

### 3 Top-quark-physics

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#### Introduction

- The top quark is the heaviest fermion  
( $m_t \approx 175 \text{ GeV} \sim v$ )
- In the SM it is just the isospin partner of the b-quark
- however in some models it plays a special role in electroweak symmetry breaking
- ▶ it is very important to study the top properties
- since  $m_t^2 \gg m_W^2$  the top width is very large

$$\Gamma_t \approx \frac{G_F m_t^3}{8\sqrt{2}\pi} \approx 1.7 \text{ GeV} \gg \lambda_{QCD} \Rightarrow$$

- there exist no toponium resonances at threshold
- the top decays before it fragments  
⇒ the top-polarization gets preserved to the decay, like the  $\tau$  at LEP

## LHC:

- $\sigma(t\bar{t}) \sim 1 \text{ nb} \Rightarrow$  huge data samples should allow very precise studies of top decays, especially rare decays
- $t\bar{t}$  production by strong interaction  $\Rightarrow$  no interesting information on  $t\bar{t}Z$  coupling
- $t\bar{t}$  production in continuum  $\Rightarrow$  no threshold scan possible

## LC:

- $\sigma(t\bar{t}) \sim 1 \text{ pb} \Rightarrow \sim 10^5$   $t\bar{t}$  events allow for precise studies
- $t\bar{t}$  production via  $\gamma, Z$  exchange  $\Rightarrow$   $t\bar{t}Z$  couplings can be measured
- $\sqrt{s}$  can be adjusted at will  $\Rightarrow$  threshold scan possible

## Measurement of the top-mass

Why do we want to know the top-mass as accurate as possible?

- a future theory of flavor hopefully predicts fermion masses or mass/ratios  
→ The mass of the heaviest quark should be known as close as possible to the precision of the  $\tau$ -mass
- in precision tests of the SM  $m_t$  enters quadratically:
  - $\Delta m_W / \Delta m_t = 0.006$   
ultimately:  $\Delta m_W = 6 \text{ MeV}$
  - $\Delta \sin^2 \theta_{\text{eff}}^{\ell} / \Delta m_t = 0.00003 / \text{GeV}$   
ultimately:  $\Delta \sin^2 \theta_{\text{eff}}^{\ell} = 0.00002$ $\Rightarrow$  need  $\Delta m_t < 1 \text{ GeV}$
- In SUSY models radiative corrections to light Higgs (h) mass:  $\Delta m_h / \Delta m_t \approx 1$   
 $\Rightarrow$  aim for  $\Delta m_t \approx \Delta m_h \approx 50 \text{ MeV}$

## What is the top mass?

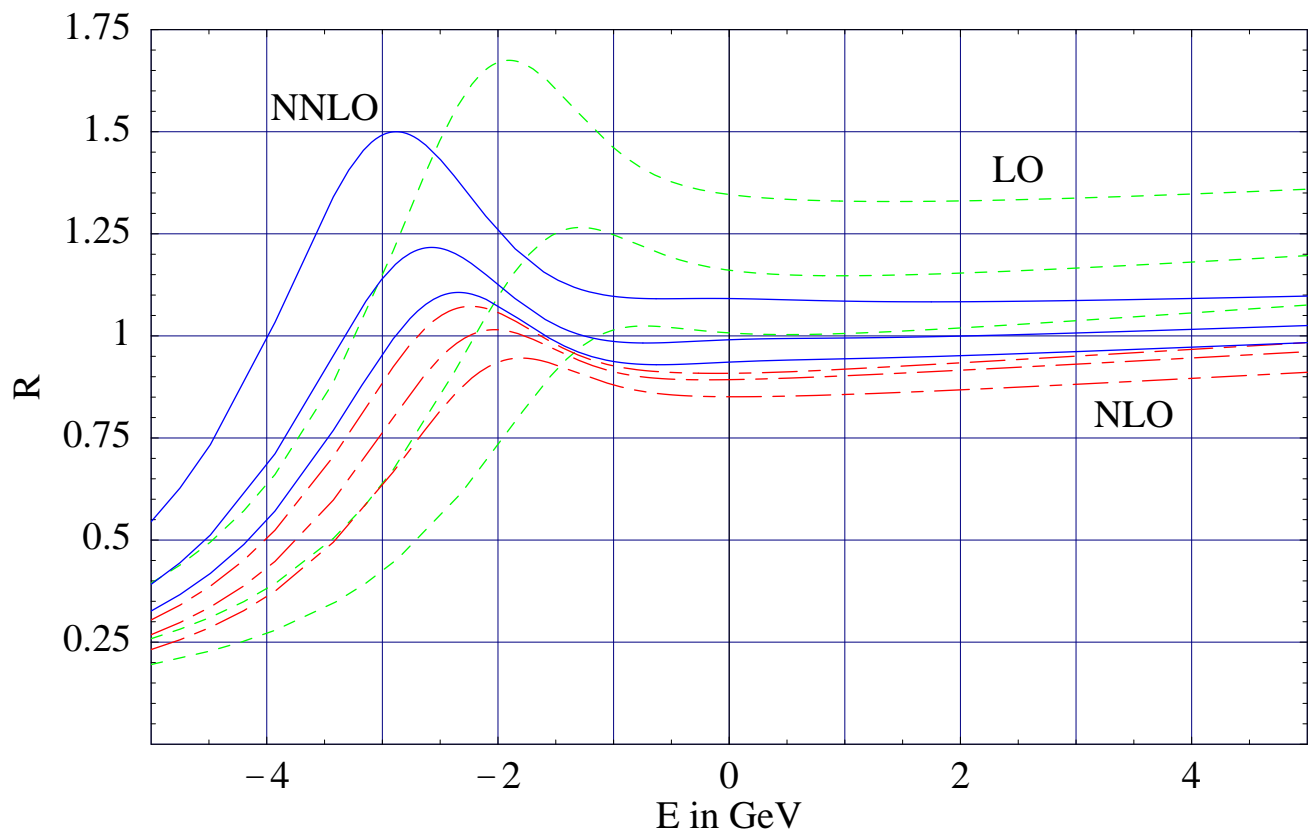
- Quarks are not free particles  $\Rightarrow$  their mass is not unambiguously defined
- pole mass: pole of propagator, natural definition if top-decays are reconstructed
- $\overline{\text{MS}}$ -mass: running mass in QCD (like coupling constant), needed in radiative corrections
- conversion pole mass  $\rightarrow$   $\overline{\text{MS}}$ -mass has theoretical uncertainties of  $\mathcal{O}(1 \text{ GeV})$
- limit of all top-reconstruction methods
- additional ambiguity of same order for reconstruction methods since only color neutral objects can be reconstructed
- threshold scans: no natural mass definition, can do calculations in several ones

