

TESTING THE ELECTROWEAK
STANDARD THEORY
AND
MEASURING STRUCTURE FUNCTIONS
AND QUARK DISTRIBUTIONS AT
HERA

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PHE 87-3; CONTR. TO HERA PHYSICS WORKING
GROUPS

- EW. TESTS (HAIDT, RÜCKL)
- STRUCTURE FUNCTIONS & DIS
(FELTESSE, INGELMANN)

1. INTRODUCTION
2. BASIC RELATIONS
3. TESTS OF THE EW-THEORY (GWS):
STATISTICAL & THEOR. ERRORS
4. TESTING EXTENSIONS OF THE
STANDARD THEORY
5. STRUCTURE FUNCTIONS &
QUARK DISTRIBUTIONS : ep
6. STRUCTURE FUNCTIONS :
THE DEUTERON OPTION
7. CONCLUSIONS

ELECTROWEAK TESTS

A. EXPERIMENTAL OPTIONS

0. $\lambda=0$, e^\pm -MEASUREMENT, $Y > .1 \rightarrow B(0)$

$$M_Z = \text{const.}$$

$$\Delta \sin^2 \theta_W = .041$$

$$\Delta M_W = 2.1 \text{ GeV } c^{-2}$$

$$L_{\text{tot}} = 200 \text{ pb}^{-1}$$

(e^+e^- - each 100 pb^{-1})

1. JET-MEASUREMENT CC & NC, $Y > .01$

$$R_-(0)$$

2. $|\lambda| \approx .8$ $R_-(\lambda), A_-(\lambda)$

FIGS.

WA: $Q^2 > 500$

$$X > .1$$

$$Y > .01$$

$$|\lambda| = .8$$

$$R_- < 50, \geq 2\% \text{ CC}, 5\% A_-$$

RED. SEA UNCERTAINTY

JET'S & e^- -MEASUREMENT

($Y > .1$, e^- -MEASUREMENT ONLY)

TOOLS OF INVESTIGATION

(i) STANDARD MODEL TESTS AT
HIGH SPACELIKE Q^2

$$SU_3^C \times [SU_2^L \times U_1^Y]$$

(ii) SEARCH FOR NOVEL PHENOMENA (CASHMERE REPORT)

- DEVIATIONS IN INCLUSIVE DISTRIBUTIONS
(STRUCTURE FUNCTIONS)
- NEW (EFFECTIVE) COUPLINGS AT HIGH Q^2
- INDIVIDUAL EVENTS WITH NEW SIGNATURES

INTRODUCTION

HERA : STUDY OF $e^\pm(\lambda)p$ -COLLISIONS
IN 2 EXPERIMENTS : H1 & ZEUS

AT : $4 \cdot E_e \cdot E_p = 4 \times 30 \times 820 \text{ GeV}^2 = S$

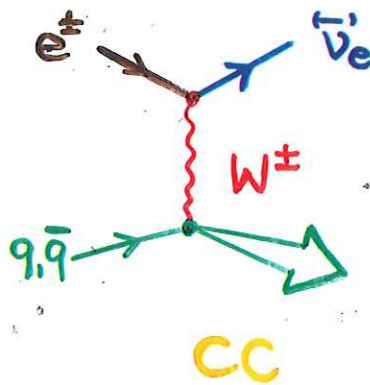
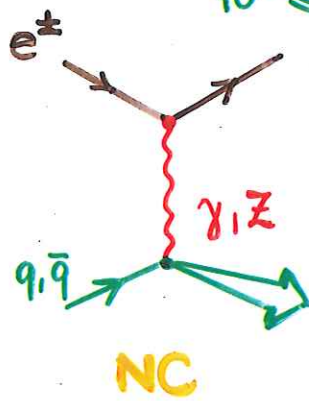
$\mathcal{L} = 1.5 \dots 2 \cdot 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$

$\int dt \mathcal{L} = \underline{200 \text{ pb}^{-1}}$ in 154... 116 days

$\sqrt{s} = 315 \text{ GeV}$

$30 \leq Q^2 \leq 30000 \text{ GeV}^2$

$10^{-4} \leq x \leq .8$



$|\lambda_e| \lesssim .8$

(d-TARGET ?)

BASIC RELATIONS

A. CROSS SECTIONS

$$\frac{d^2\sigma_{NC}^{\pm}}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4 x} \left\{ Y_+ [F_2 + \kappa_z C_1 G_2 + \kappa_z^2 C_2 H_2] \right. \\ \left. + Y_- [\underset{\substack{\uparrow \\ |y|^2}}{\kappa_z C_3 x G_3} + \underset{\substack{\uparrow \\ |z|^2}}{\kappa_z^2 C_4 x H_3}] \right\}$$

$$\frac{d^2\sigma_{CC}^{\pm}}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4 x} \kappa_W^2 \frac{1 \mp \lambda}{2} \left\{ Y_+ W_2^{\pm} \pm Y_- x W_3^{\pm} \right\}$$

$$Y_+ = 2 - 2y + y^2 / (1+R) ; Y_- = 1 - (1-y)^2$$

$$C_i = C_i(v_j, a_j, \lambda)$$

$$v_j = I_{3j}^L - 2Q_j \sin^2\theta$$

$$a_j = I_{3j}^L$$

B. STRUCTURE FUNCTIONS

NC
S: $(F_2, G_2, H_2) = x \sum (q + \bar{q}) \{ Q_q^2, 2Q_q v_q, v_q^2 + a_q^2 \}$
NS: $(xG_3, xH_3) = x \sum (q - \bar{q}) \{ Q_q a_q, v_q a_q \}$

CC
S: $(W_2^+, W_2^-) = (x \sum d_i + \bar{u}_i, x \sum u_i + \bar{d}_i)$
NS: $(xW_3^+, xW_3^-) = (x \sum d_i - \bar{u}_i, x \sum u_i - \bar{d}_i)$

9 STRUCTURE FUNCTIONS FROM

≈ 6 PARTON DISTRIBUTIONS

$u_v, d_v, u_s = d_s, s, c, b$, ~~+~~ at HERA (EHLQ NEW)

C. PROPAGATORS'

$$K_Z(Q^2) = \frac{Q^2}{Q^2 + M_Z^2} \frac{1}{4s_\theta^2 c_\theta^2}$$

$$K_W(Q^2) = \frac{Q^2}{Q^2 + M_W^2} \frac{1}{4s_\theta^2}$$

D. ELECTROWEAK PARAMETERS

$$\{\alpha, g, G_F, M_W, M_Z, s_\theta^2\}$$

GWS:

3 INDEPENDENT PARAMETERS!

$\alpha,$
or
 G_F
or
 $M_{Z,LEP}$

FIX,

MEASURE:

M_W OR $\sin^2\theta$ OR g, \dots

↑

$$\sin^2\theta = \sin^2\theta(\alpha, G_F, M_W)$$

$$g = g(\alpha, G_F, M_W)$$

⋮

etc.

E. OBSERVABLES

$$A^-(\lambda) = \frac{\sigma_{NC}^-(-\lambda) - \sigma_{NC}^-(\lambda)}{\sigma_{NC}^-(-\lambda) + \sigma_{NC}^-(\lambda)}$$

