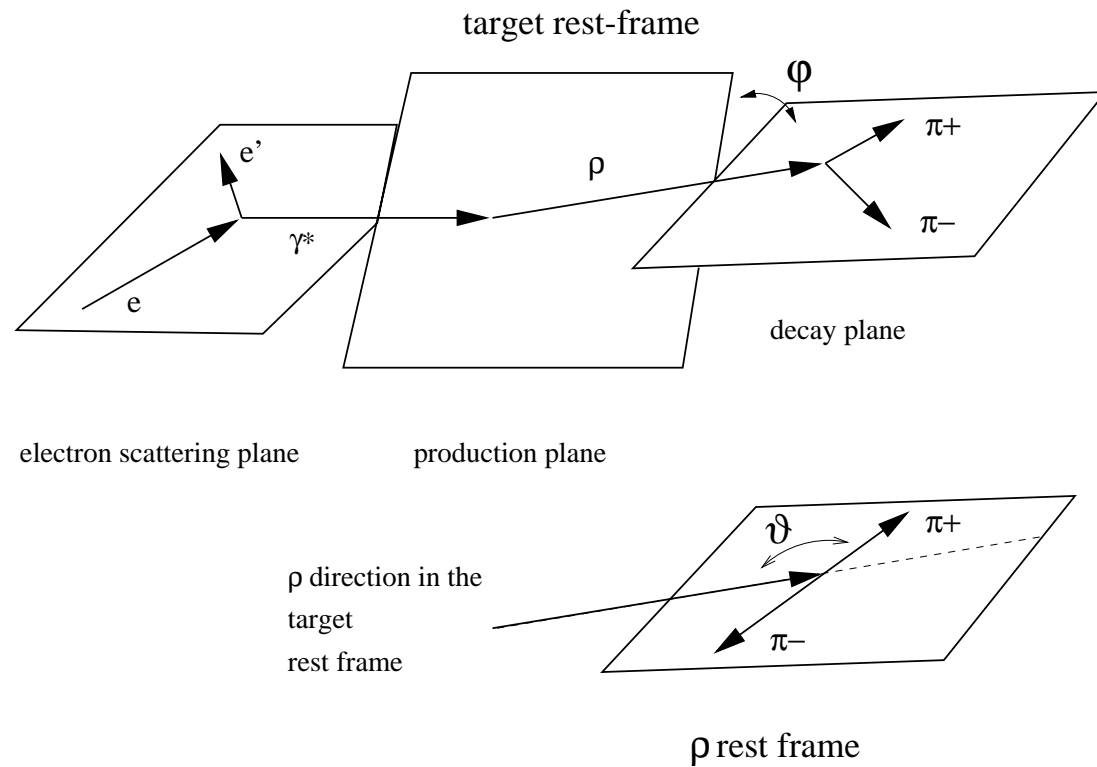
$$\Delta_{ij; \Lambda \Lambda'}(k, P_h) = \int \frac{d^4\xi}{(2\pi)^4} e^{-ik \cdot \xi} \langle 0 | \psi_i(0) | P_h, \Lambda \rangle \langle P_h, \Lambda' | \bar{\psi}_j(\xi) | 0 \rangle$$

where Λ and Λ' are hadron's helicities (1,0,-1)



$$\Delta(k, P_h, A_h^i, A_h^{ij}) = \text{Tr} \left[R_h^{(\text{decay})} (A_h^i, A_h^{ij}) \Delta(k, P_h) \right]$$