## Most promising method: threshold scan

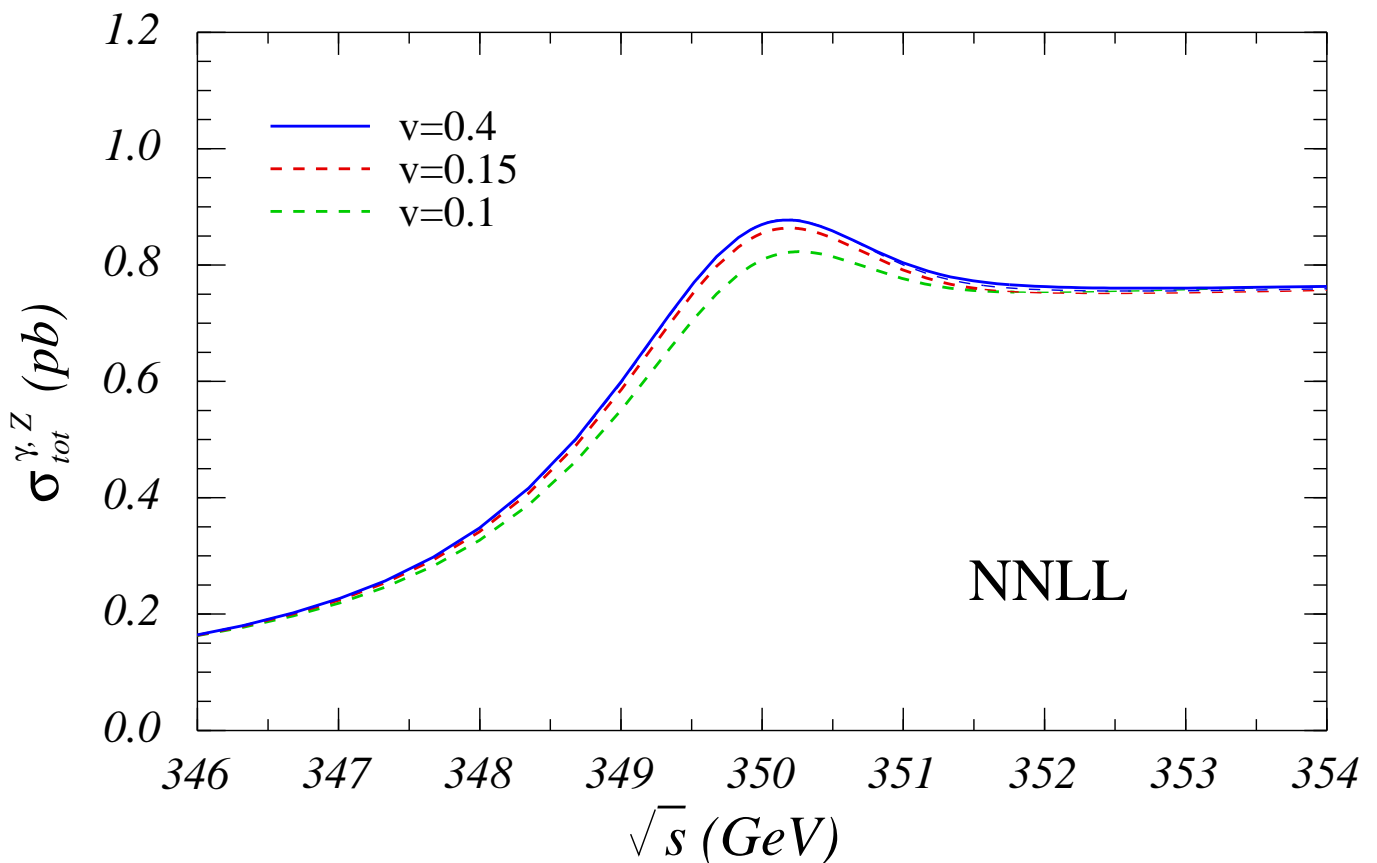
- Due to large top mass and the corresponding short lifetime no toponium resonances are existing any more
- However still large corrections due to Coulomb-like QCD potential:

$$E_{\text{tot}}(r) = 2m_t + V(r)$$
$$V(r) \propto \frac{\alpha_s(1/r)}{r}$$

- QCD corrections known to 3rd order (pole mass)