100 pb⁻¹ each

$$R^-(\lambda) = \frac{\sigma_{NC}^-(\lambda)}{\sigma_{CC}^-(\lambda)}$$

200 pb⁻¹ each

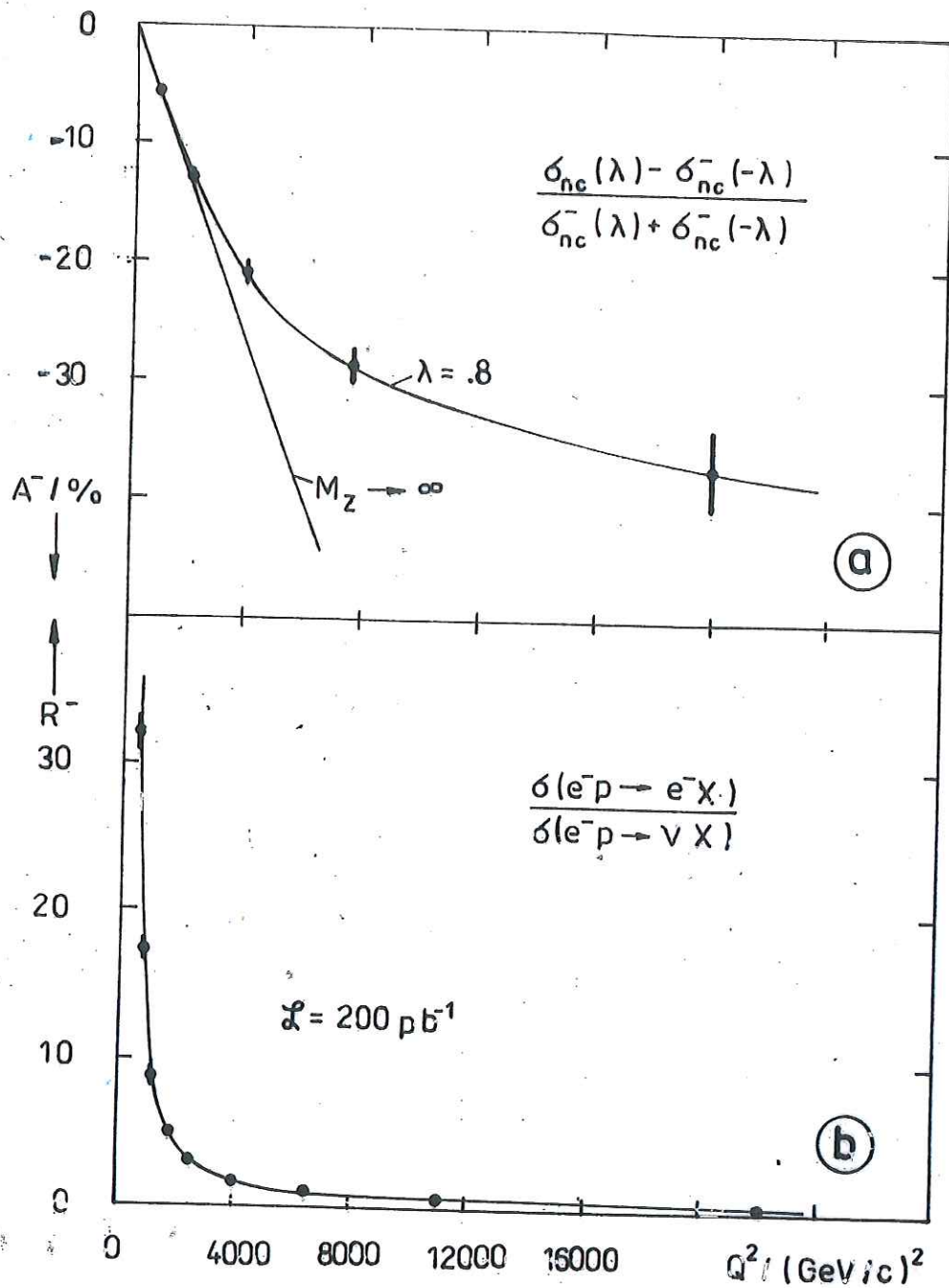
$$B(\lambda) = \frac{\sigma_{NC}^+(-\lambda) - \sigma_{NC}^-(\lambda)}{\sigma_{NC}^+(-\lambda) + \sigma_{NC}^-(\lambda)}$$

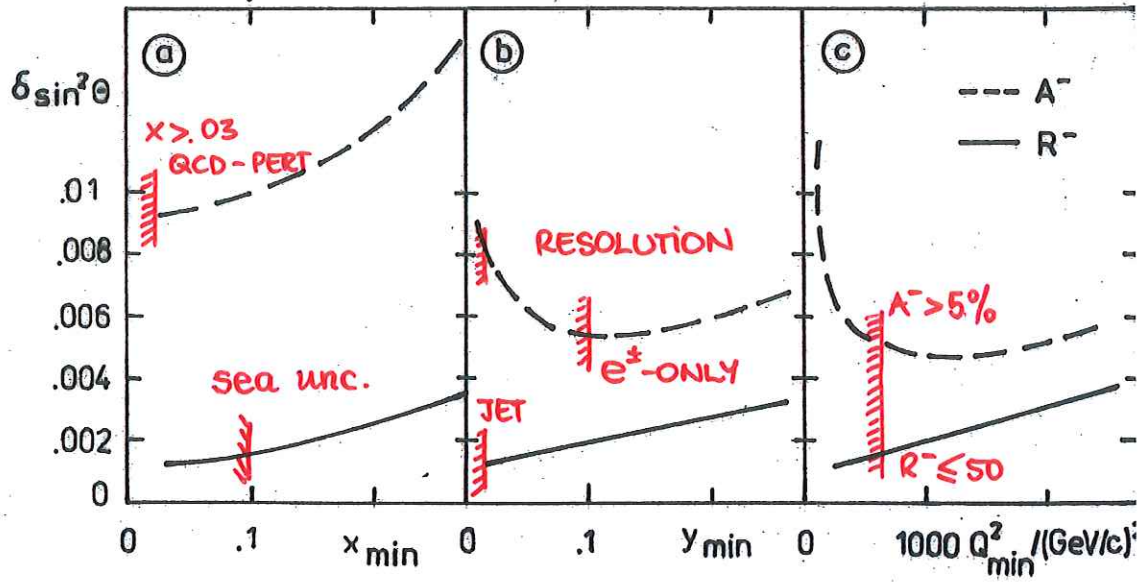
ONLY e's
 $\lambda = 0$ FINITE

$$A^+ \rightarrow I_{3R}e$$

OTHER CHOICES ARE EQUALLY POSSIBLE,
STAT. ERRORS ARE GREATER I.G.

FIG.





$M_z = \text{CONST.}$
OMS

Fig. 2

MADE USE OF MC-CALCULATIONS
OF THE H1-RESOLUTIONS FOR

$E_{\text{JET}}, \theta_{\text{JET}}; E_e, \theta_e$ (J. FELTESSE '85
'86)

: H1 PROPOSAL
MARCH 1986

B. STATISTICAL PRECISIONS: 1 PARAMETER FITTED

Tab. 1

fixed error	s^2_θ	M_Z	M_W	G_F
$\delta \sin^2 \theta$ A^-	-	.005	.005	.007 *)
R^-	-	.002 *)	.002	.005
$\delta M_Z / \text{GeV}$	3.02 .51	-	.25 .10	1.09 .76
$\delta M_W / \text{GeV}$	2.53 .45	.27 .10	-	1.38 .95

$M_Z^{\text{LEP}} \uparrow$

Statistical errors of $\sin^2 \theta$, M_Z , M_W in one-parameter fits assuming α and either s^2_θ , M_Z , M_W or G_F , resp., to be fixed. The upper values refer to A^- and the lower ones to R^- at $|\lambda| = .8$, $L = 200 \text{ pb}^{-1}$.

*) A RECENT ANALYSIS (APRIL 1987) BY D. HAIDT & H. PAUL HAS CONFIRMED THESE VALUES.

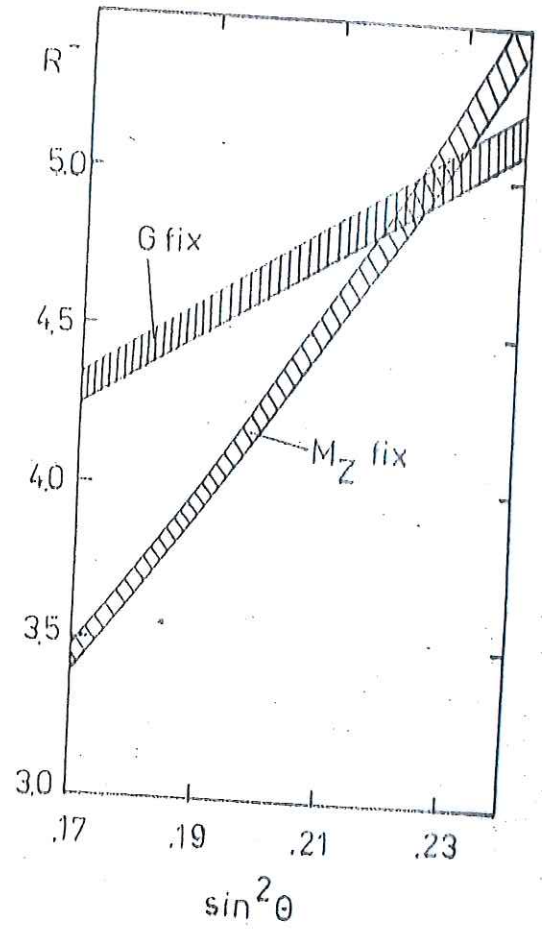
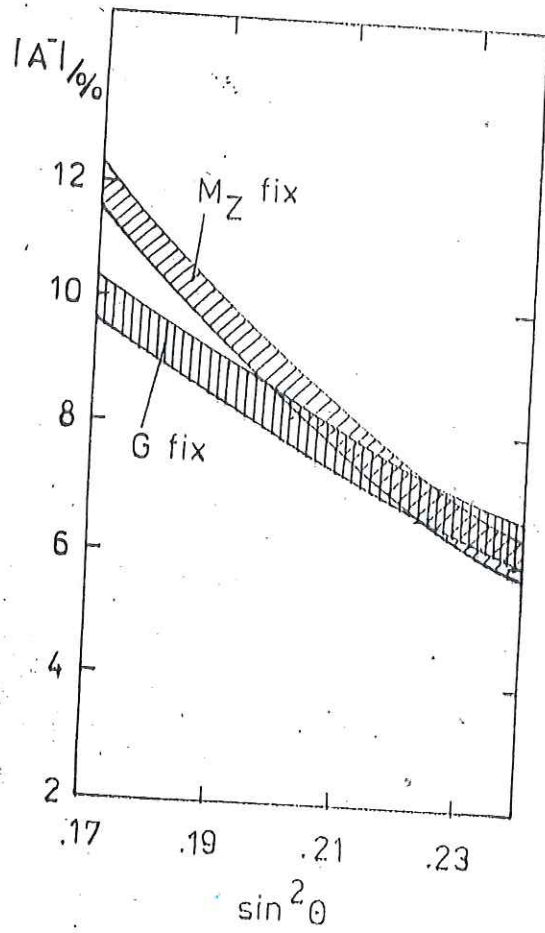