Analyzing powers



$$A_{LL} = \frac{1}{3} (\cos^2 \theta + \cos 2\theta)$$

$$A_{LT}^x = -\sin 2\theta \cos \varphi$$

$$A_{TT}^{xx} = -\sin^2 \theta \cos 2\varphi$$

$$A_{LT}^y = -\sin 2\theta \sin \varphi$$

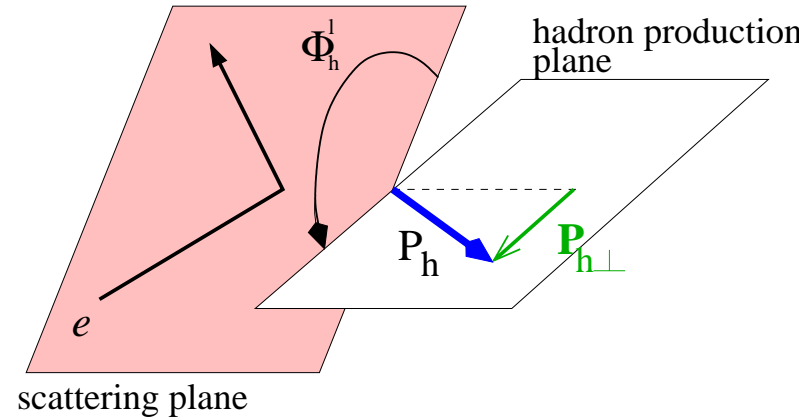
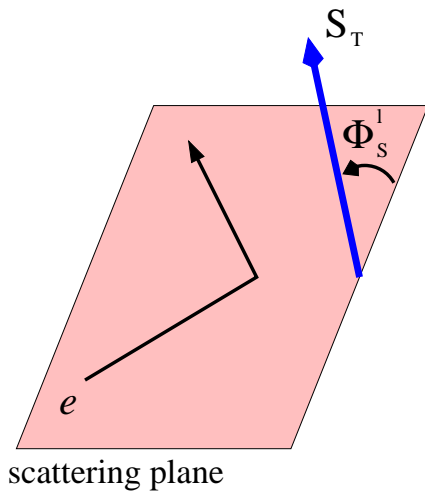
$$A_{TT}^{xy} = -\sin^2 \theta \sin 2\varphi$$

Transversity and asymmetries

$$d\sigma^\uparrow - d\sigma^\downarrow = f(\phi_S^l, |P_{h\perp}|, \phi_h^l, x_B, z_h)$$

$$\left\langle \begin{array}{c} \text{Weight} \\ \text{Function} \end{array} \right\rangle_T(x_B, z_h) = \int d\phi_S^l d|P_{h\perp}| d\phi_h^l \text{Weight Function}(d\sigma^\uparrow - d\sigma^\downarrow)$$

$$\propto h_1(x_B) \times \text{?}(z_h)$$

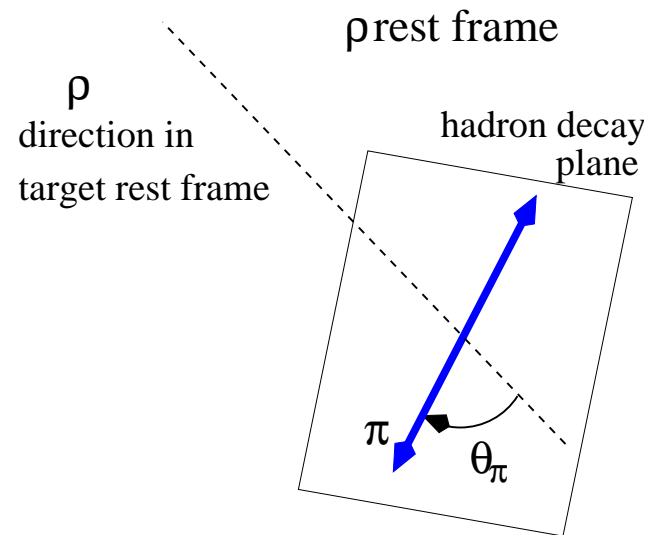
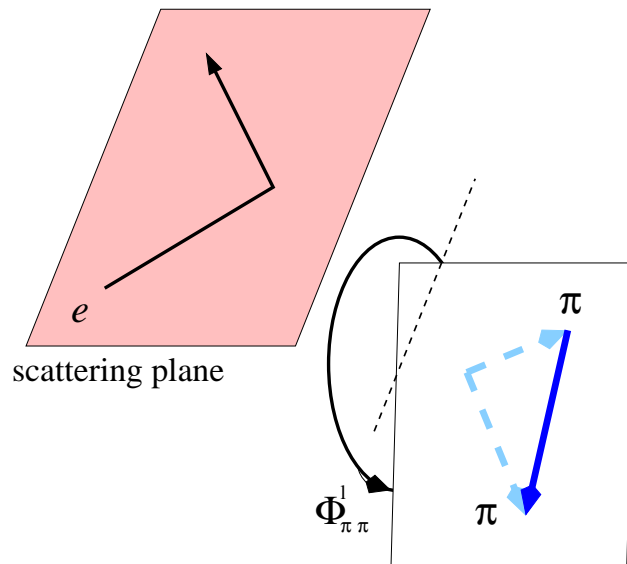