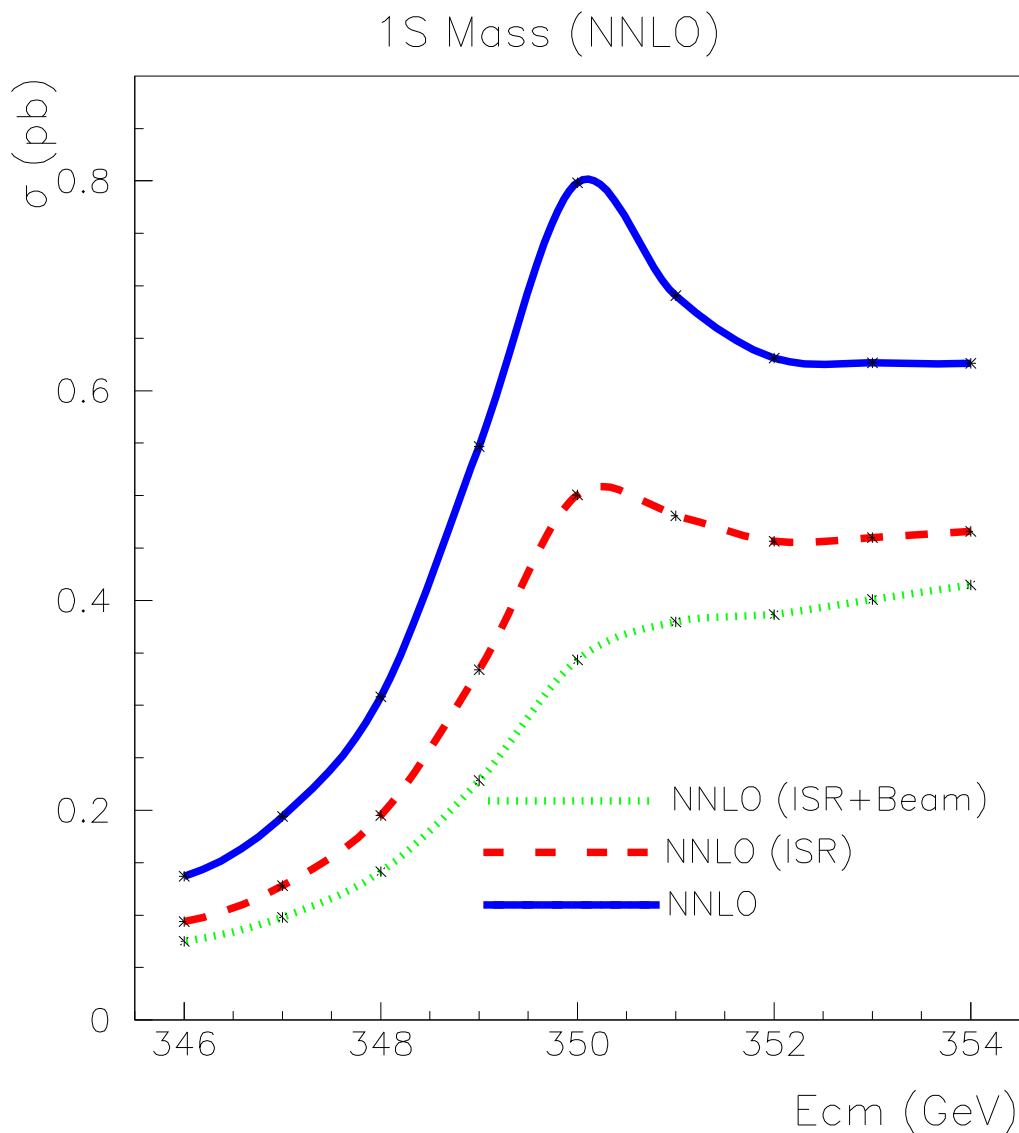
- large scale dependence and huge shift of the peak from order to order using pole mass
- theoretical error on  $m_t$  unclear
- in addition the same uncertainty appears going from  $m_{\text{pole}}$  to  $m_{\overline{\text{MS}}}$
- both problems can be solved by redefining the mass, shifting part of the potential to the mass definition



( $v$ =top velocity parameter, should be  $> 0.15$ )

Threshold cross section now very well under control

- In the QCD-corrected cross section some remnant of the 1S peak remains visible



- This peak is completely washed out by ISR, beamstrahlung and beamenergy-spread
- However uncertainties in beam parameters do not effect precision of  $m_t$  measurement

## Additional information:

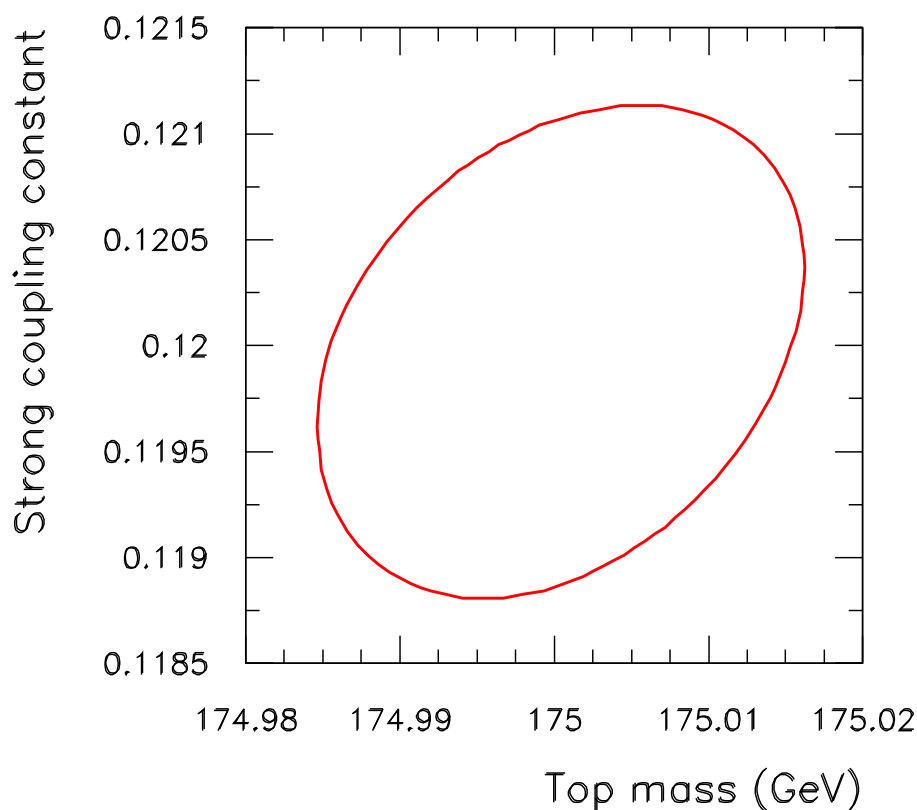
- absolute value of the  $t\bar{t}$  cross section: sensitive to  $\alpha_s, \Gamma_t$
- Momentum distribution of top quarks near threshold sensitive to  $m_t$
- Forward backward asymmetry: sensitive to  $\Gamma_t$
- Can try multi-parameter fits