Fig. 2

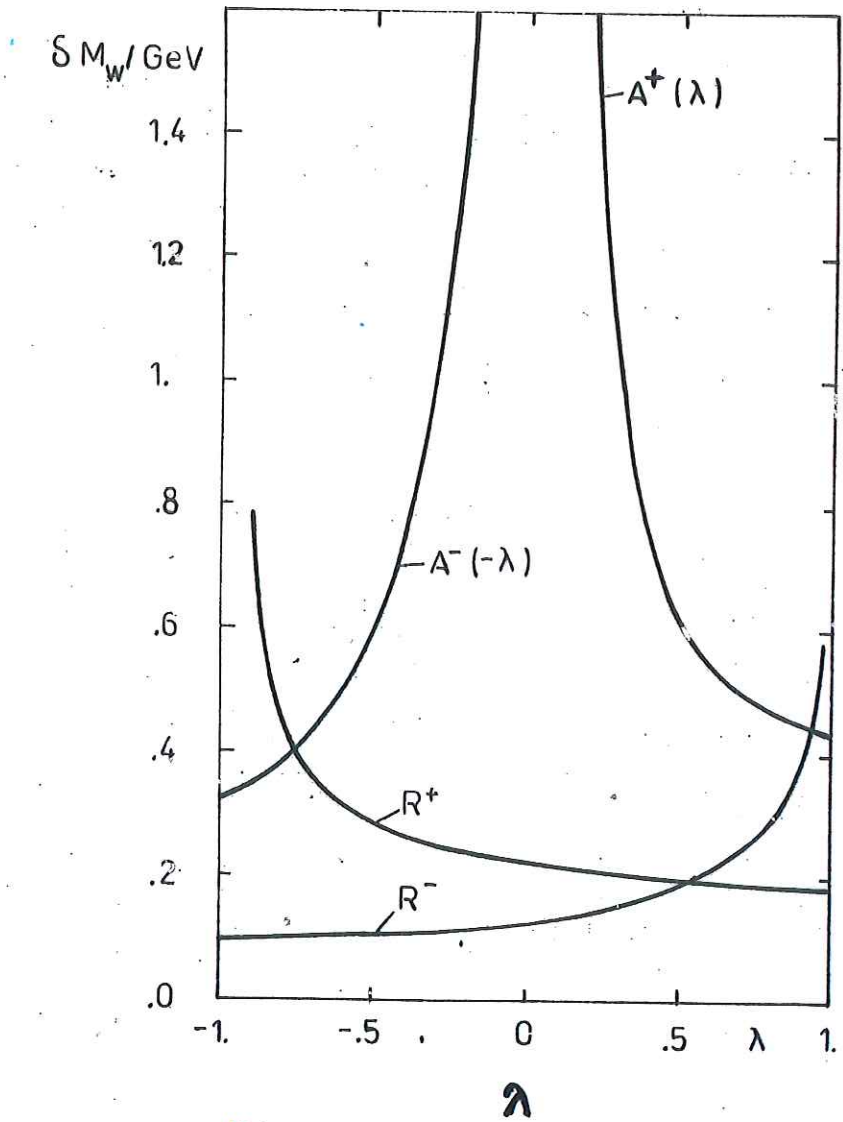
FIXING: $\alpha, \sin^2\theta$ AND FITTING

EITHER M_W OR M_Z :

R_- : $\Delta M_W = \pm 450 \text{ MeV}, \Delta M_Z = \pm 3 \text{ GeV}$

A_- : $\Delta M_Z = \pm 5.6 \text{ GeV}$

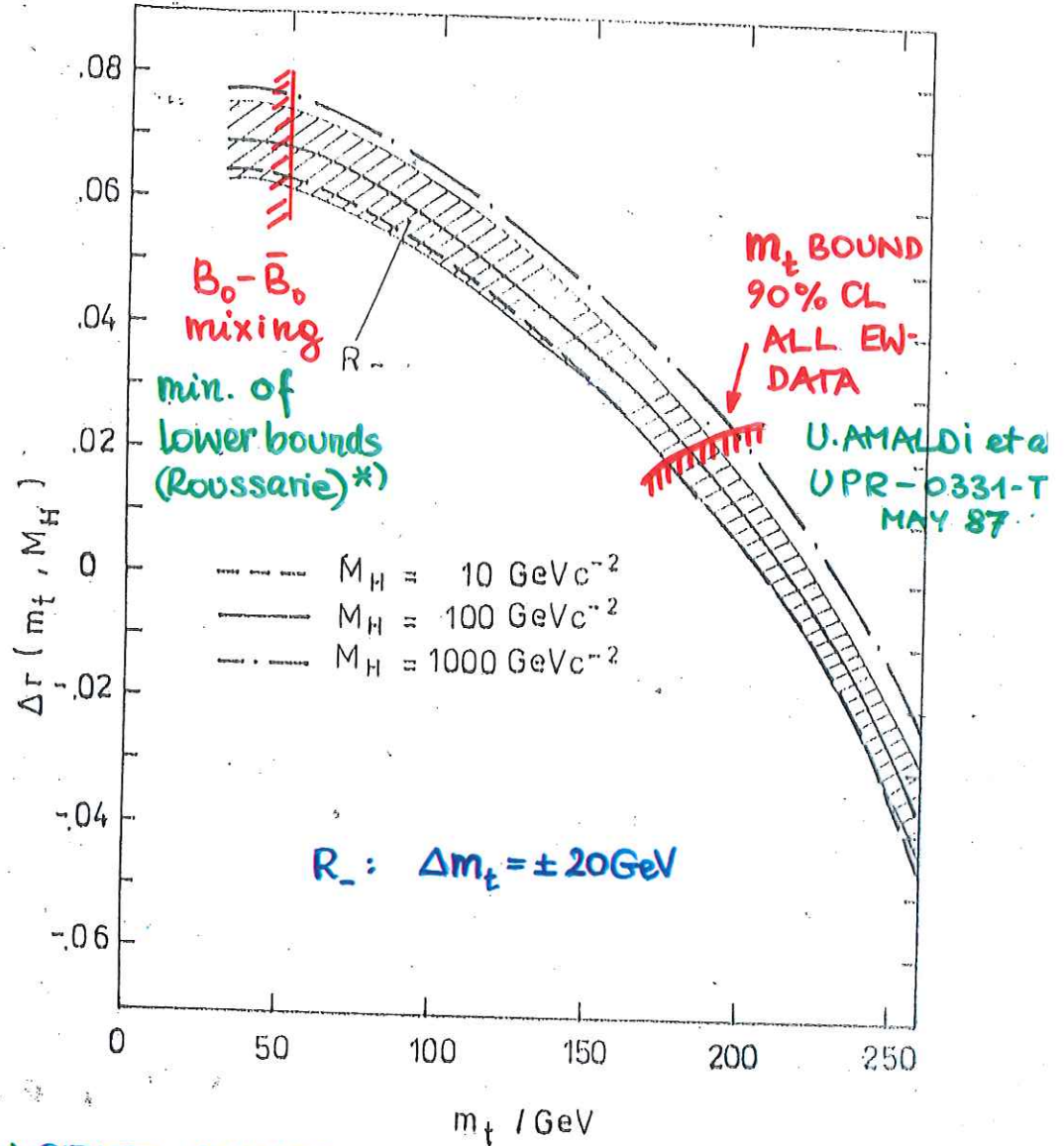
C. POLARIZATION EFFECTS



λ

	$ \lambda = .8$	$\lambda = 0$	
$\Delta M_W :$.104 GeV	→ .130 GeV	R-
$\Delta S_\theta^2 :$.002	→ .0023	R-

D. BOUNDING m_t & m_H



*) OTHER BOUNDS:
Bigi, Ellis et al.,
Roos et al.

Fig. 5

E. TWO PARAMETER FITS

eg.
FIX α , FIT M_W & M_Z .