Asymmetries in spin-1 production

$$d\sigma^\uparrow - d\sigma^\downarrow = f(\theta_\pi, \phi_{\pi\pi}^\ell, \phi_S^\ell, |P_{\rho\perp}|, \phi_\rho^\ell, x_B, z_\rho)$$

$$\left\langle \begin{array}{c} \text{Weight} \\ \text{Function} \end{array} \right\rangle_T(x_B, z_\rho) = \int d\theta_\pi d\phi_{\pi\pi}^\ell d\phi_S^\ell d|P_{\rho\perp}| d\phi_\rho^\ell \text{Weight Function}(d\sigma^\uparrow - d\sigma^\downarrow)$$

$$\propto h_1(x_B) \times \text{?}(z_\rho)$$



Asymmetries containing the transversity distribution

$$\langle \sin(\phi_{\pi\pi}^\ell + \phi_S^\ell) \rangle_T \propto h_1(x_B) H_{1LT}(z_\rho) \begin{matrix} \chi\text{-odd} \\ T\text{-odd} \end{matrix}$$

- Requires measurement of $\phi_{\pi\pi}^\ell, \phi_S^\ell, x_B, z_\rho$
- No complications due to presence of hadron transverse momentum

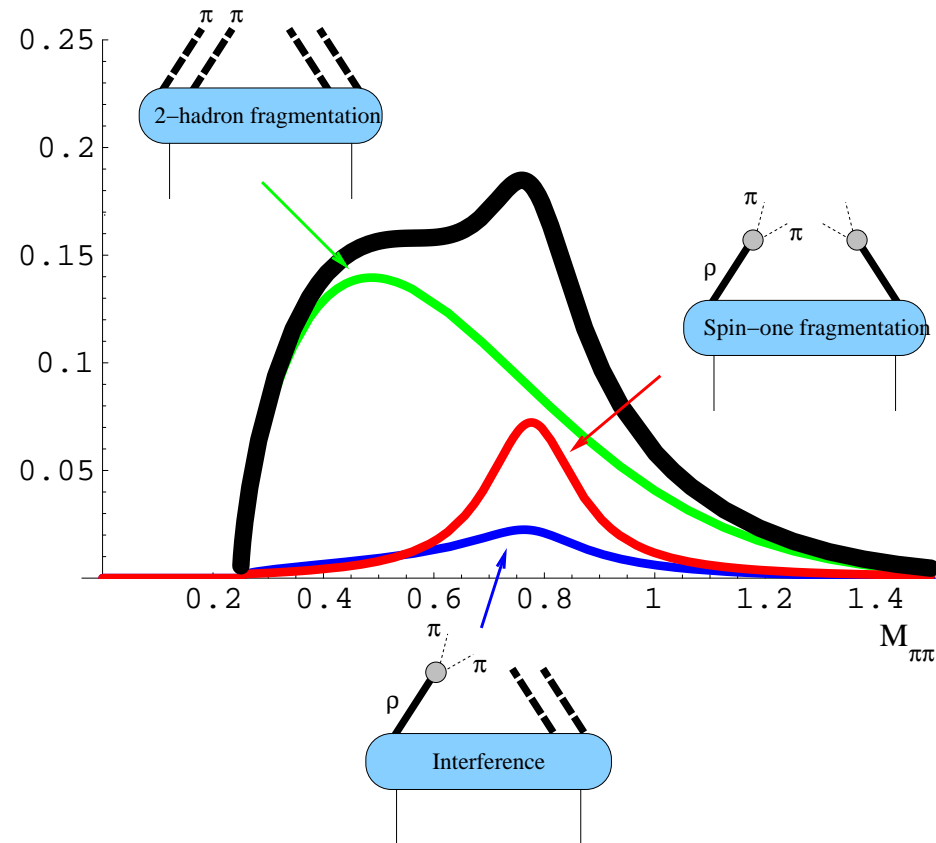
$$\left\langle \frac{|P_{\rho\perp}|}{2z M_h} \sin(2\phi_{\pi\pi}^\ell + \phi_S^\ell - \phi_\rho^\ell) \right\rangle_T \propto h_1(x_B) H_{1TT}^{(1)}(z_\rho)$$