Results (10 scan points with  $\mathcal{L} = 30 \text{ fb}^{-1}$  each):

$$\Delta m_t = 34 \text{ MeV}$$

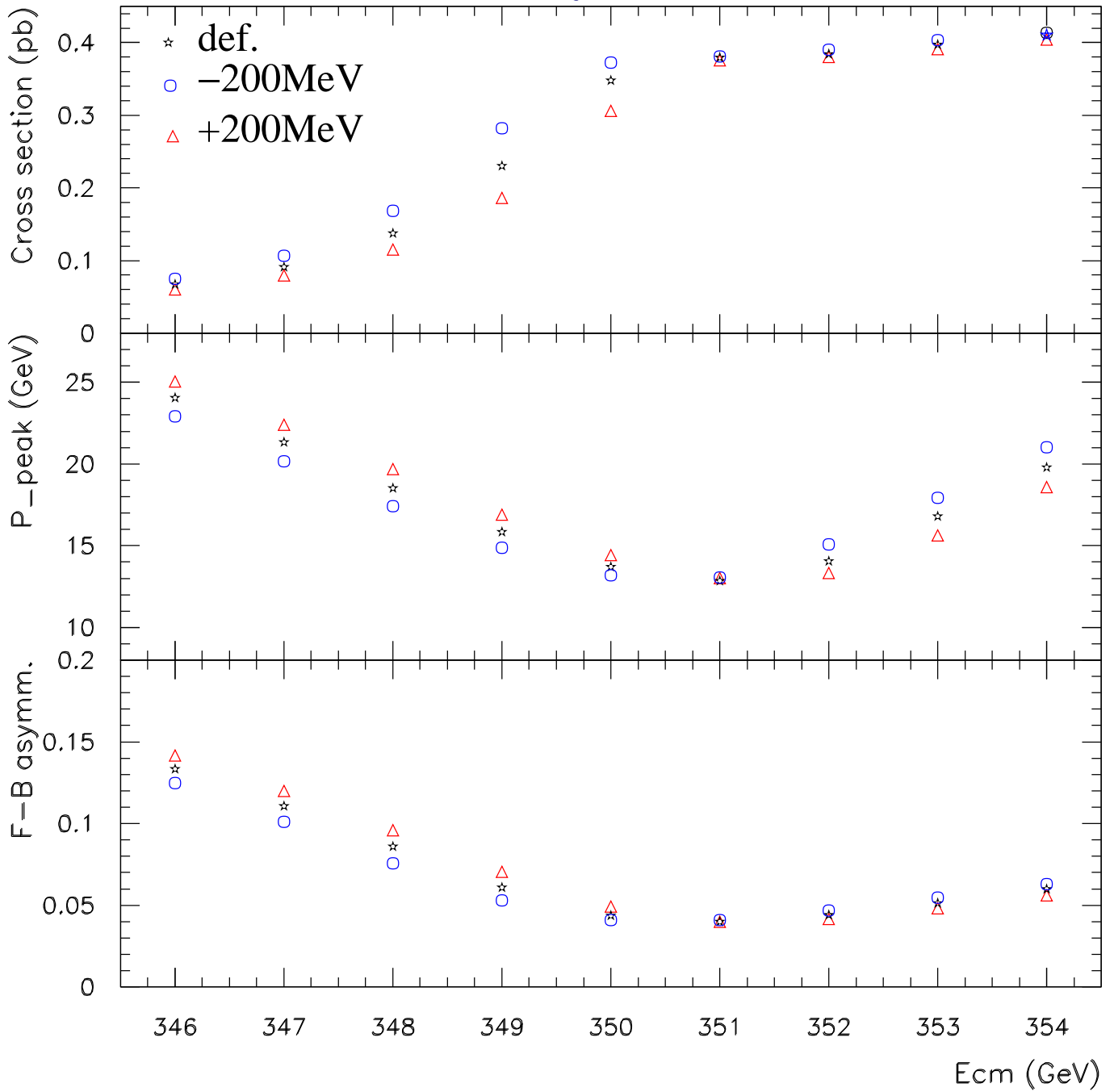
$$\Delta \Gamma_t = 42 \text{ MeV}$$