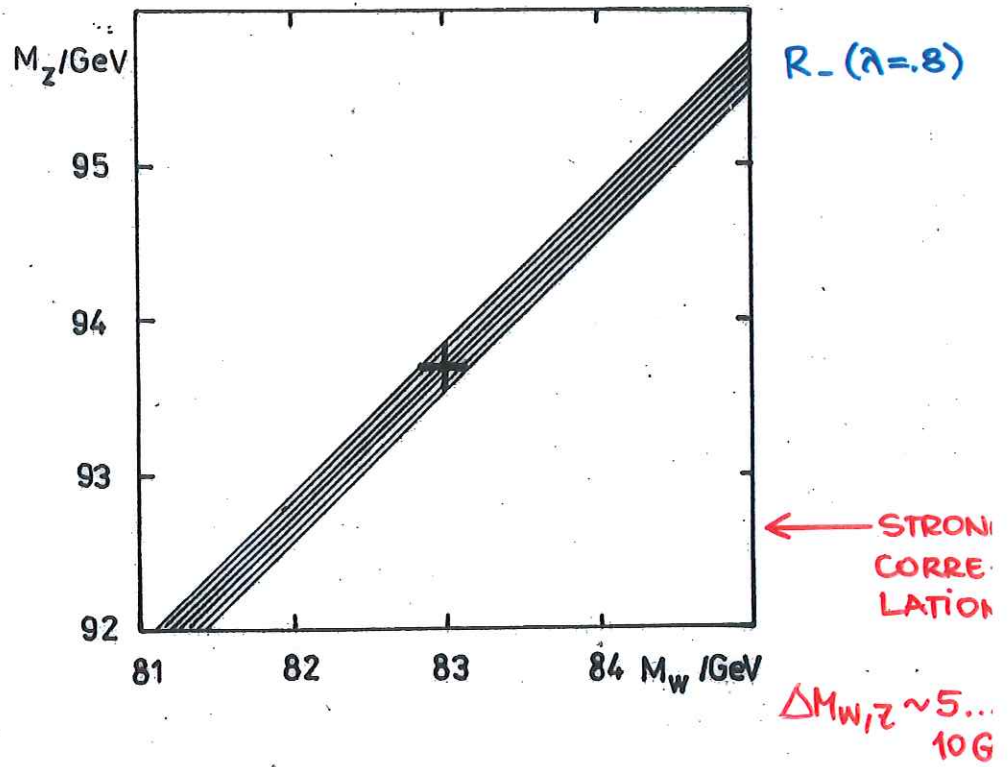


Fig. 4

- DO THE ERROR CONTOURS FOR A_- , R_- INTERSECT IN A WAY THAT DIMINISHES THE GLOBAL ASYM. ERRORS ?

→ AT HERA NOT.

$$|z|^2 \ll |\gamma|^2 + |\gamma z|$$

$$Q^2 \ll k_w^2, M_z^2$$

$$R_- = k_w^2 \phi_1 \rightarrow k^2 \phi_1$$

$$A_- = k_z \phi_2 \rightarrow k \phi_2 + O(k^2), k^2 \ll k^4 !!$$

$$\delta R_- \sim 2k \phi_1 \left\{ (\partial_{x_1} k)^2 \delta x_1^2 + (\partial_{x_2} k)^2 \delta x_2^2 \right\}^{1/2}$$

$$\delta A_- \sim 2\phi_2 \left\{ (\partial_{x_1} k)^2 \delta x_1^2 + (\partial_{x_2} k)^2 \delta x_2^2 \right\}^{1/2}$$

→ NEARLY PARALLEL ERROR CONTOURS

UNCERTAINTIES OF THE ANALYSIS

Tab. 2

A. PARAMETER-FIXING:

Fitted Fixed	$\delta \sin^2 \theta_W$		$\delta M_Z/\text{GeV}$		$\delta M_W/\text{GeV}$	
	A ⁻	K ⁻	A ⁻	R ⁻	A ⁻	K ⁻
$\Delta \sin^2 \theta_W$ = .001	-	-	- .58 + .57	\pm .26	- .54 + .53	\pm .20
ΔM_Z = \pm 50 MeV	\pm .0001	\pm .0002	-	-	+ .05 - .04	\pm .05
ΔM_W = \pm 50 MeV	\pm .0001	\pm .0002	+ .04 - .05	\pm .05	-	-

LEP:

Errors in $\sin^2 \theta_W$, M_Z , M_W due to the experimental error of fixed input quantities.

B. $R = \sigma_L / \sigma_T$

$$R=0 \rightarrow R=.1$$

$$A_- : \Delta S_\theta^2 = .0003$$

$$R_- : \Delta S_\theta^2 = .0008$$

C. m_q finite (m_b) ($Q^2 > 500 \text{ GeV}^2$)

$$R_- : \Delta S_\theta^2 = .0001$$

(m_c is already massless at this Q^2 .)

D. PARTON DISTRIBUTIONS

10% CHANGES OF EITHER u_v, d_v, s INDUCE:

δS_0^2	s	u_v	d_v
A ₋	$\mp .0003$	$\mp .0014$	$\pm .0011$
R ₋	$\pm .0017$	$\mp .0031$	$\pm .0015$

DO, $\Lambda = .2 \text{ GeV}$

$\sin^2 \theta = .217$

AN ANALYSIS OF THE MEASUREMENT OF F_2 AND $\sigma_{cc}^+ / \sigma_{cc}^-$ SHOWS, THAT THESE QUANTITIES CAN BE CONTROLLED ON THE LEVEL OF $\Delta/Q \sim 5\%$, I.E. THE PARTON-DISTRIBUTION UNCERTAINTIES ARE LESS THAN THE STAT. PRECISION.

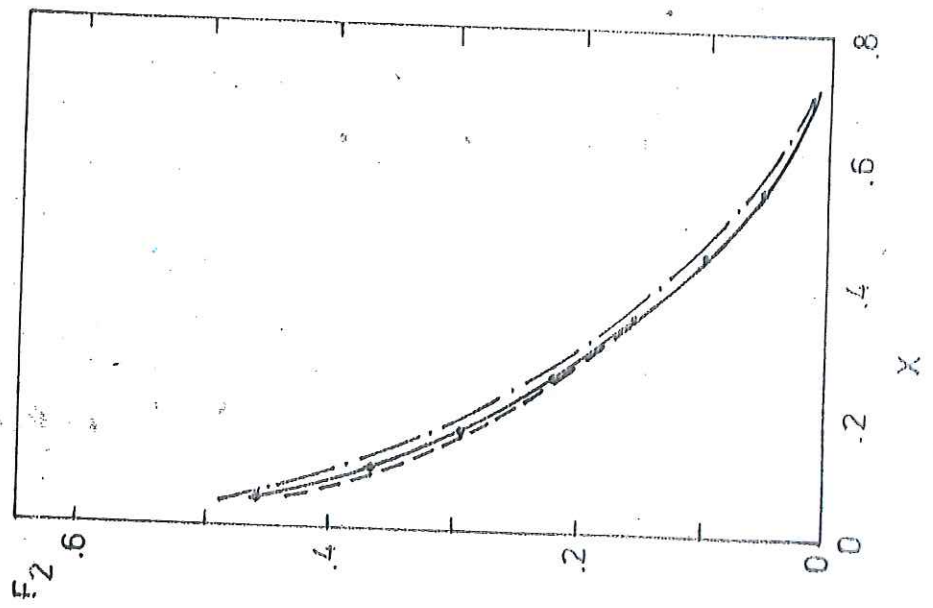


Fig. 6

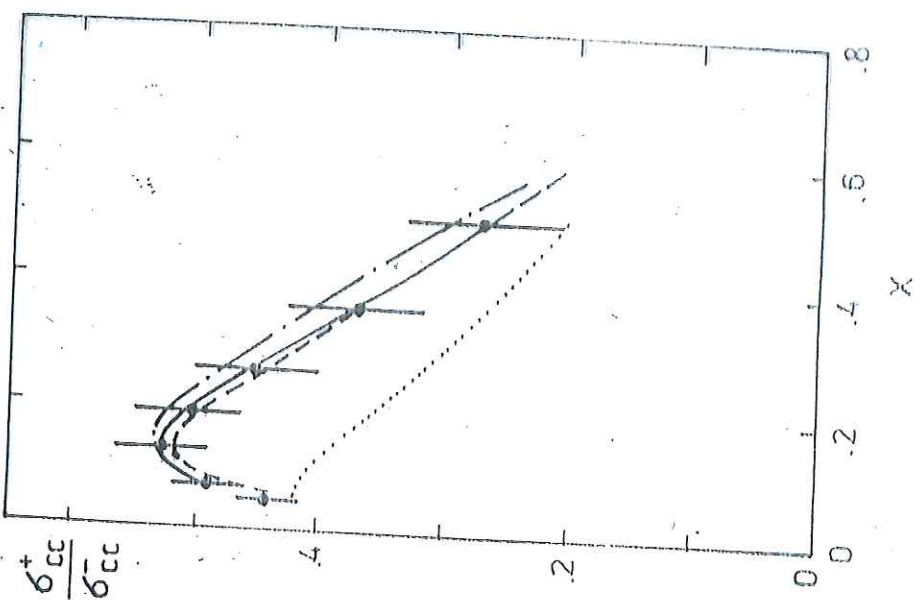
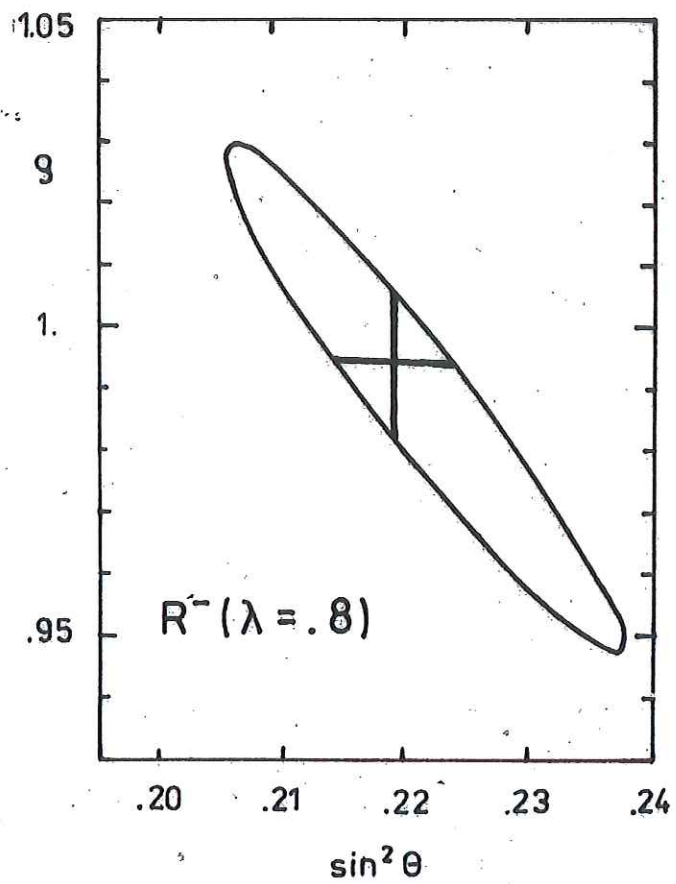