$$\left\langle \frac{|P_{\rho\perp}|}{2z M_h} \sin(\phi_\rho^\ell + \phi_S^\ell) \right\rangle_T \propto h_1(x_B) H_{1LL}^{\perp(1)}(z_\rho)$$

$$\left\langle \frac{|P_{\rho\perp}|^2}{2z^2 M_h^2} \sin(\phi_{\pi\pi}^\ell - \phi_S^\ell - 2\phi_\rho^\ell) \right\rangle_T \propto h_1(x_B) H_{1LT}^{\perp(2)}(z_\rho)$$

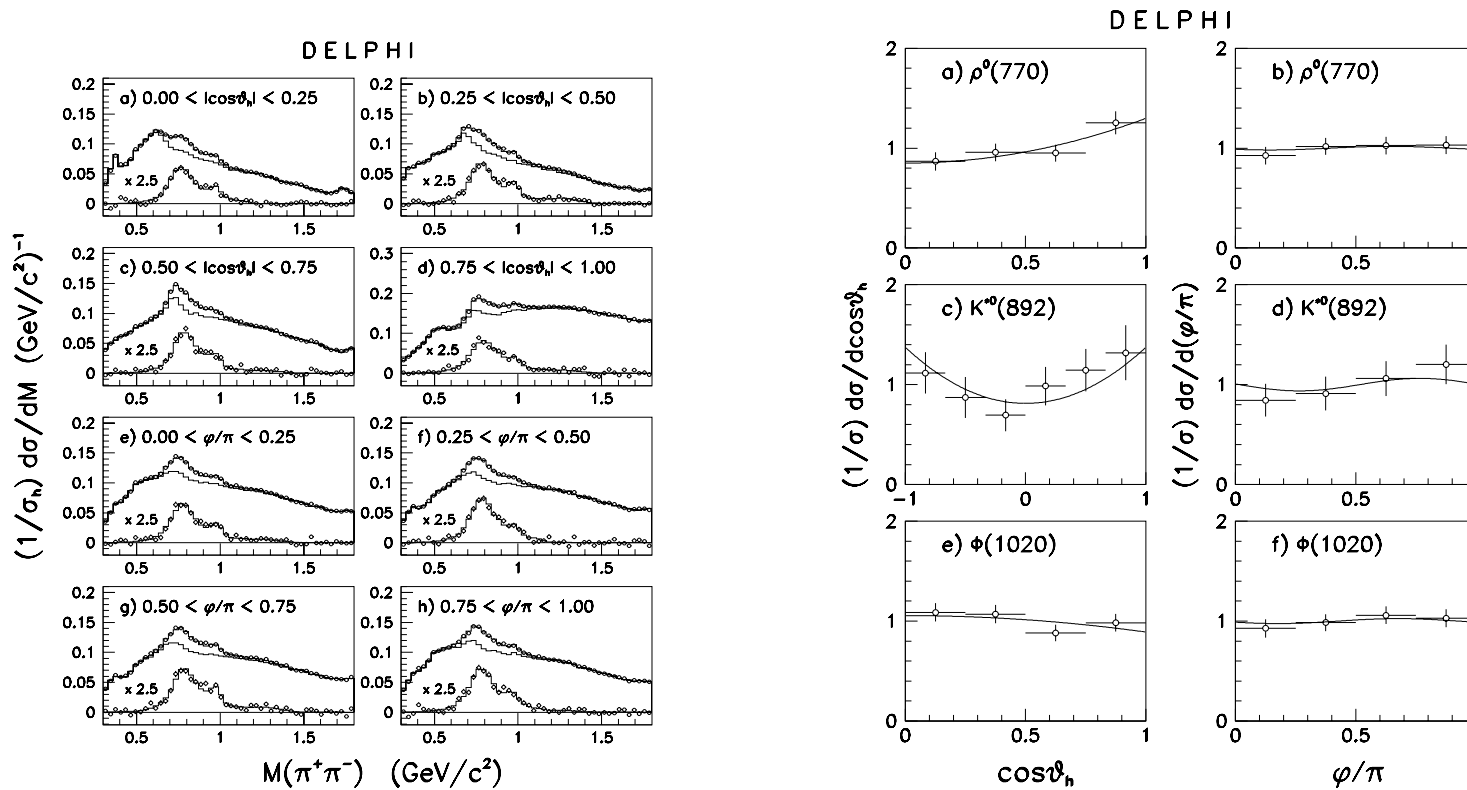
Where to observe these functions?

