

- Contact interactions
- Models with Z's
- Extra dimensions
- Conclusions

Contact Interactions

Very heavy exchange-particle: Propagator $\propto \frac{1}{M^2}$ Effective Lagrangian:

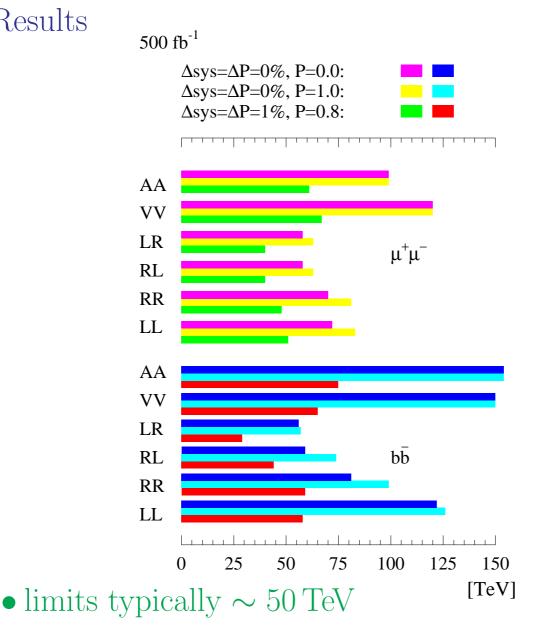
 $\mathcal{L}_{eff} = \sum_{i,k=L,R} \lambda_{ik}^2 / M^2 \alpha^{ik} (\bar{e}_i \gamma^{\mu} e_i) (\bar{f}_k \gamma^{\mu} f_k)$ with $\alpha^{ik} = \pm 1$ Scale-parameter $\Lambda^2 = \frac{4\pi M^2}{\lambda^2}$ (e.g. $\mu \operatorname{decay} \Lambda = (\sqrt{2}G_{\mu})^{-1/2} \sim 250 \,\mathrm{GeV})$ $\frac{d\sigma}{d\cos\theta} = SM(s,t) + C_2(s,t)\frac{1}{\Lambda^2} + C_4(s,t)\frac{1}{\Lambda^4}$

(Equivalent to t-channel exchange of a heavy scalar with mass M and coupling λ)

Main sensitivity is in interference term, so large dependence on helicity structure Assumptions

- $\sqrt{s} = 500 \,\text{GeV}, \, \mathcal{L} = 500 \,\text{fb}^{-1}$
- b-tagging efficiency $\varepsilon_{\rm b} = 60\%$
- systematic error 0, 1% (pessimistic)

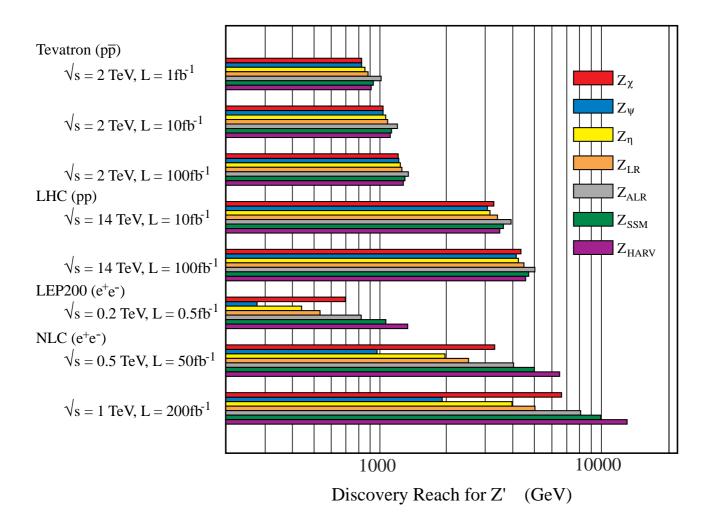
Results



- systematics will dominate, otherwise $\Lambda_{\rm lim} \propto \mathcal{L}^{1/4}$
- polarization helps little
- LHC reach similar but in different channels