Fig. 7

EW-TESTS BEYOND GWS

A. g AS INDEPENDENT PARAMETER IN OMS



$$|\lambda| = .8$$

Fig. 8

$$\text{PARABOLIC ERROR: } \delta g = .01$$

$$\text{FIXING } \sin^2 \theta = \sin^2 \theta_{\text{LEP}} : \delta g = .003$$

B. FINITE RH- ISOSPIN COMPONENTS

$$V_i = I_{3Li} - 2Q_i \sin^2 \theta_W + I_{3Ri}$$

$$Q_i = I_{3Li} - I_{3Ri}$$

↑ GWS
≡ 0

$i \in [u, d, e]$

HIGHEST SENSITIVITY : $A_{\pm} (\lambda \rightarrow .8)$

i	δI_{3Ri}	
	$A^- (.8)$	$A^+ (.8)$
u	.016	.015
d	.034	.036
e	.056	.030

C. ADDITIONAL W'S OR Z'S

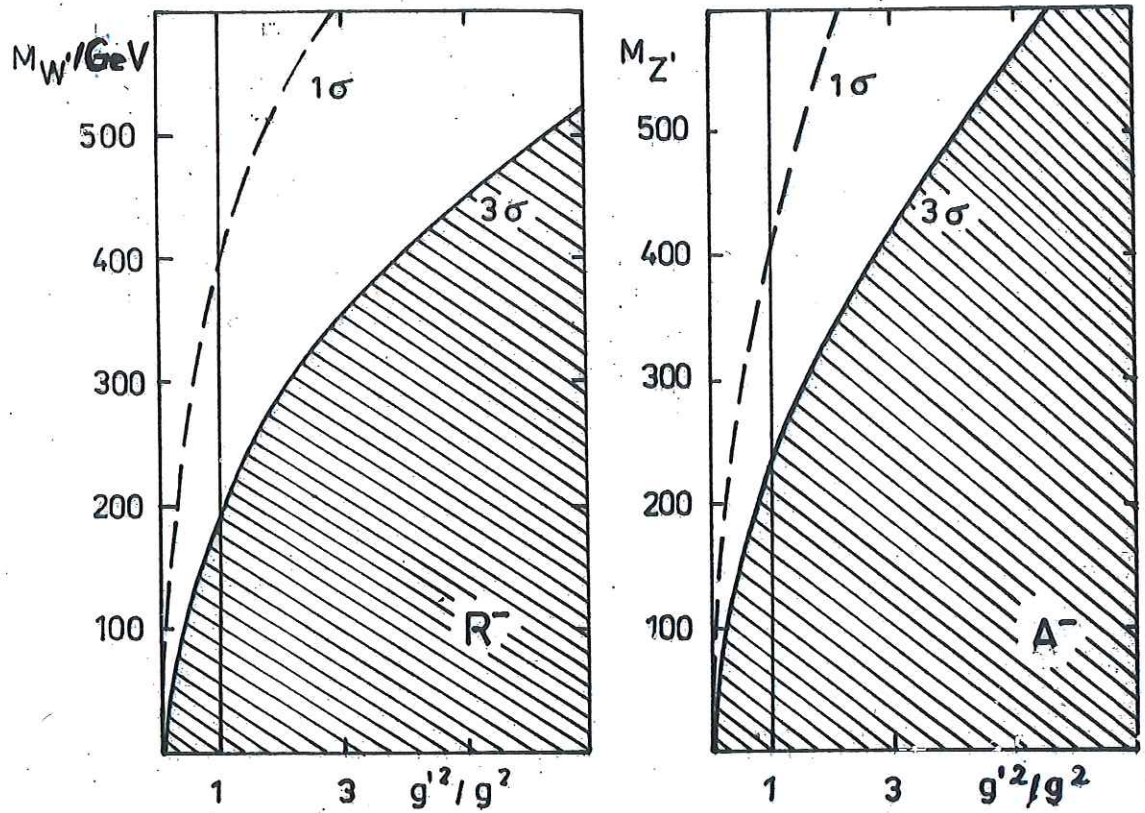


Fig. 8

$$K(Q^2) = g^2 \frac{M^2}{M^2 + Q^2} + g'^2 \frac{M'^2}{M'^2 + Q^2}$$

$$M = M_W, M_Z$$

D. COMPOSITENESS SCALE OF W'S & Z'S

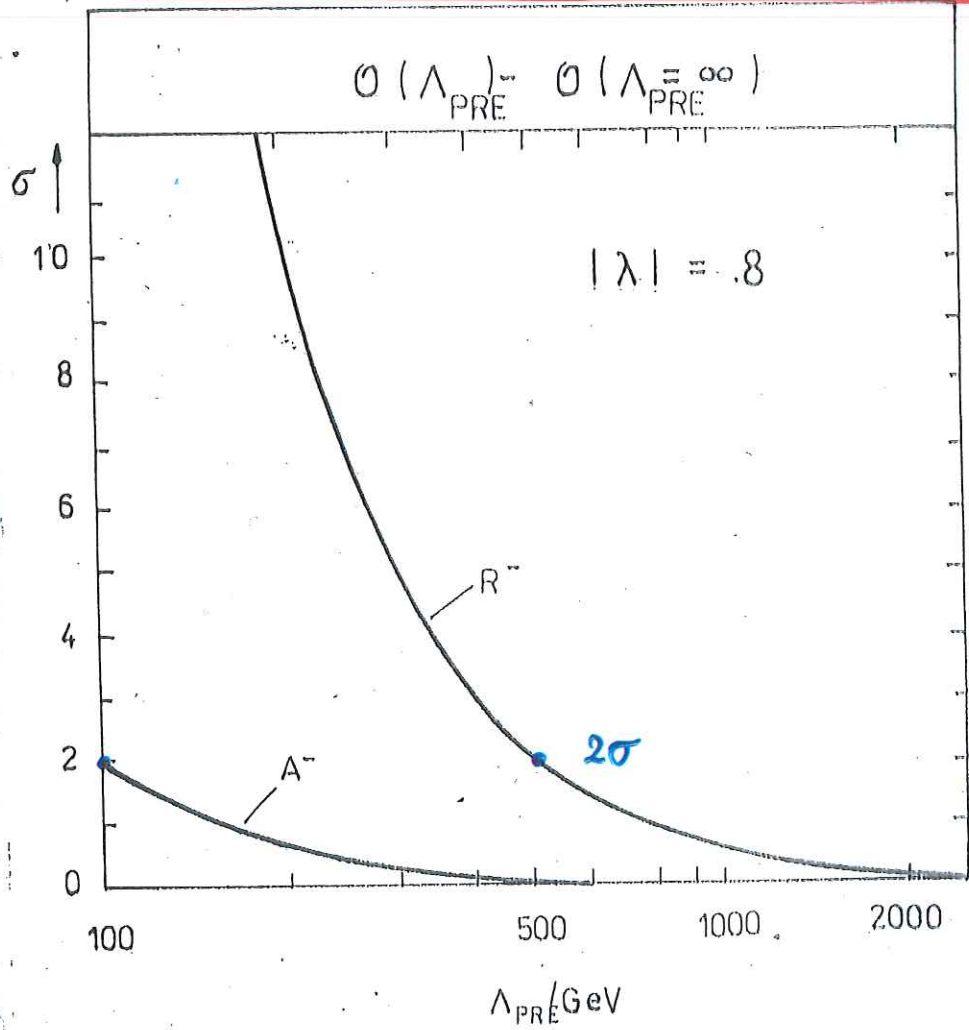


Fig. 10

$$K_{Z,W} = K_{Z,W} \cdot 1 / (1 + Q^2 / \Lambda_{PRE}^2)$$

MEASUREMENT OF STRUCTURE FUNCTIONS: $e^{\pm}p$

NC: $B_{\pm}(\lambda) = \frac{1}{2} [\sigma_{NC}^{+}(\lambda) \pm \sigma_{NC}^{-}(\lambda)]$ (KLEIN, RIEMANN '8)

$$B_{+} \cdot \frac{Q^4 x}{2\pi\alpha^2 Y_{+}} = F_2 + f_1(k_z, G_2, H_2, \lambda)$$

↑
CONTR. VANISHES FOR
 $\lambda = v/q \approx 0$

- CONTAINS ONLY SINGLET FUNCTIONS
- COUPLINGS, k 's CAN NOT BE SEPARATED AS FACTORS ($q \cdot e^{\pm}D$)

FIG.

$$B_{-} \cdot \frac{Q^4 x}{2\pi\alpha^2 Y_{-}} \cdot \frac{1}{k_z} = xG_3 + f_2(k_z, xH_3, \lambda)$$

↑
CONTR. VANISHES FOR
 $\lambda = 2vq/(v^2 + q^2)$

- CONTAINS ONLY NON-SINGLET FUNCTIONS

FIG.