Mass spectrum of pion couples



Measurements in $e^+ e^-$ annihilation

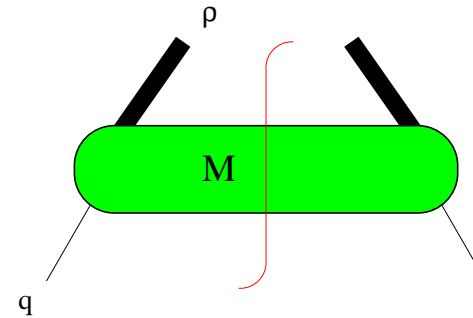
Delphi collaboration Phys. Lett. B 406 (1997) 271



Measurement of decay distribution seems to be feasible.

Quark+hadron spin structure

Full scattering matrix M for $q \rightarrow \rho + X$



$D_1 + G_1 - \frac{D_{1LL}}{3}$	0	0	0	$\sqrt{2}(H_1 + iH_{1LT})$	0	
0	$D_1 + \frac{2D_{1LL}}{3}$	0	0	0	$\sqrt{2}(H_1 - iH_{1LT})$	
0	0	$D_1 - G_1 - \frac{D_{1LL}}{3}$	0	0	0	
0	0	0	$D_1 - G_1 - \frac{D_{1LL}}{3}$	0	0	
$\sqrt{2}(H_1 - iH_{1LT})$	0	0	0	$D_1 + \frac{2D_{1LL}}{3}$	0	
0	$\sqrt{2}(H_1 + iH_{1LT})$	0	0	0	$D_1 + G_1 - \frac{D_{1LL}}{3}$	

Positive definiteness of this matrix requires:

A.B., P.J. Mulders, hep-ph/0104176

Positivity of diagonal elements \Rightarrow trivial bounds

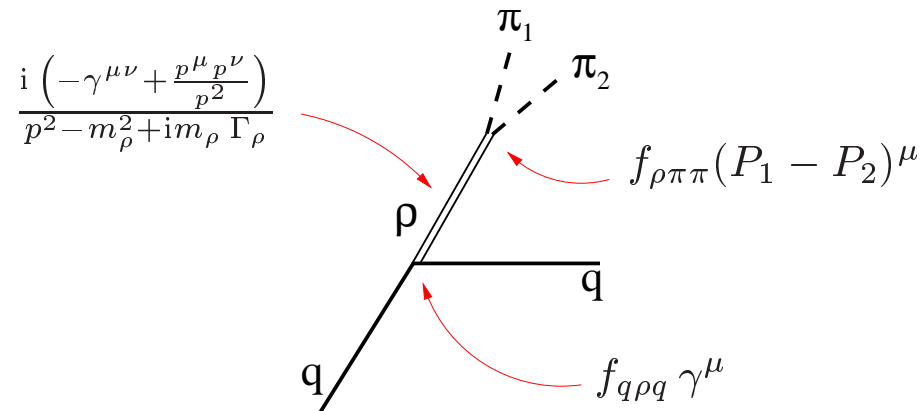
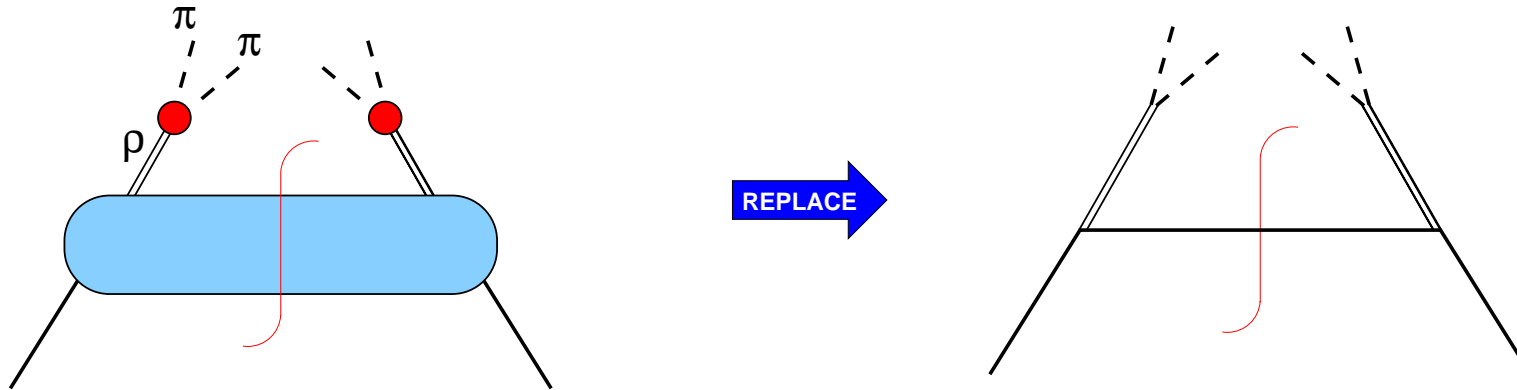
$$\begin{aligned}
 D_1(z) &\geq 0 \\
 -\frac{3}{2} D_1(z) &\leq D_{1LL}(z) \leq 3 D_1(z) \\
 |G_1(z)| &\leq D_1(z) - \frac{1}{3} D_{1LL}(z) \leq \frac{3}{2} D_1(z)
 \end{aligned}$$

Positivity of 2-dimensional minors \Rightarrow (Soffer-like) bound

$$[H_{1LT}(z)]^2 \leq \left[D_1(z) + \frac{2}{3} D_{1LL}(z) \right] \left[D_1(z) - \frac{1}{3} D_{1LL}(z) \right] \leq \frac{9}{2} [D_1(z)]^2.$$

Generating non-zero T-odd functions

Take a **simple model** for the fragmentation process



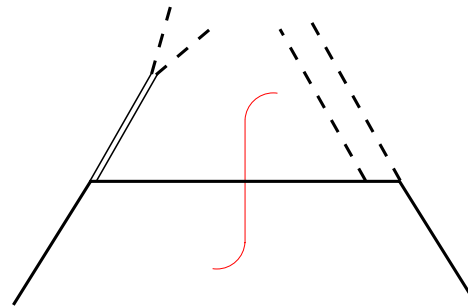
Generating non-zero T-odd functions

Add **one-loop corrections** (see also talk by Rajen Kundu)



T-odd functions $\neq 0$

“Competing” or “complementary” mechanism: **interference fragmentation functions** (see talk by Marco Radici)



Conclusions

➡ Fragmentation to spin-1 hadron is a viable option for transversity measurements, at the same level as Collins function and interference fragmentation functions.

➡ Since polarimetry on the final hadron requires the observation of two decay products, it represents a specific contribution to the more general case of two hadron production.

➡ Information on the unknown fragmentation functions can be extracted from positivity bounds and simple models.