$$\Delta \alpha_s = 0.0023$$

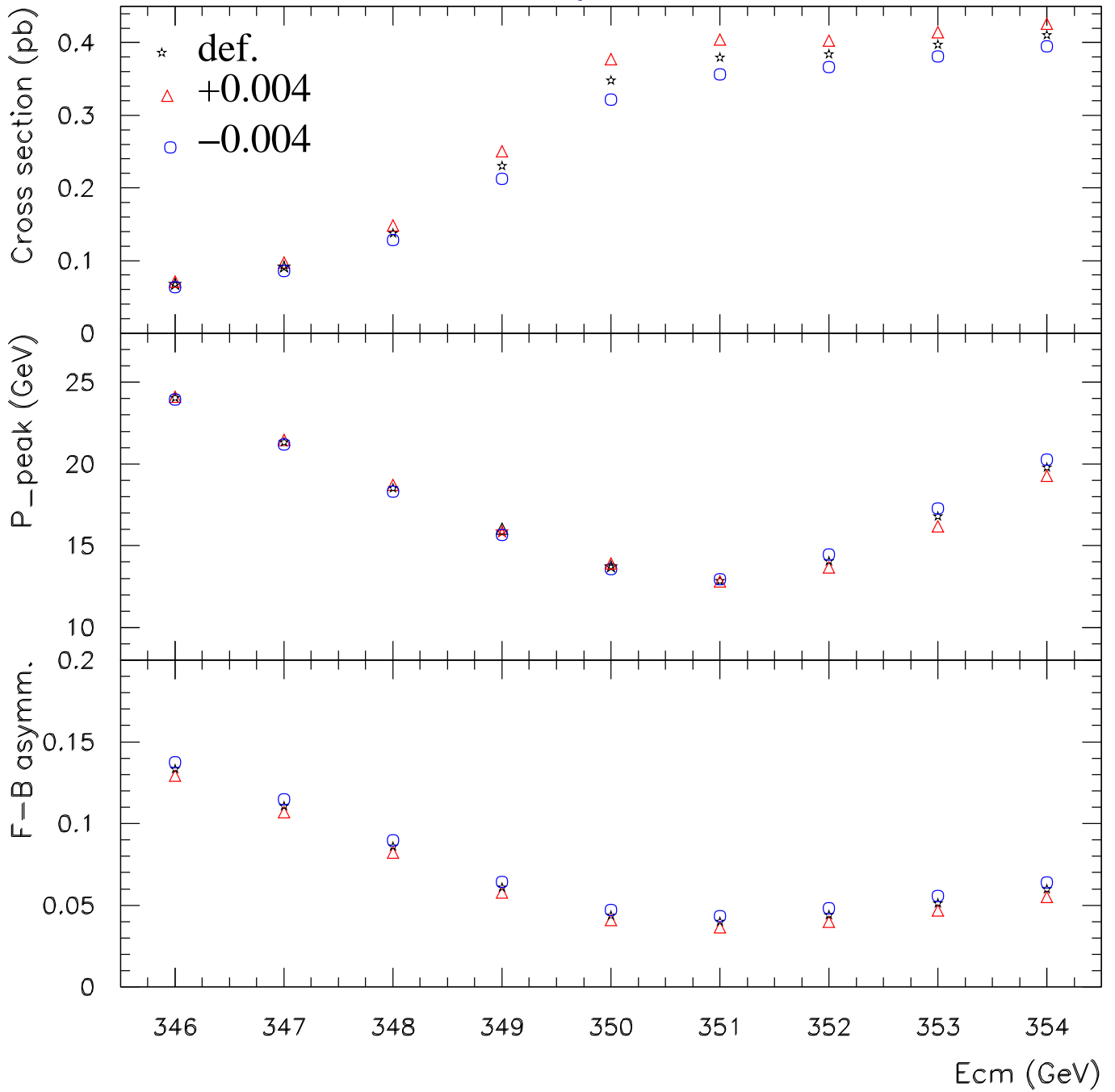




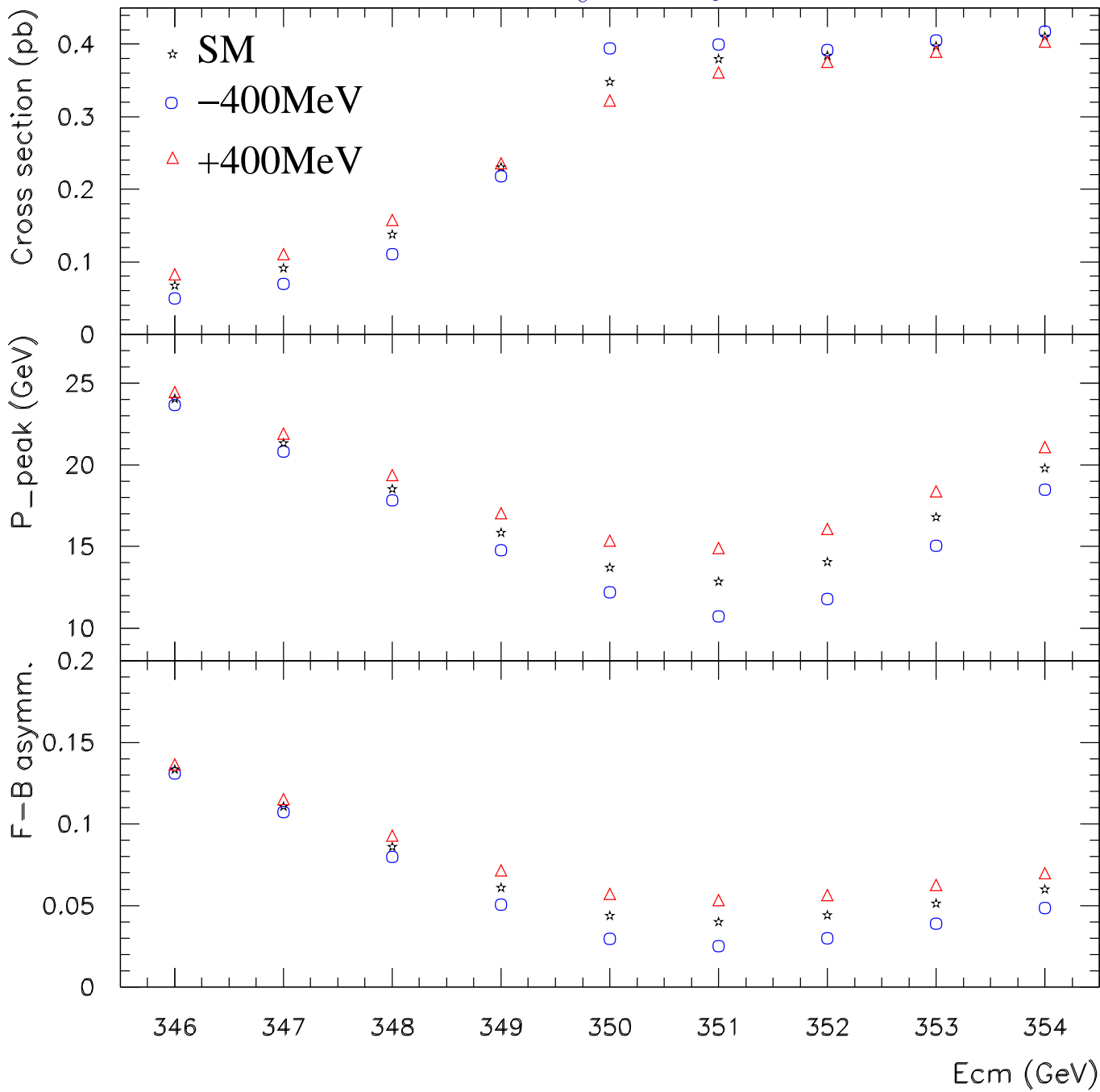
# Sensitivity to $m_t$



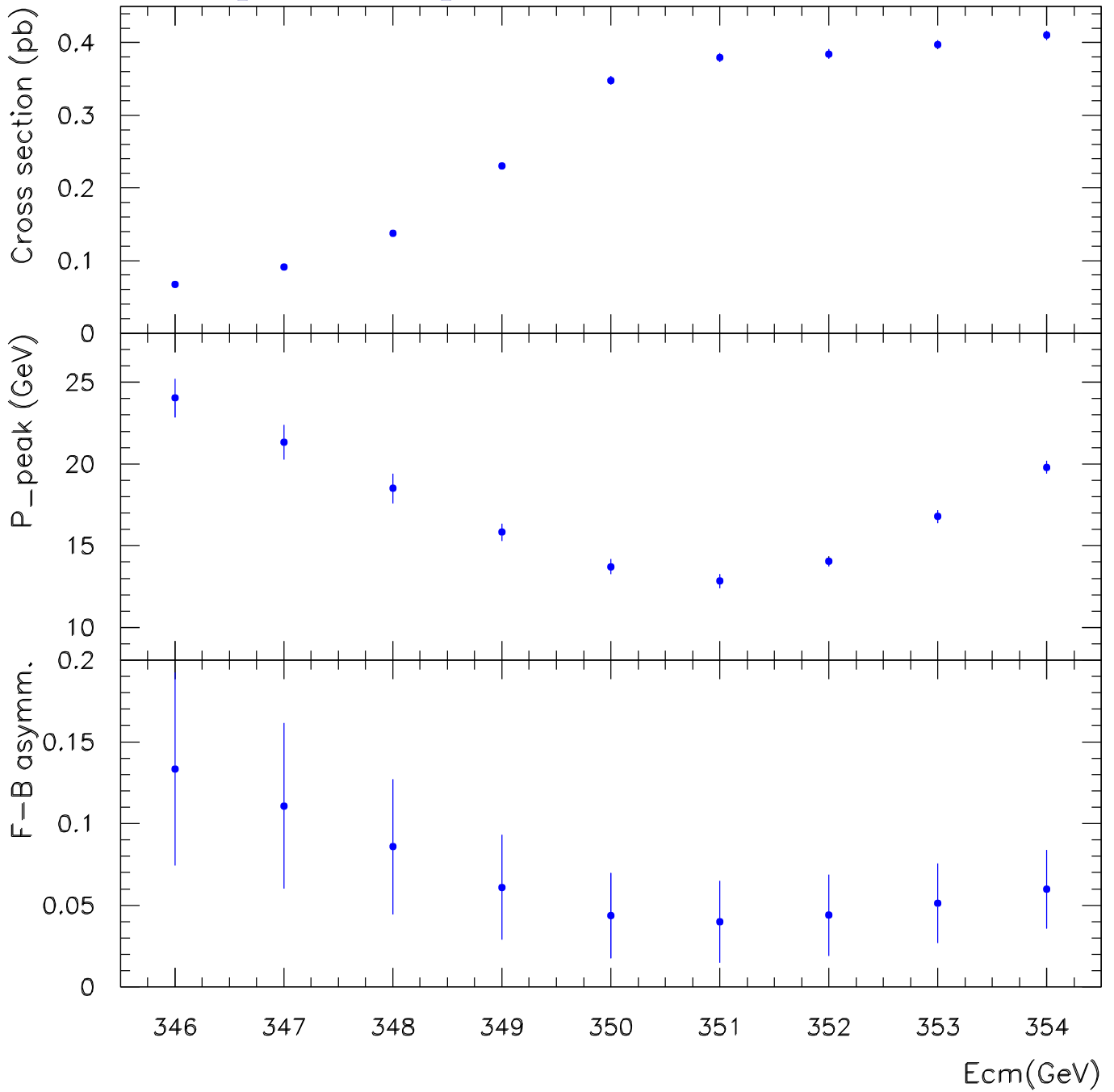
# Sensitivity to $\alpha_s$



# Sensitivity to $\Gamma_t$



# Expected experimental scan results



One step further:

The absolute value of the cross section is also sensitive to the Ht Yukawa coupling ( $y_t$ )