CC: 4 STRUCTURE FUNCTIONS
 W_2^{\pm}, xW_3^{\pm} VS. 2 CROSS SECTION

(B.) MEASUREMENT OF STRUCTUREFUNCTIONS

- MORE SIZEABLE EFFECTS: (NO RATIOS OF CROSS SECTIONS)
- RC_{1L} HAS TO BE CORRECTED FOR TO MEASURE
EG. F_2 , I.E. $\Delta\sigma_{1L}(x,y) / \sigma_{1L}(x,y) \lesssim 30\%$
- CONSTRAINS (x,y) TO MEDIUM VALUES

RADIATIVE CORRECTIONS

∃ CALCULATIONS OF: $\frac{d\sigma^{\text{tr}}}{dx dy} \cdot (1 + RC_{\text{IL}}(x, y))$ BY:

- STILL NOT IN AGREEMENT
- BÖHM & SPIESBERGER (WÜRZBURG UNIV.)
 - BARDIN, FEDORENKO, RIEMANN et al. (DUBNA)

EARLY '80ies

RECALCULATION ≤ JUNE 1987

- MÖHRING, KRIPFGANZ (LEIPZIG UNIV.) (PARTS ONLY)
- ⋮

EFFECTS ON:

(A.) TESTS OF EWTHY: $\left\{ \begin{array}{l} A_{\pm} \rightarrow A_{\pm}^{\text{tr}} (1 + \delta_{\text{IL}} A_{\pm}) \\ R_{\pm} \rightarrow R_{\pm}^{\text{tr}} (1 + \delta_{\text{IL}} R_{\pm}) \end{array} \right.$

(DIFFERENCE) & RATIOS

$$S_{\text{eH}}^2 \rightarrow S_{\text{eH}}^2 + \Delta S^2(m_t, m_H, x_c, y_c; \bar{x}_i^q)$$

$$M_W^{\text{tr}} \rightarrow M_W^{\text{tr}} + \Delta M_W(\dots)$$

$m_t, m_H(x, y) - \text{CU}$
~ (PARTON DISTR'S)

LET $\{Q_i \pm \Delta_i\}_{i=1}^N$ BE MEASUREMENTS
OF Q . $\langle Q \rangle \pm \Delta$ IS THEN OBTAINED BY

$$\langle Q \rangle \pm \Delta = \sum_i \frac{Q_i}{\Delta_i^2} / \sum_i \frac{1}{\Delta_i^2} \pm 1 / \left(\sum_i \frac{1}{\Delta_i^2} \right)^{1/2}$$

Q_i ENTERS $\langle Q \rangle$ WITH THE WEIGHT $\frac{1}{\Delta_i^2}$

THE FITTING PROCEDURE

- CALCULATE : A., R. FOR THE x & Q^2 -BINS
- DEFINE : $\chi^2 = \sum_{\text{BINS}} \left[\frac{Q_i - Q(i; EW_1, EW_2, EW_3)}{\Delta Q_i} \right]^2$
- MINIMIZE : χ^2 FIXING : EW_2 & EW_3 VARYING EW_1 .



$$EW_1 \pm \Delta EW_1$$

$$\Delta EW_1 \approx \left[2 / \left[\partial^2 \chi^2 / \partial EW_1^2 \right] \right]^{1/2} = \frac{\partial EW_1}{\partial Q} \left(\sum_i \Delta Q_i^{-2} \right)^{-1}$$

THE VALUE OF $\frac{\partial Q}{\partial EW}$ DOES STRONGLY DEPEND ON $Q = Q(EW)$.

(A. BÖHM, JB MKTR, HIOKI)

$e^+e^- \rightarrow s_0^2$

: THE RATIO OF A DISPLACEMENT (BIAS) OF E

$$Q_{\text{meas}} \rightarrow Q_{\text{nature}} + \Delta$$

$$\delta E_{\text{meas}} = \frac{\partial E}{\partial Q} \Delta$$

AND ΔE_{STAT} DEPENDS NOT ON $\frac{\partial E}{\partial Q}$ (J. KUHN)

EW-TESTS BEYOND GWS

SUMMARY: (DETAILS IN: PHE 87-3)

A: $g_{(0MS)}$ AS INDEPENDENT PARAMETER

$g - \sin^2\theta$: $\delta g = .01$, $\delta \sin^2\theta = .006$, R
 ONLY g : $\delta g = .003$

B: I_3^R :

e	$A^+ (.8)$.030
μ	$A^+ (.8)$.015
d	$A^- (.8)$.034

C:

W'	$R^- \geq 3\sigma$	SENSITIVE	$M_{W'} < 200 \text{ GeV}$
Z'	$A^- \geq 3\sigma$		$M_{Z'} < 240 \text{ GeV}$

D: \wedge_{PRE}^{WIZ}

R^-	$\geq 2\sigma$	SENSITIVE	$\Lambda < 500 \text{ GeV}$
A^-			$\Lambda < 100 \text{ GeV}$

DEUTERONS AT HERA

NC:

$$F_2^{eN} = \frac{Q_u^2 + Q_d^2}{2} \sum_i (xq + x\bar{q}) + \frac{Q_u^2 - Q_d^2}{2} 2x(c-s)$$

SIMILAR FACTORIZATION OF COUPLINGS
FOR : $B_+(\lambda) / Y_+ \cdot Q^4 x / 2\pi\alpha^2$



- CONTAINS SINGLET FUNCTIONS ONLY
- PROPAGATORS FACTORIZE IF ONE CORRECTS FOR (C-S)

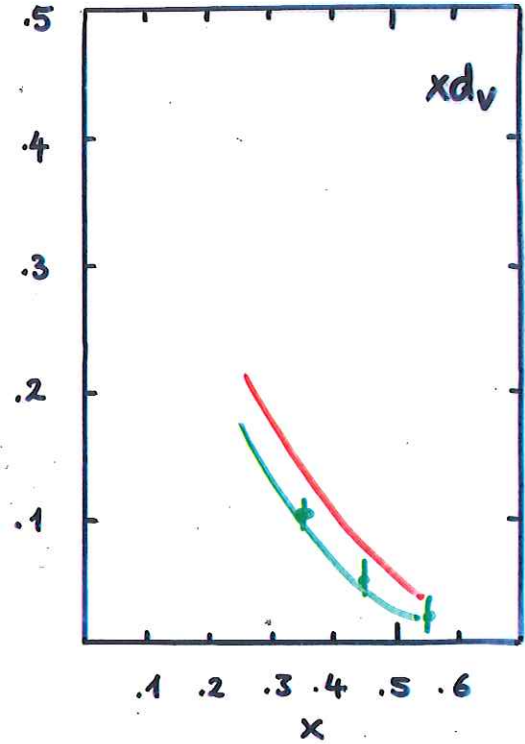
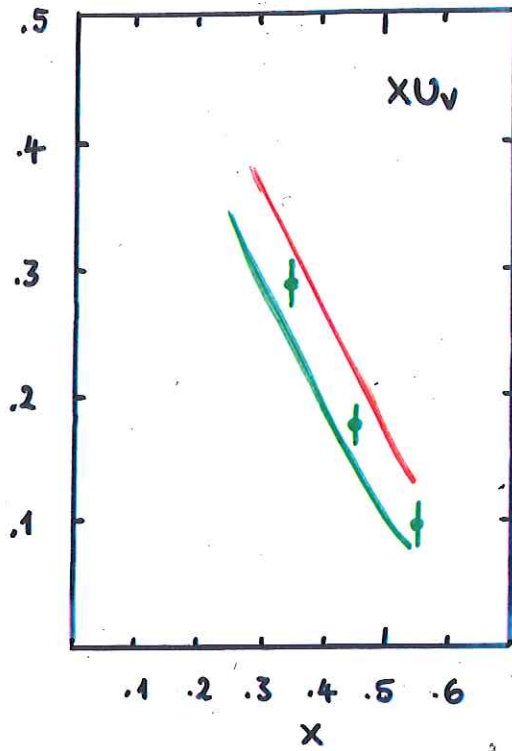
FIG.

$$xG_3 = (Q_u a_u + Q_d a_d) x (u_v + d_v)$$

SIMILARLY : $B_-(\lambda) / Y_- \cdot k_z \cdot Q^4 x / 2\pi\alpha^2$

FIG.