Models with Z's

- models with extended gauge groups (left-right-symmetric, E_6) normally require additional Z-bosons
- in principle Z and Z' mix, however Z Z' mixing angle tightly constrained by Z-precision data
- for direct production LHC reaches much higher Z'-limits than LC ($\sim 3\,{\rm TeV})$
- however for ff-production Z'-exchange interferes with Z and γ exchange so that Z'-effects remain visible for $m_{Z'} \gg \sqrt{s}$ (in the same way PEP and PETRA could measure properties of the Z)
- measurement of cross sections and asymmetries gives access to vector- and axial-vector-couplings separately
- model dependent analyzes:
 - -assume a given model
 - \implies all couplings are defined
 - $-\operatorname{can}$ use leptonic and hadronic events
 - -deviations from SM prediction translate directly into Z'-mass



- (very moderate Luminosity assumptions for LC, however statistical scaling only with $\mathcal{L}^{1/4}$ and large contributions from Luminosity systematics)
- on average limits comparable to LHC
- however much larger difference between models, since sensitivity is in interference term
- on the contrary LC is not sensitive to the total width of the Z'

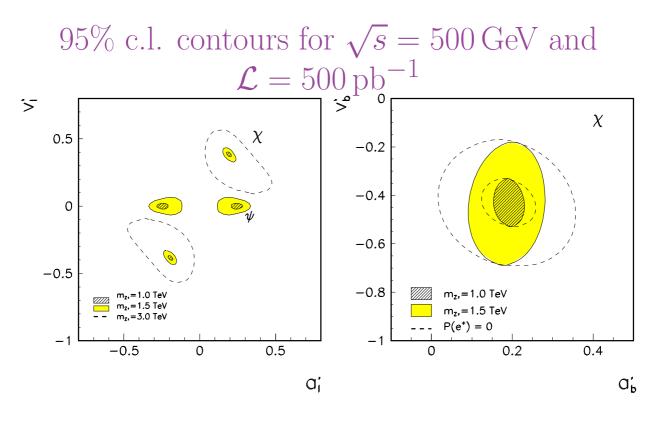
• model independent analyzes:

-LC sensitive to normalized couplings

$$a_f^N = a'_f \sqrt{\frac{s}{m_{Z'}^2 - s}}$$
$$v_f^N = v'_f \sqrt{\frac{s}{m_{Z'}^2 - s}}$$

- for leptons can obtain model independent limits/measurements on normalized couplings
- -all hadronic observables depend on product of leptonic couplings (Z'-production) and hadronic couplings (Z'-decay)
- ➡ can measure hadronic couplings only if leptonic couplings deviate significantly from zero
- experimental assumptions:
 - -beam polarizations 90/60% with $\Delta \mathcal{P}/\mathcal{P} = 1\%$
 - $-\operatorname{luminosity}$ known to 0.5%
 - -leptons can be tagged with $\varepsilon = 95 \pm 0.5\%$
 - $-\,\mathrm{b}$ quarks can be tagged with $\varepsilon = 60 \pm 0.6\%$
 - -measure cross sections, A_{LR} and A_{FB}^{ℓ}

Ideal case: LHC discovers a Z', so mass is known and LC can measure the couplings



- measure leptonic couplings to few % and b-couplings to $\sim 10\%$ for $m_{Z'}=1.5\,{\rm TeV}$
- \bullet limits should roughly stay constant for $m_{Z'}/\sqrt{s} = {\rm const}$
- the LC can distinguish the models over basically the full LHC discovery range

Hierarchy-problem:

Why is $m_{\rm H} \sim 100 \,{\rm GeV} \ll {\rm M_{pl}} \sim 10^{19} \,{\rm GeV}?$

Possible answers:

- SUSY (already seen)
- in reality is $M_{\rm pl} \sim 100 \,\text{GeV}$ but it appears so large because gravity lives in 4 + n dimensions

$$\mathrm{M}_{\mathrm{pl}}^2 = \mathrm{M}_{\mathrm{D}}^{2+\mathrm{n}} \mathrm{R}^{\mathrm{n}}$$