⇒ can take  $\alpha_s$  from other measurements and fit  $y_t$  instead

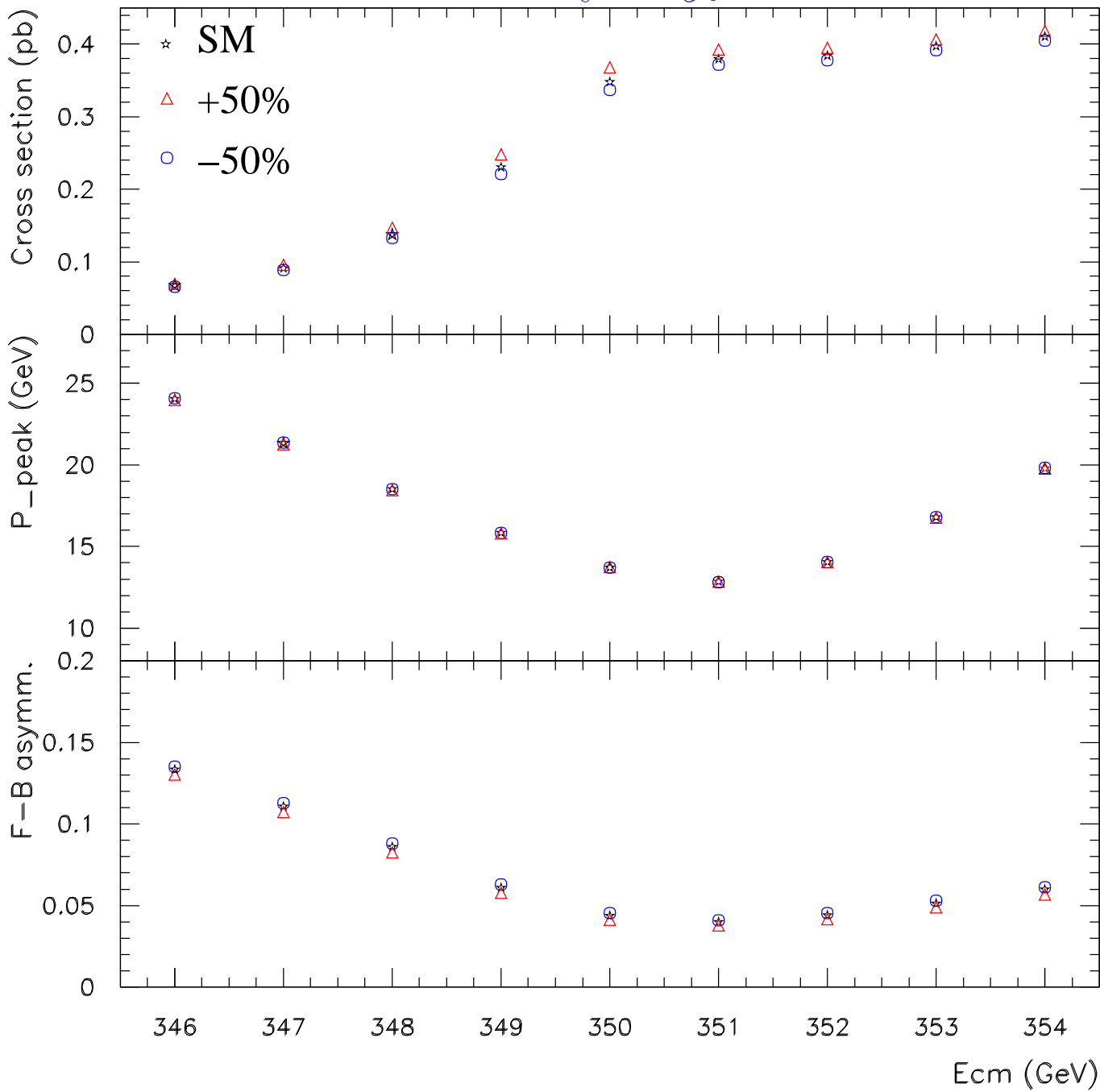
Result:

$$\frac{\Delta y_t}{y_t} = \begin{matrix} +0.35 \\ -0.65 \end{matrix}$$

for a 3-parameter fit

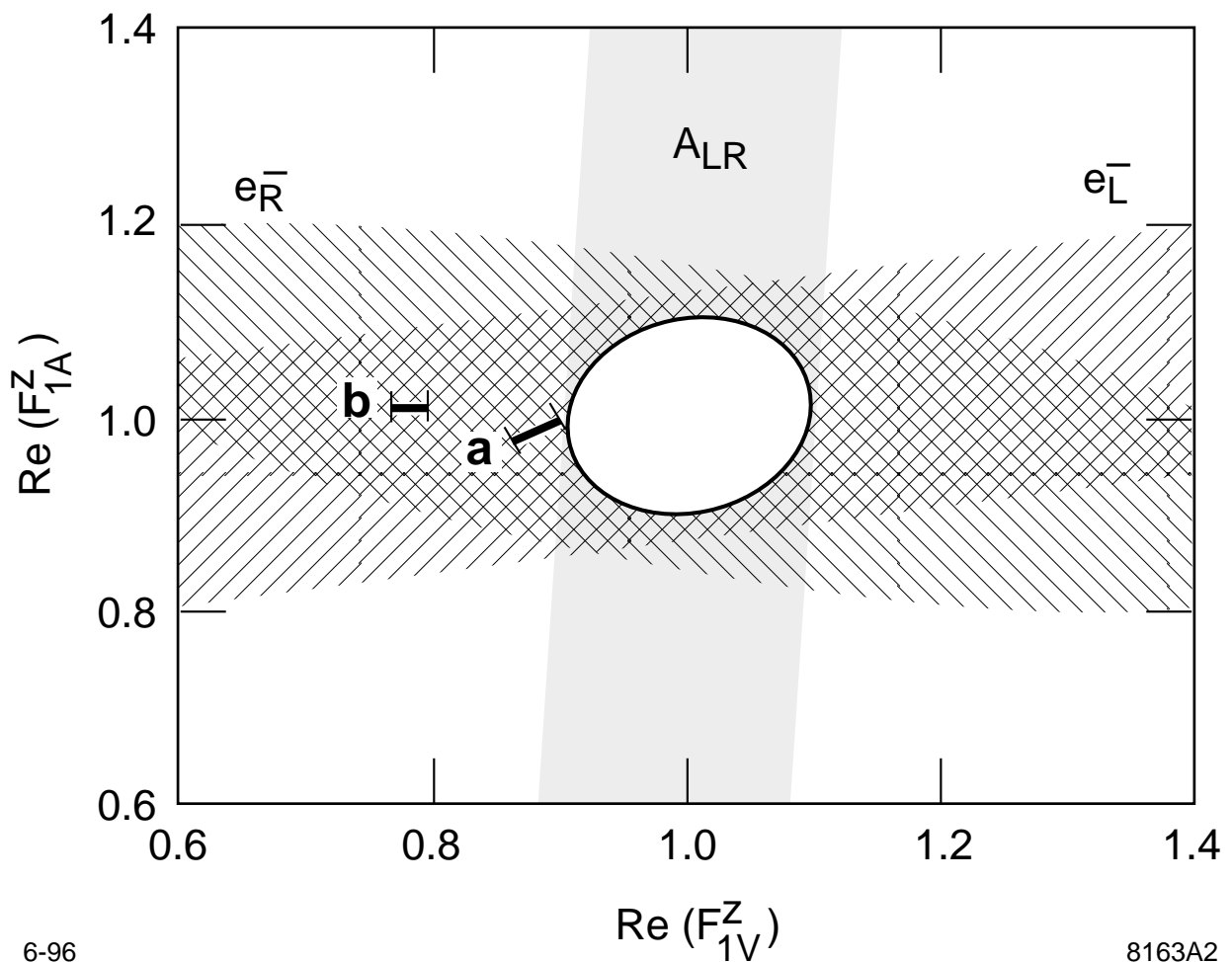
(LHC:  $\Delta m_t \approx \pm 1.5 \text{ GeV}$ )