VALENCE APPROXIMATIONS



$\sigma_{nc}^{\pm}, \sigma_{cc}^{\pm}, x \geq 0.3$

RECONSTRUCTION OF QUARK DISTRIBUTIONS : $e^\pm p$

9 STRUCTURE FUNCTIONS

6 QUARK DISTRIBUTIONS

4 CROSS SECTIONS. $\sigma_{NC}^\pm, \sigma_{CC}^\pm$

$$\text{MAP: } \begin{pmatrix} \sigma_{NC}^+ \\ \sigma_{NC}^- \\ \sigma_{CC}^+ \\ \sigma_{CC}^- \end{pmatrix} \xrightarrow{A_{ij}} \begin{pmatrix} \Sigma u \\ \Sigma d \\ \Sigma \bar{u} \\ \Sigma \bar{d} \end{pmatrix}$$

$$\det_4 A_{ij} \sim K(Q^2) [1 - (1-y)^4]$$

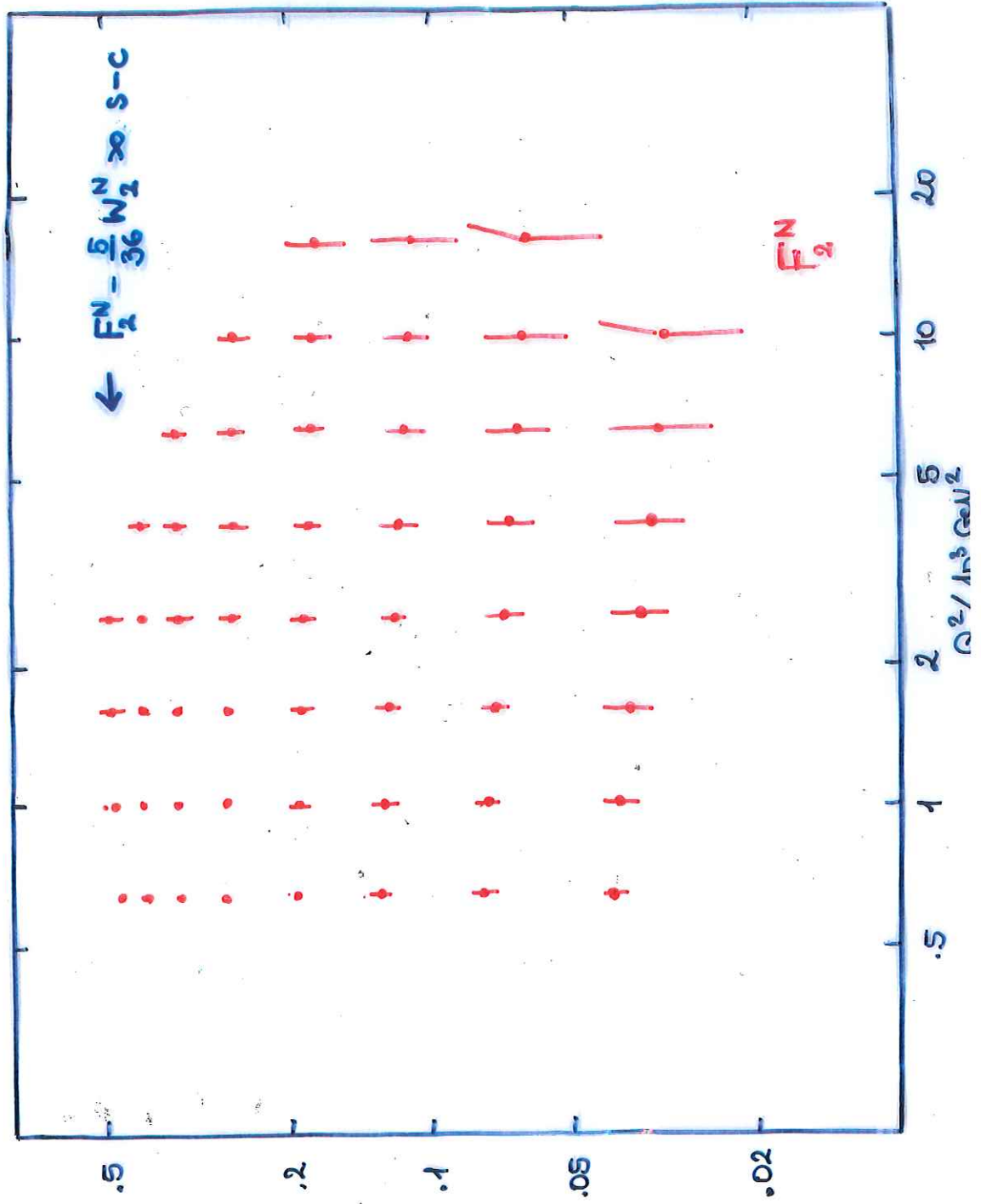
SINGULARITY AT LOW Q^2, y

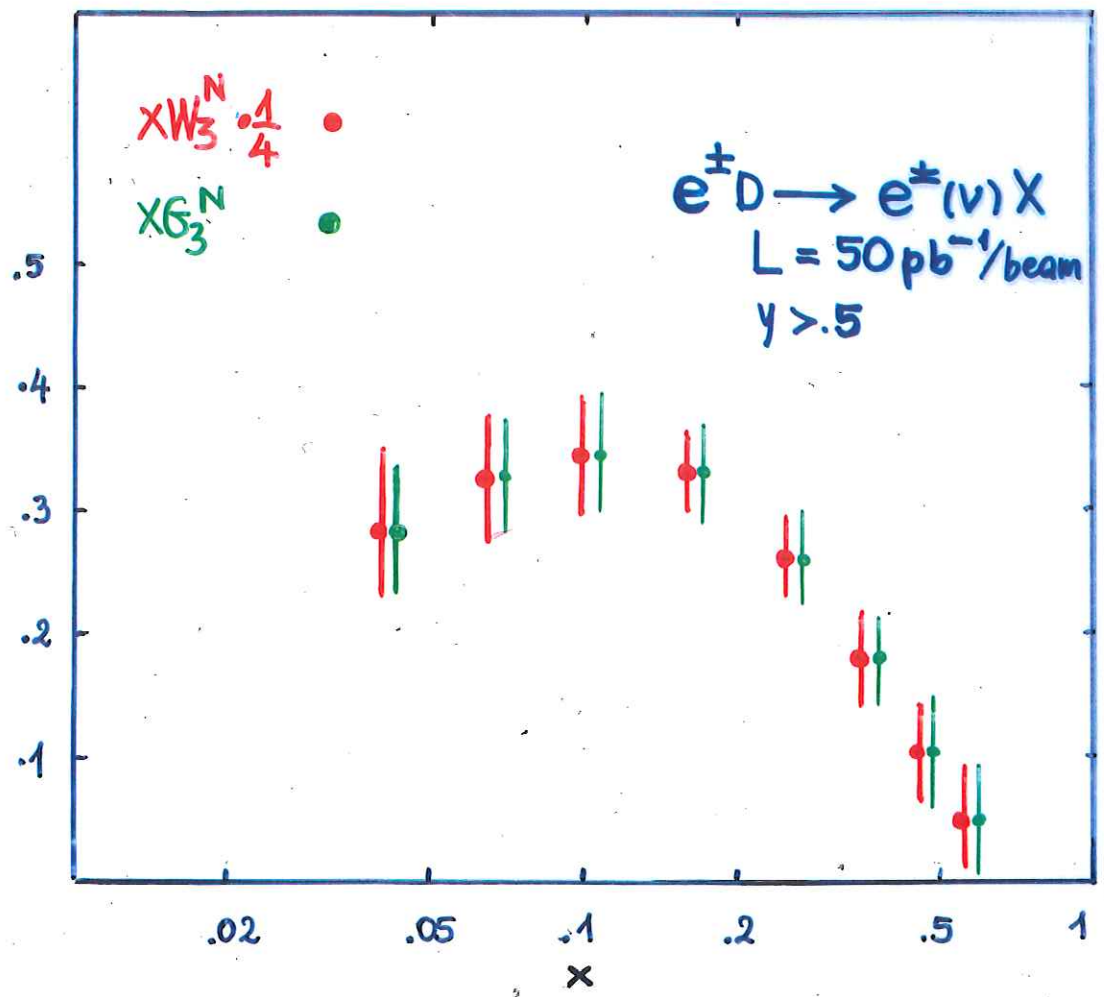
$$\sigma_{NC}^+ \rightarrow \sigma_{NC}^-$$