R : compactification radius of extra dimensions

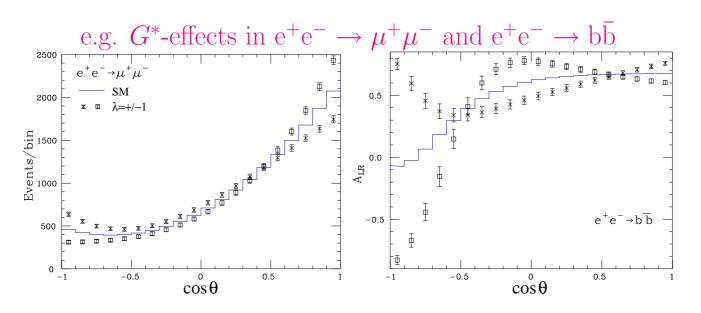
$$\Rightarrow R = M_{\rm pl}^{\frac{2}{\rm n}} M_{\rm D}^{-(\frac{2}{\rm n}+1)}$$

$$\sim 10^{\frac{30}{n}-17} \left(\frac{1\,{\rm TeV}}{M_{\rm D}}\right)^{1+\frac{2}{n}} [{\rm cm}]$$

n = 1	$R = \mathcal{O}(10^{13} \text{cm})$	excluded
n=2	$R = \mathcal{O}(1\text{mm})$	\sim excluded
n = 7	$R = \mathcal{O}(1 \text{fm})$	

Experimental signatures:

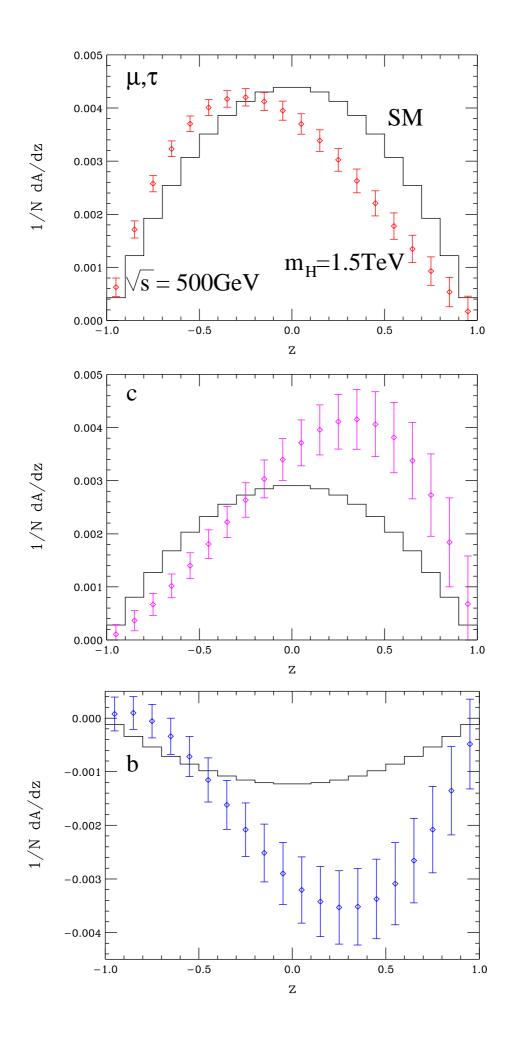
- In the bulk of the extra dimensions there live a huge number of graviton states (Kaluza-Klein towers G^*)
- → Expect effects in single γ production (e⁺e⁻ → γG^{*}, G^{*} invisible) and fermion pair production (e⁺e⁻ → G^{*} → ff)



- LC limit $M_D < 4(7)$ TeV for $\sqrt{s} = 0.5(1)$ TeV
- LHC comparable
- $\cos \theta(=z)$ dependence very different from Z'

Additional possibility: transverse polarization

- with transverse beam polarization there exists an azimuthal asymmetry depending on $\cos \theta \rightarrow \text{plot}$
- this asymmetry is symmetric in $\cos \theta$ for vector or scalar particle exchange
- for tensor exchange (gravitons) it receives an asymmetric component
- \twoheadrightarrow Graviton and Z' exchange can be distinguished up to $M < 10 \sqrt{s}$
 - extra dimensions can be excluded up to $M_D < 10(22)$ TeV for $\sqrt{s} = 0.5(1)$ TeV (highest reach at next generation colliders)



Conclusions on alternatives

- The LC is sensitive to a "General new physics scale" of order 50 TeV
- In concrete models (Z', extra dimensions) this translates into mass scales of few TeV
- LC and LHC have similar reach but are highly complementary
 - The LC is mainly sensitive to $e^+e^-\ell^+\ell^-$ and $e^+e^-b\bar{b}$ couplings while LHC is sensitive to $\ell^+\ell^-q\bar{q}$ (q=u,d)
 - LHC mainly sees the pure new physics while
 LC sees its interference with the SM
 - The LHC can discover that there is "something new" by seeing a resonance, then the LC can distinguish models by measuring the couplings