# Sensitivity to $y_t$



## $t\bar{t}Z$ couplings

- the top-couplings to the Z can be obtained from  $t\bar{t}$ -production in the continuum
- due to the interference between Z and  $\gamma$  exchange the total cross section and the left-right asymmetry are sensitive to the Z-couplings
- a very conservative analysis at  $\sqrt{s} = 400$  GeV gives 90% c.l. limits on the 10% level for anomalous couplings



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- this precision is sensitive to some ETC-models

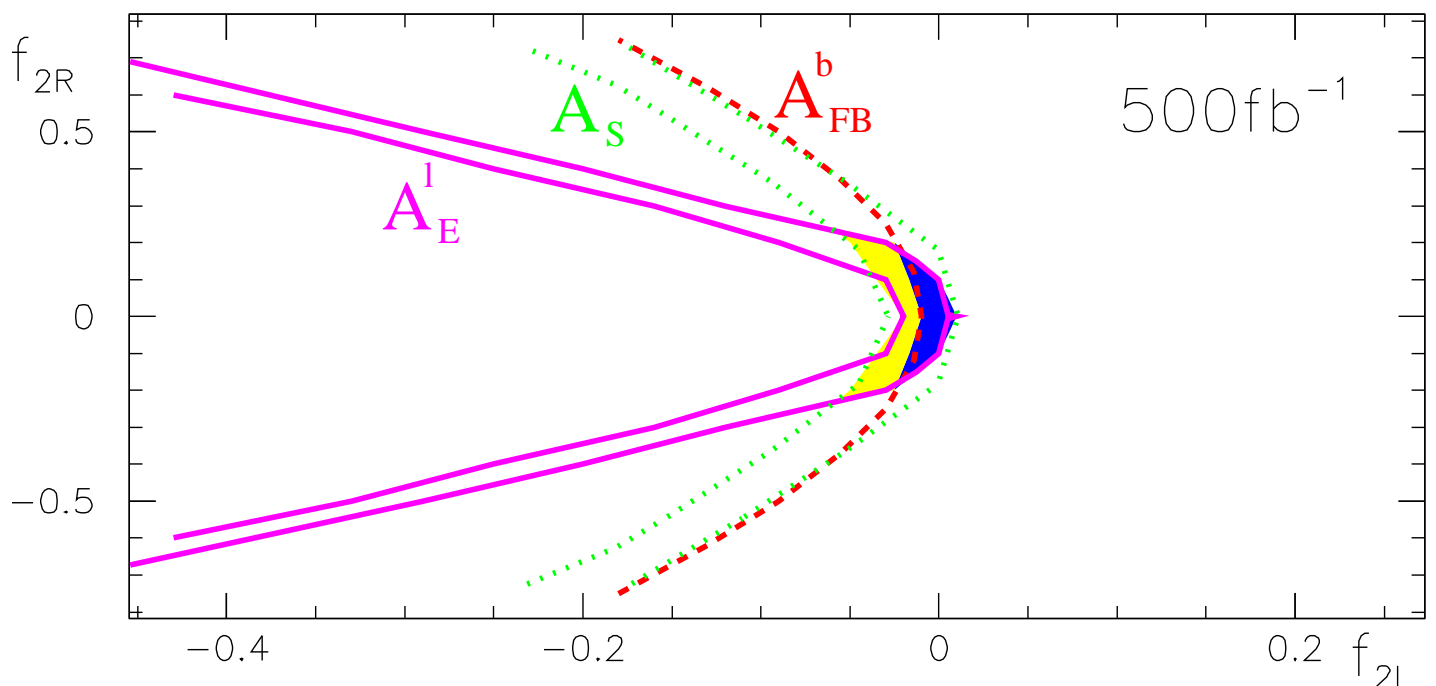
## Wtb couplings

- effective Lagrangian:

$$\mathcal{L} = \frac{g}{\sqrt{2}} [W_\mu^- \bar{b} (\gamma_\mu f_{1L} P_- + \gamma_\mu f_{1R} P_+) t - \frac{1}{2M_W} W_{\mu\nu} \bar{b} \sigma^{\mu\nu} (f_{2R} P_- + f_{2L} P_+) t] + \text{h.c.}$$

( $P_\pm = 1/2(1 \pm \gamma_5)$ ) SM:  $f_{1L} = 1$ , rest=0

- Present data put tight constraints on  $f_{1L}, f_{1R}$ , so try to measure  $f_{2L}, f_{2R}$
- use  $e^+e^- \rightarrow t\bar{t} \rightarrow X\ell\nu$  at  $\sqrt{s} = 500$  GeV and assume  $t\bar{t}Z$ -vertex to be standard
- analyze  $A_{\text{FB}}^b, A_{\text{FB}}^\ell$  and lepton energy in top rest frame
- results:  $\Delta f_{2L} \approx 0.02, \Delta f_{2R} \approx 0.2$





## Conclusions on top-quark physics

- The  $t\bar{t}$  threshold seems theoretically well under control
- The top quark mass can be measured to  $\sim 50$  MeV which is more than one order of magnitude better than what LHC can do
- the top width can be measured on the 3% level in the threshold scan
- the LC is the unique place to test  $t\bar{t}Z$ -couplings and can do that with a precision to better than 10%
- for t-decay physics probably the LHC is better due to much higher statistics