USE: 3 CROSS SECTIONS: $\sigma_{NC}^-, \sigma_{CC}^+, \sigma_{CC}^-$

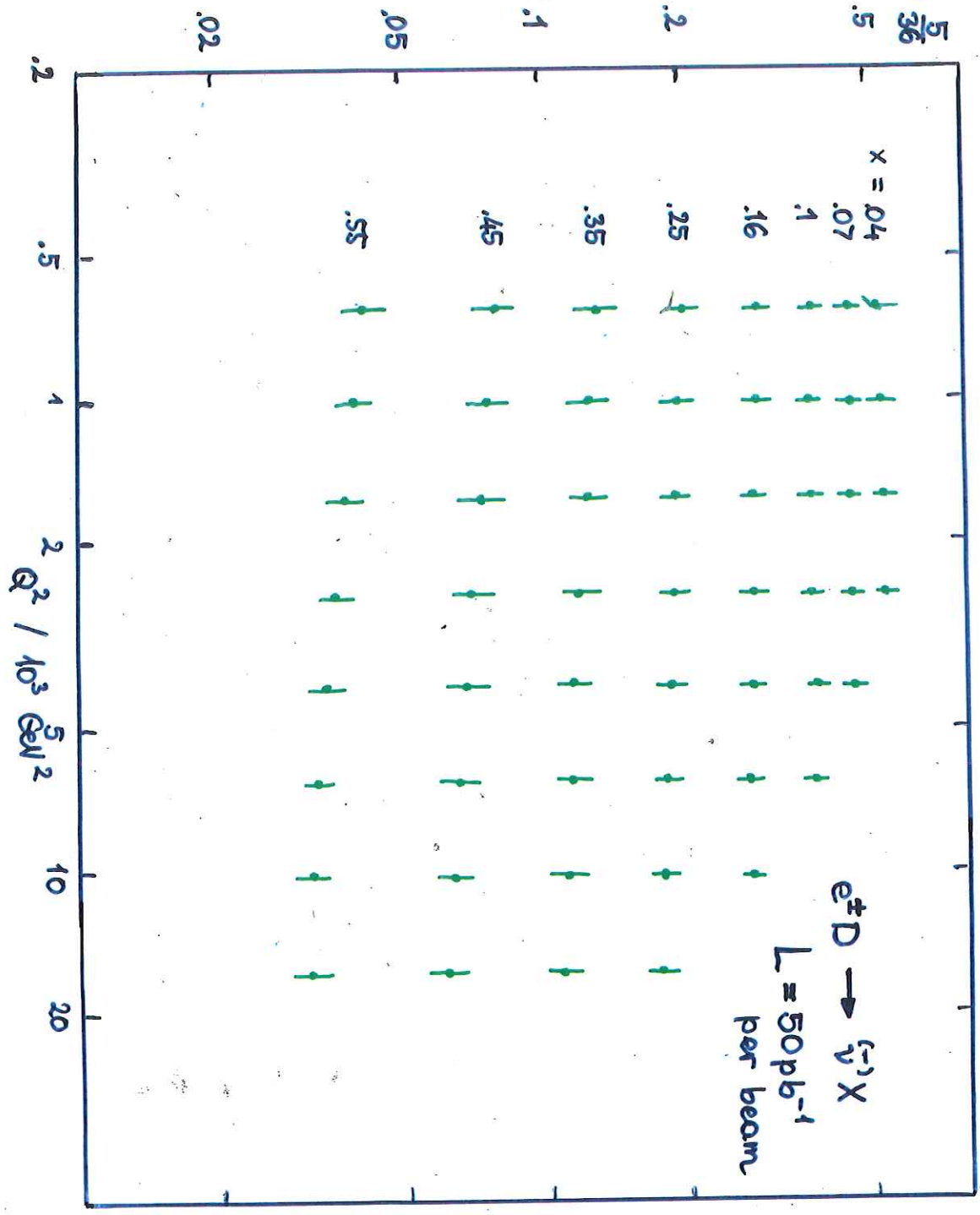
ASSUME: $u_s = d_s = \text{const} \cdot s$, "ignore" c, b

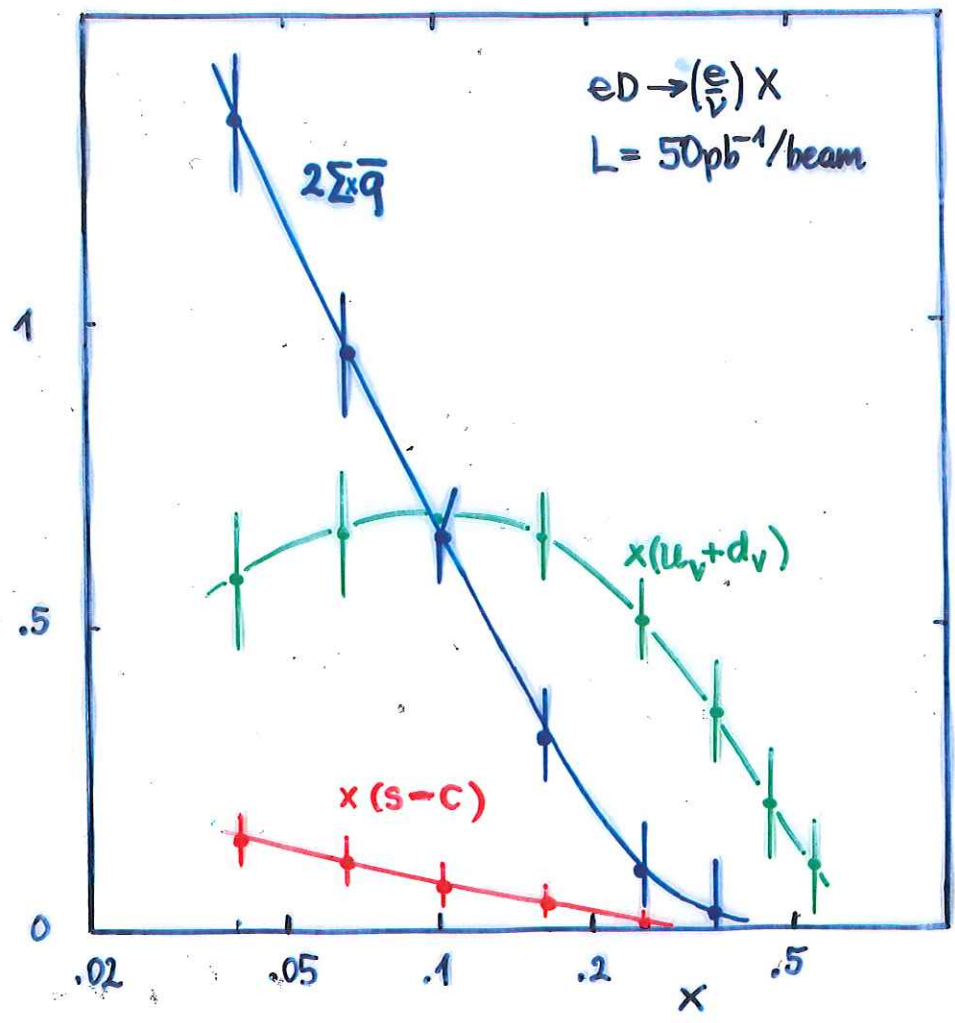
$$\longrightarrow u_v, d_v, s(u_s)$$





$W_2^N \cdot \frac{5}{36}$





CONCLUSIONS

- A STUDY HAS BEEN PRESENTED OF HERA'S POTENTIAL TO TEST THE STANDARD MODEL AND POSSIBILITIES TO EXTRACT QUARK DISTRIBUTIONS.
- R_L APPEARS TO BE THE BEST SUITED QUANTITY TO MEASURE $\sin^2\theta$, M_W , m_t IN THE GWS-MODEL AND TO TEST MODEL EXTENSIONS AS: $\exists W', Z', \Lambda_{PRE} < \infty$
- POLARIZATION IS, FROM THIS POINT OF VIEW, LESS IMPORTANT THAN A PRECISE JET-MEASUREMENT AND A GOOD SEPARATION OF NC & CC-EVENTS
- $\mathcal{L} = 200 \mu\text{b}^{-1}$ MAY YIELD:

$$\Delta S_\theta^2 = \pm 0.002, \quad \Delta M_W = \pm 100 \text{ MeV} \quad (\text{STAT.})$$

THE MAJOR UNCERTAINTY-OF UP TO THE ORDER OF STAT. PRECISION-IS INDUCED BY PARTON DISTRIBUTION UNCERTAINTIES.

- BEYOND THE GWS MODEL ONE EXPECTS TO OBTAIN:

$$\delta g = \pm 0.003, \quad \delta I_3^{R_0} = .03, \quad \left. \begin{array}{l} M'_W \leq 200 \text{ GeV} \\ M'_Z \leq 240 \text{ GeV} \end{array} \right\} \begin{array}{l} 3\sigma \\ \text{SENSITIVITY} \\ \text{FOR } g = g' \end{array}$$

$$\Lambda_{\text{PRE}} \leq 500 \text{ GeV} \quad 2\sigma \text{ SENSITIVITY}$$

- DIRECT MEASUREMENT OF PARTON DISTRIBUTIONS IS ONLY POSSIBLE IN THE VALENCE REGION.
- UNDER SOME ASSUMPTIONS $\langle xq \rangle_{Q^2}$ CAN BE EXTRACTED WITH $O(25\%)$ ERRORS.
- INTERFERENCE EFFECTS CAN BE USED TO MEASURE xG_3 AT THE SAME LEVEL.
- THE DEUTERON OPTION AT HERA WOULD ALLOW TO MEASURE $x(u_v + d_v), xS, x(S-C)$ AS A FUNCTION OF x